



czechLISA

Czech participation in the LISA mission

Alexandre M. Pombo



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Boson Stars that mimic Black Holes

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Introduction: Black Hole mimicker

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- The shadow is associated with the LR and illumination source
- LRs are a generic feature of BHs
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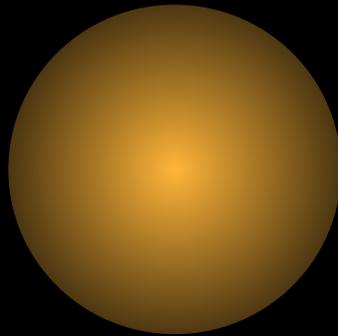
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Can a dynamically robust, horizonless object mimic a BH?

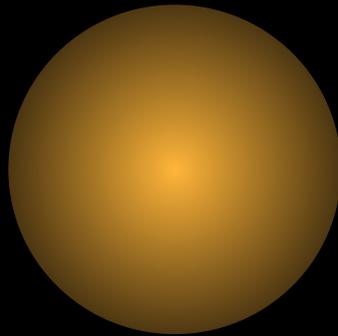
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- A Boson Star is an hypothetical astronomical object made out of bosons
- For this stars to exist, one needs a stable boson
- Compact BSs are usually contain a massive complex scalar fields with $U(1)$ global symmetry
- They are everywhere regular lumps of scalar or vector boson



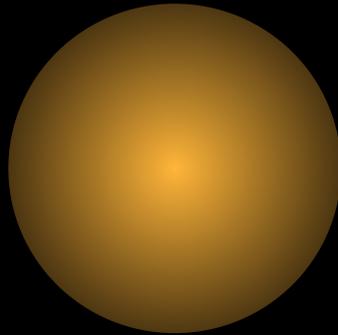
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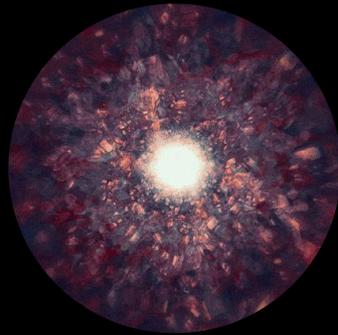
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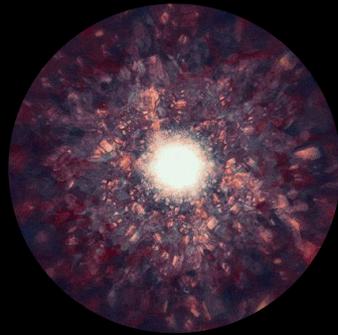
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- Boson Stars have been shown to be able to form dynamically
- And have been proposed as candidate dark matter objects
- They would interact very weakly with electromagnetic radiation
- The gravity of a compact BS can bend light and create an accretion disk



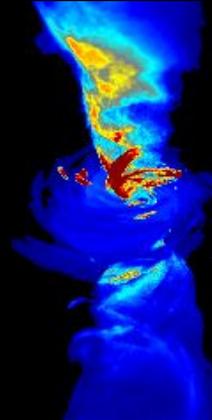
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The Model

The model: Lagrangean

- The Einstein-matter action, with a spin- $s = 0, 1$ classical field

$$\mathcal{S} = \int d^4x \sqrt{-g} \left[\frac{R}{16\pi G} + \mathcal{L}_s \right] ,$$

