

Kuchařka ke kanonickému ensemble

1, $H(p, q)$ & fázový prostor

2, $Z_c(\beta, V, N) = \int d\Gamma_N e^{-\beta H(p, q)}$

$\frac{1}{N!}$? ... exkluzivita F

3, $F(T, V, N) = - \frac{1}{\beta} \log Z_c(\beta, V, N) / \beta = \frac{1}{k_B T}$

4, termodynamika ...

$S = - \left(\frac{\partial F}{\partial T} \right)_{V, N}$

4b, $u = \langle H \rangle = - \frac{\partial \log Z_c}{\partial \beta}$ & $\langle (\Delta H)^2 \rangle = \frac{\partial^2 \log Z_c}{\partial \beta^2}$

1, Dvoučásticový systém

$N \gg 1$ nezávislých částic, každá se může nacházet ve

$|0\rangle \leftrightarrow \epsilon_0 = 0$

$|1\rangle \leftrightarrow \epsilon_1 = \epsilon$

a, $S(T, V, N)$

b, $C(T)$

$Z_c = \sum_{\text{stavy}} e^{-\beta H}$

\downarrow
 $\{n_1, \dots, n_N\}$ $n_i \rightarrow 0$
 $\rightarrow 1$

$H = \sum_{i=1}^N H_1(i)$

$H_1 = |0\rangle 0 \langle 0| + |1\rangle \epsilon \langle 1|$

$e^{-\beta H} = \prod_i e^{-\beta H_1(i)}$

$\sum_{\text{stavy}} = \sum_{i=1}^N \sum_{j=0}^1$

$Z_c = \prod_{i=1}^N Z_1 = (Z_1)^N$

$\sum_{i=1}^2 \sum_{j_i=0}^1 = e^{-\beta(0+0)} + e^{-\beta(\epsilon+0)} + e^{-\beta(0+\epsilon)} + e^{-\beta(\epsilon+\epsilon)}$

$= e^{-\beta 0} (e^{-\beta 0} + e^{-\beta \epsilon}) + e^{-\beta \epsilon} (e^{-\beta 0} + e^{-\beta \epsilon})$

$= (e^{-\beta 0} + e^{-\beta \epsilon}) (e^{-\beta 0} + e^{-\beta \epsilon})$

$Z_{i=2}$

$Z_{i=1}$

• faktORIZACE PARTIČNÍ FUNKCE

pro nezávislé skupně volnosti je $Z_c = \prod_i Z_i$

$$Z_1 = \sum_{\substack{\text{1 část} \\ \text{stavu}}} e^{-\beta H} = e^{-\beta 0} + e^{-\beta \epsilon} = 1 + e^{-\beta \epsilon}$$

$$Z_c = (1 + e^{-\beta \epsilon})^N \Rightarrow F(\beta, V, N) = -k_B \log Z_c = -N k_B \log(1 + e^{-\beta \epsilon})$$

$$U = \langle H \rangle = - \frac{\partial}{\partial \beta} \log Z_c = -N \frac{\partial}{\partial \beta} \log(1 + e^{-\beta \epsilon})$$

$$= - \frac{N}{1 + e^{-\beta \epsilon}} \cdot (-\epsilon e^{-\beta \epsilon}) = N \frac{\epsilon e^{-\beta \epsilon}}{1 + e^{-\beta \epsilon}} = N \epsilon p(\epsilon)$$

0 · p(0) + ε p(ε)

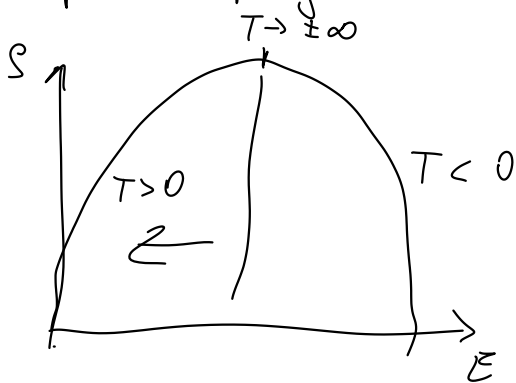
$$\langle 1 | \frac{1}{Z_c} e^{-\beta H} | 1 \rangle = \frac{e^{-\beta \epsilon}}{1 + e^{-\beta \epsilon}}$$

mikro

$$\boxed{\epsilon = \frac{N \epsilon}{1 + e^{\epsilon/k_B T}}}$$

$$C = \frac{\partial U}{\partial T} = \dots \text{Schottky}$$

• Záporné teplo:



$$\frac{1}{T} = \frac{\partial S}{\partial E}$$

rezervuar ↔ měnící systém

T > 0 !

