

Method of Configuration Interaction- Simple examples of HF failurea.) dissociation of  $\text{H}_2$ b.) ionization of  $\text{N}_2$ c.) existence of  $\text{H}^-$ - Correlation energy

$$E_{\text{cor}} = E - E^{\text{HF}} \quad (\text{nonrelativistic})$$

HF is a variational method with an exact nonrelativistic electronic Hamiltonian on a poor set of Slater determinants.

→ HF gives the upper guess of the energy  $\rightarrow E_{\text{cor}} \leq 0$

(equality in a case of 1-electron system)

- Idea of configuration interaction

- HF spin-orbitals form a complete basis set in 1-electron Hilbert space  $H$  of dimension  $2K$ .  $N$ -tuple direct product of this basis forms a basis in  $\underbrace{H \times H \times \dots \times H}_{N \times}$ . Here  $N$  is the number of electrons.
- From this  $N$ -order direct product of the HF spin-orbitals we choose a subset satisfying the Pauli exclusion principle. Namely, the antisymmetric basis set products without repeated 1-electron basis elements (spin-orbitals). These are exactly excited Slater

determinants of HF spinorbitals (occupied + virtual).

We have  $\binom{2K}{N}$  of such ~~of~~ configurations.

- Moreover, we define an "order of importance" on the product basis set elements. This order comes in form of groups according to the number of excitations.

$$\{B\}_N = |1\psi_0\rangle \cup |1\psi_a^r\rangle \cup |1\psi_{ab}^{rs}\rangle \cup |1\psi_{abc}^{rst}\rangle \cup \dots$$

reference	single excitations	double excitations	triple excitations	...
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- Emphasize that CI uses 2 different ordering at different levels:

a.) Ordering of 1-electron basis according to 1-electron mean energy (in contrast to simple GS-orthogonalization)

b.) Ordering of N-products where we assume singly excited determinants more important than triply excited ...

- With this N-electron basis set the  $^{\text{CI}}$  ground state can be written:

$$|\phi_0\rangle = c_0|1\psi_0\rangle + \sum_{ar} c_a^r |1\psi_a^r\rangle + \sum_{\substack{a < b \\ r < s}}^{rs} c_{ab}^r |1\psi_{ab}^{rs}\rangle + \sum_{\substack{a < b < c \\ r < s < t}}^{rst} c_{abc}^r |1\psi_{abc}^{rst}\rangle + \dots$$

$$|\phi_0\rangle = c_0|1\psi_0\rangle + \sum_{ar} c_{ar}^r |1\psi_a^r\rangle + \left(\frac{1}{2!}\right)^2 \sum_{\substack{a < b \\ r < s}}^{rs} c_{ab}^r |1\psi_{ab}^{rs}\rangle + \left(\frac{1}{3!}\right)^2 \sum_{\substack{a < b < c \\ r < s < t}}^{rst} c_{abc}^r |1\psi_{abc}^{rst}\rangle + \dots$$

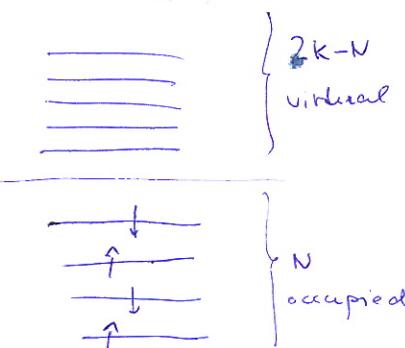
- How many of m-order excitations do we have?

(N .. number of electrons, 2K basis spin-elements)

$$= \binom{N}{m} \binom{2K-N}{m}$$

Number of possibilities to choose m-electrons from N occupied orbitals.

Number of possibilities to put m-electrons into 2K-N virtual orbitals



Method of CI  
(continued)

Variations with respect to  $c_0, c_a^r, c_{ab}^{rs}, c_{abc}^{rst}, \dots$  coefficients leads to diagonalization of the CI Hamiltonian matrix:

$$\langle \Phi_0 | = c_0 | 0 \rangle + c_s | S \rangle + c_D | D \rangle + c_T | T \rangle + \dots$$

CI Hamiltonian

HF theory	$\langle 0   H   0 \rangle$	0	$\langle 0   H   D \rangle$	0	0	0
CIS only used for inaccurate excited states	0	$\langle S   H   S \rangle$	$\langle S   H   D \rangle$	$\langle S   H   T \rangle$	0	0
CISD (most common CI method)	$\langle D   H   0 \rangle$	$\langle D   H   S \rangle$	$\langle D   H   D \rangle$	$\langle D   H   T \rangle$	$\langle D   H   R \rangle$	0
CISDT	0	$\langle T   H   S \rangle$	$\langle T   H   D \rangle$	$\langle T   H   T \rangle$	$\langle T   H   Q \rangle$	...
Full CI	0	0	$\langle Q   H   D \rangle$	$\langle Q   H   T \rangle$	$\langle Q   H   Q \rangle$	

Two variants of the CI idea:

CAS CI and MCSCF

1.) CAS CI one-electron basis (the set of HF orbitals) is severely reduced. This defines the active space. Then the FCI is carried ~~out~~ in the resulting active space.

CAS CI = complete active space CI.

## 2.) MCSCF (Multi-configuration self-consistent field)

$$|\Psi_{\text{MCSCF}}\rangle = \sum_I c_I |\Psi_I\rangle$$

$|\Psi_I\rangle$  are excited determinants in the active space. Active space is again a reduced 1-electron set of HF orbitals. In contrast to CASSCF, where the excited determinants are fixed, in MCSCF method also the ~~state~~ 1-electron basis <sup>forming</sup> in the active space is optimized:

- - New  $c_I$  coefficients
- $c_I$  coefficients give HF equations
- New  $\psi_i$  from HF
- New  $|\Psi_I\rangle$ , new FCI matrix
- New  $c_I$  after the FCI diagonalization

Note: MCSCF has more applications:

- 1.) It gives decent results on its own
- 2.) It is used to generate reduced density matrix, that can be diagonalized → natural orbitals. These orbitals can define a different excitation space (instead of the canonical HF orbitals) for a following CISD calculations. The convergence is usually faster in the natural orbitals.