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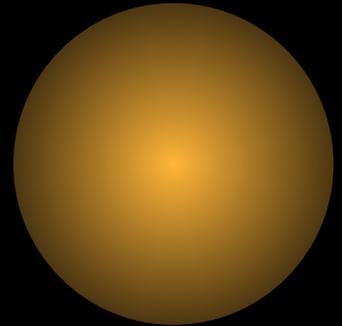
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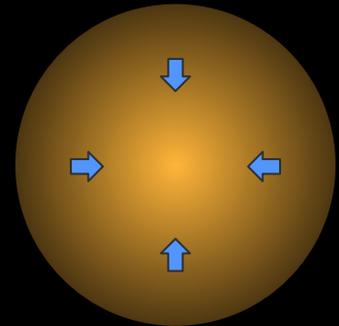
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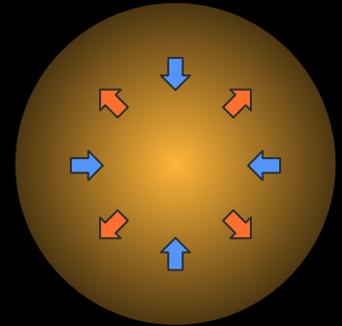
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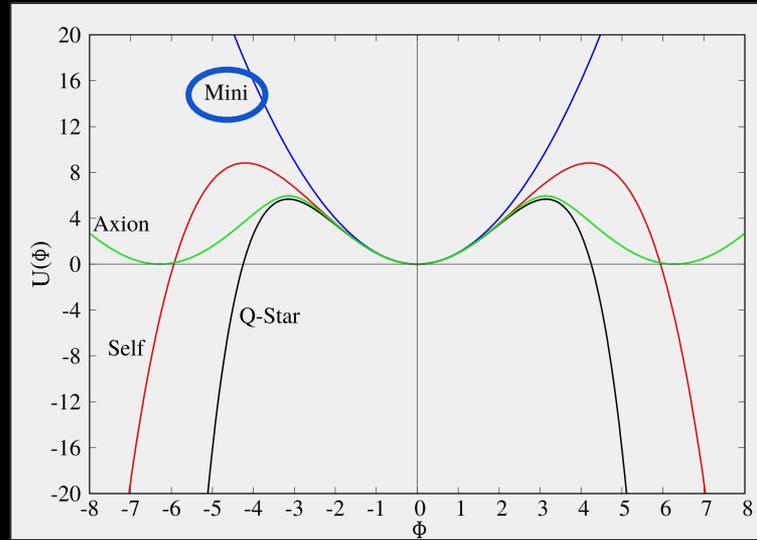
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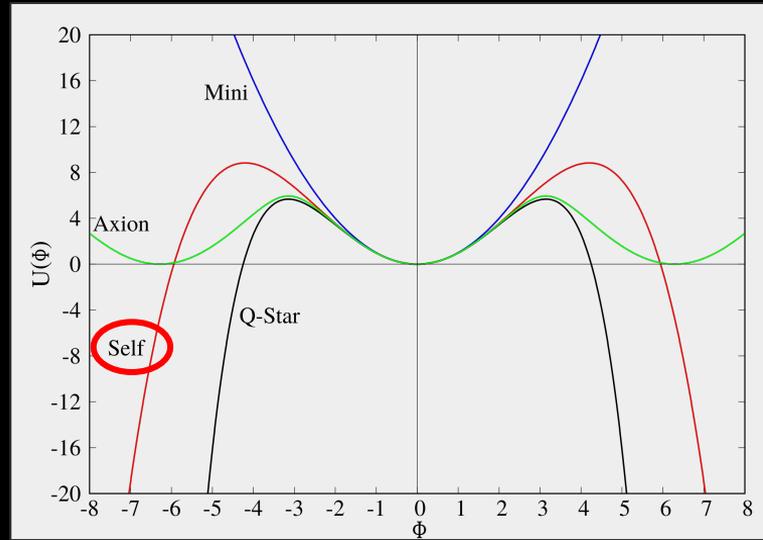
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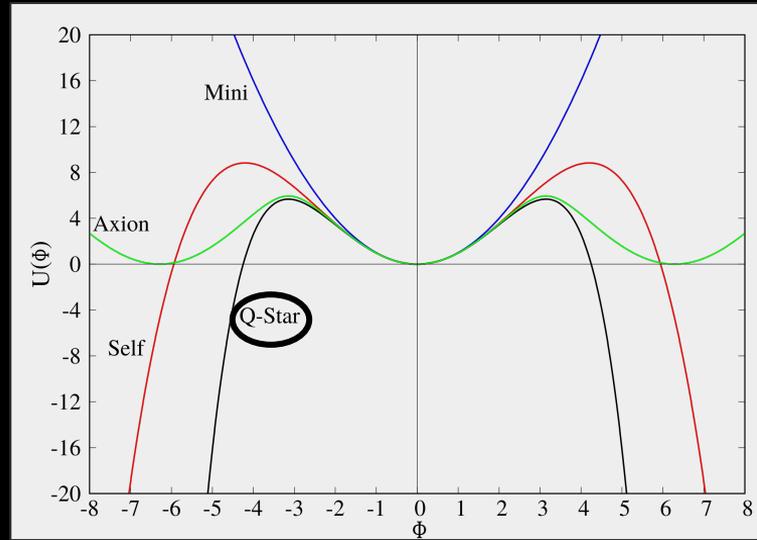
$$U_{\text{poly}} = \mu_S^2 \Phi^2 + \lambda \Phi^4$$