Ideální plyn kanonicky

$$H = \sum_{i=1}^N H_i = \sum_{i=1}^N \frac{p_i^2}{2m}$$

$$Z_c = \int d^3x^N e^{-\beta H} = \frac{1}{N!} \prod_{i=1}^N \int \frac{d^3p_i d^3q_i}{h^3} e^{-\beta \frac{p_i^2}{2m}} = \frac{1}{N!} \left(\int \frac{d^3p d^3q}{h^3} e^{-\beta \frac{p^2}{2m}} \right)^N$$

$$\Rightarrow Z_c = \frac{1}{N!} (Z_1)^N$$

↑
faktORIZACE PARTIČNÍ FUNKCE

$$E = \sum_{i=1}^{3N} \frac{p_i^2}{2m} + \sum_{i < j} \Phi(q_i, q_j)$$

$$Z_1 = \frac{1}{h^3} \int d^3q d^3p e^{-\beta \frac{p^2}{2m}} = \frac{V}{h^3} \int d^3p e^{-\beta \frac{p^2}{2m}} = \left(\begin{array}{l} \text{8 fe'x.} \\ \text{swiadczanie} \end{array} \right)$$

$$= \frac{4\pi V}{h^3} \int_0^\infty dp p^2 e^{-\beta \frac{p^2}{2m}} = \left(\begin{array}{l} p \sqrt{\frac{1}{2m}} = x \\ dp = dx \sqrt{\frac{2m}{1}} \end{array} \right)$$

$$= \frac{4\pi V}{h^3} \sqrt{\frac{2m}{1}}^3 \int_0^\infty dx x^2 e^{-x^2}$$

$$Z_1 = \frac{V}{h^3} \left(\int_{-\infty}^{\infty} dp_i e^{-\frac{\beta}{2m} p_i^2} \right)^3 = \frac{V}{h^3} \left(\frac{2\pi m}{\beta} \right)^{3/2} = V \left(\frac{2\pi m}{h^2 \beta} \right)^{3/2}$$

$$Z_c = \frac{1}{N!} (Z_1)^N = \frac{1}{N!} V^N \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3N}{2}} (k_B T)^{3/2}$$

$$\Rightarrow F = -k_B T \log Z_c = -N k_B T \log \left[V \beta^{-3/2} \left(\frac{2\pi m}{h^2} \right)^{3/2} \right] + \log N!$$

$$= -N k_B T \log \left(T^{3/2} \left(\frac{V}{N} \right) \right) - \frac{3}{2} N k_B \log \left(\frac{2\pi m k_B}{h^2} \right) - N$$

$$F(T, V, N) = -N k_B T \log \left(T^{3/2} V N^{-1} \right) + N T f_0$$

$$\psi(|\vec{v}|=v) \sim e^{-\alpha v^2} \leftarrow \epsilon_k$$

Co: $F(T, V, N)$ vs. f_0
 v mikrokanon. formula cisna

• $S(T, V, N)$ vs. $S_e(T, V, N)$

Ideální plynu v grav. poli

$$E\langle H \rangle = U + \langle H_I \rangle$$

$$H = \sum_i H_1(i)$$

$$H_1 = \frac{p^2}{2m} + mgz = \frac{p^2}{2m} + Fz$$

$$= H_0 + H_I$$



$$\Sigma = \int dx dy$$

$$Z_c = \frac{1}{N!} (Z_1)^N$$

$$U = \langle H_0 \rangle$$

$$Z_1 = \frac{1}{h^3} \int d^3p d^3q e^{-\beta H_1} = \sum \left(\frac{2\pi m}{h^2 \beta} \right)^{3/2} \int_0^\infty dz e^{-\beta mgz}$$

$$\left(\sqrt{2\pi m k_B T} = \lambda_T \dots \text{termální vlnová délka} \right)$$

$$= \sum \left(\frac{2\pi m}{h^2 \beta} \right)^{3/2} \frac{e^{-\beta mgz}}{\beta mg} \Big|_0^\infty = \underbrace{\sum \left(\frac{2\pi m}{h^2} \right)^{3/2} \beta^{-5/2} \frac{1}{mg}}_{\sim N} = Z_1$$