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$$U_{\text{poly}} = \mu_S^2 \Phi^2 + \lambda \Phi^4 + \gamma \Phi^6$$

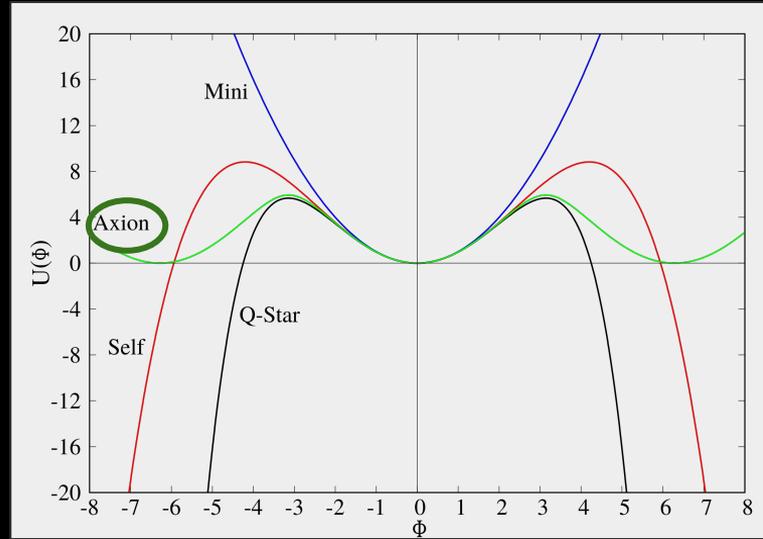


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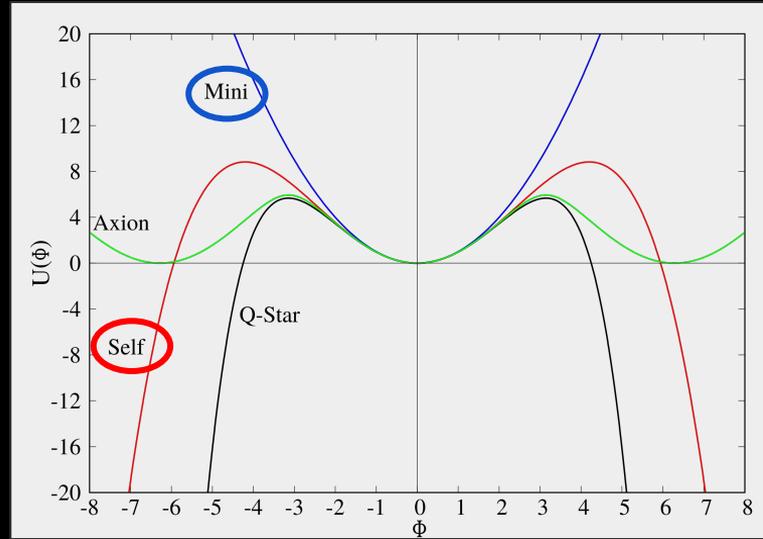
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$$V = \frac{\mu_P^2}{2} \mathbf{A}^2 + \frac{\lambda_P}{4} \mathbf{A}^4$$



Numerics

- We developed a numerical solver to tackle the systems of ODEs
- It is written in C where:
 - The integrator is an explicit 6(5)th Runge-Kutta method
 - The boundary conditions are implemented through a secant/bisection algorithm



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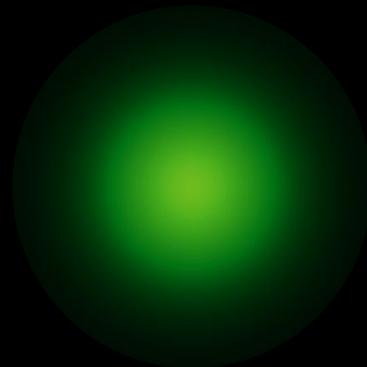
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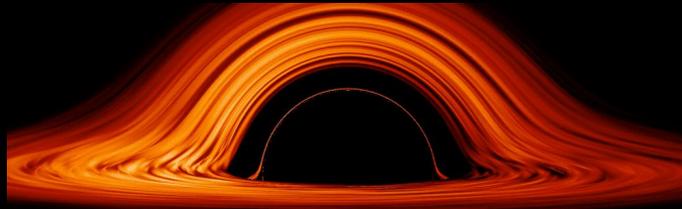
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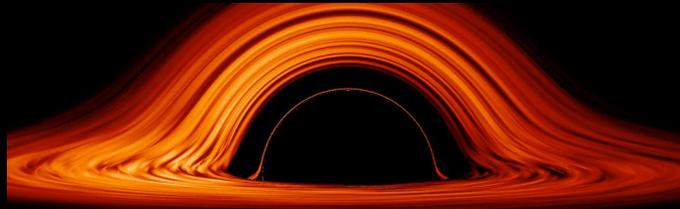
Black Hole Shadow

- If the source of light in the vicinity of the BS has the same morphology as it would have around a BH, the observational image could be similar
- A key feature is the cut-off in the emission due to the disk's inner edge
- Determined by the innermost stable circular orbit (ISCO) of the BH
- However, for spherical BSs there is no ISCO



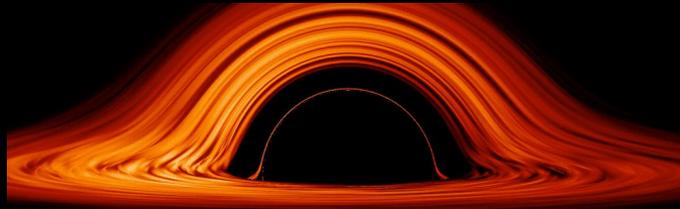
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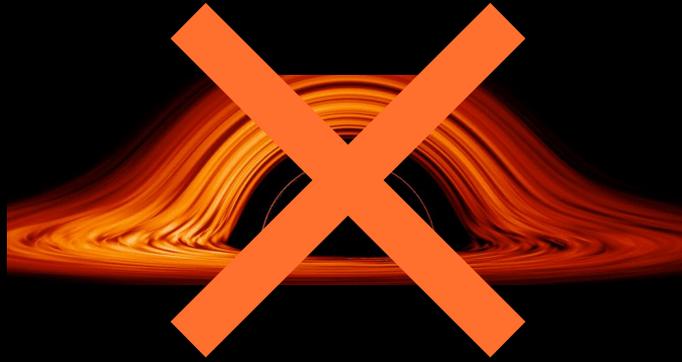
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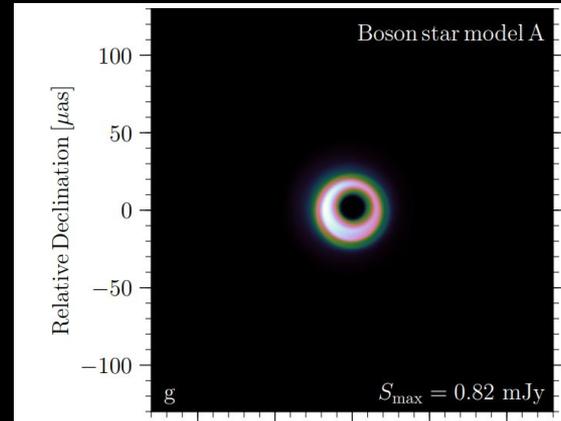
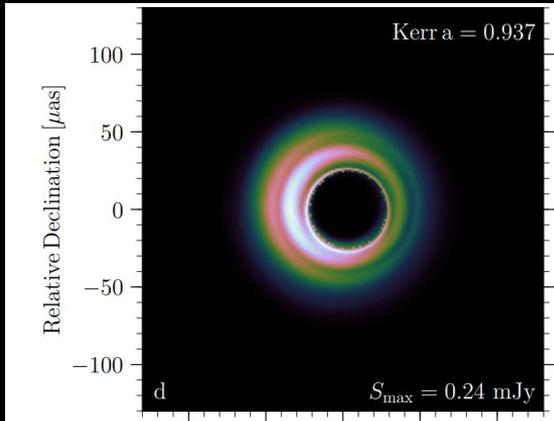
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- Show that the angular velocity of the orbits, Ω , attains a maximum at some areal radius R_Ω
- Suggesting that R_Ω determines the inner edge of the accretion disk