$N \log N - N$

$$\Rightarrow F = -\frac{1}{\beta} \log Z_c = -\frac{N}{\beta} \log Z_1 + \frac{1}{\beta} \log N!$$

$$F = -\frac{N}{\beta} \log \left(\frac{\Sigma}{N} \beta^{-5/2} \right) - \frac{N}{\beta} \log \left[\left(\frac{2\pi m}{h^2} \right)^{3/2} \frac{1}{mg} \right] - \beta N$$

$$F(T, \Sigma, N) = -N k_B T \log \left(T^{5/2} \frac{\Sigma}{N} \right) - N k_B T \log \left[\frac{k_B^{5/2}}{mg} \left(\frac{2\pi m}{h^2} \right)^{3/2} \right] - k_B T N$$

$$\omega = \frac{1}{Z_c} e^{-\beta H_0 - \beta mgz}$$

$$G(T, \Sigma, mg, N)$$

$$Z_c = \int d^3x \omega e^{-\beta H_0 - \beta mgz}$$

Legendre v TD potenciálu

$$S(u, v, N) \rightarrow S \left[\frac{1}{T} \right] \left(\frac{1}{T}, v, N \right)$$

\downarrow
 x, y, z

$$\rightarrow S \left[\frac{1}{T}, \frac{F_z}{T} \right] \left(\frac{1}{T}, \frac{F_z}{T}, y, x, N \right)$$

$$\langle z \rangle = \int dz z \omega = -\frac{1}{\beta m} \frac{\partial}{\partial g} \log Z_1$$

Teplotná kapacita

$$C_{\Sigma, N} = T \left(\frac{\partial S}{\partial T} \right)_{\Sigma, N} = - T \left(\frac{\partial^2 F}{\partial T^2} \right)_{\Sigma, N}$$

$$F(T, \Sigma, N) = - N k_B T \log \left(T^{5/2} \frac{\Sigma}{N} \right) - N k_B T \log \left[\frac{k_B^{5/2}}{m g} \left(\frac{2\pi m}{h^2} \right)^{3/2} \right] - k_B T N$$

$$F \approx - N k_B T \log T^{5/2}$$

$$\frac{\partial^2 F}{\partial T^2} = 0$$

$$\frac{\partial F}{\partial T} = - N k_B \log T^{5/2} - \frac{5}{2} N k_B \frac{T}{T}$$

$$E = \langle H \rangle = \frac{5}{2} N k_B T$$

$$\frac{\partial^2 F}{\partial T^2} = - \frac{5}{2} N k_B \frac{1}{T} \Rightarrow C_{\Sigma} = \frac{5}{2} N k_B = \frac{3}{2} N k_B + \frac{2}{2} N k_B$$

kin. energie

grav. pole

C_V vs. C_P

kvadratických /
st. volnosti

$T \uparrow \langle z \rangle \uparrow$

$$C_V = \frac{f}{2} k_B$$

$\langle z \rangle$:

$$\left[\sum \left(\frac{2\pi m}{h^2} \right)^{3/2} \beta^{-5/2} \frac{1}{m g} = Z_1 \right]$$

$$\langle z \rangle = \int dz z \omega = - \frac{1}{\beta m g} \frac{\partial}{\partial g} \log Z_1$$

$$\log Z_1 = - \log g + \log \left(\sum \left(\frac{2\pi m}{h^2} \right)^{3/2} \beta^{-5/2} \frac{1}{m} \right)$$

$$\Rightarrow - \frac{1}{\beta m g} \frac{\partial}{\partial g} \log Z_1 = \frac{1}{\beta m g} \Rightarrow \langle z \rangle = \frac{k_B T}{m g} \quad \uparrow \text{ s } T \text{ roste lineárne}$$

$$U = \langle H_0 \rangle = \langle H \rangle - m g \langle z \rangle = - \frac{\partial \log Z_c}{\partial \beta} - N k_B T$$

$$H = H_0 + m g z$$

$$\downarrow$$

$$\frac{5}{2} N k_B T$$