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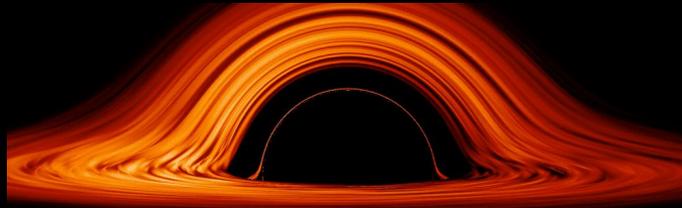
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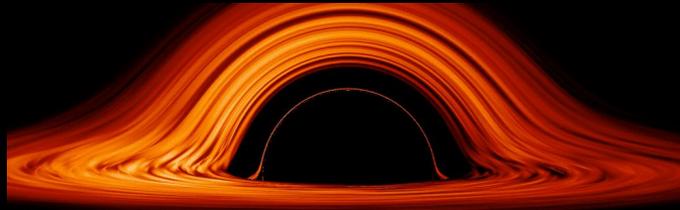
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- The objective of this work is to assess if **stable and dynamically robust** BS can yield the same shadow as a BH.

Light Rings
and
Timelike Circular Orbits

Geodesic motion

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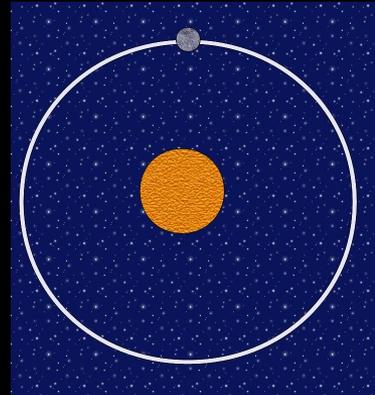
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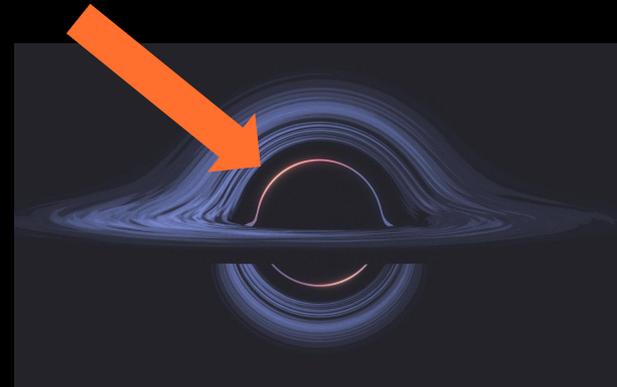
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- For a circular orbit, $\dot{r} = 0 = \ddot{r}$



Geodesic motion: Light Rings $k = 0$

- Let us first consider null geodesics ($k = 0$)

$$-r\sigma \left(\frac{-2m'}{r} + \frac{2m}{r^2} \right) + 2 \left(1 - \frac{2m}{r} \right) (\sigma - r\sigma') = 0.$$



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- We wish to find the first BS solution containing a LR
- In other words, the first ultracompact BS

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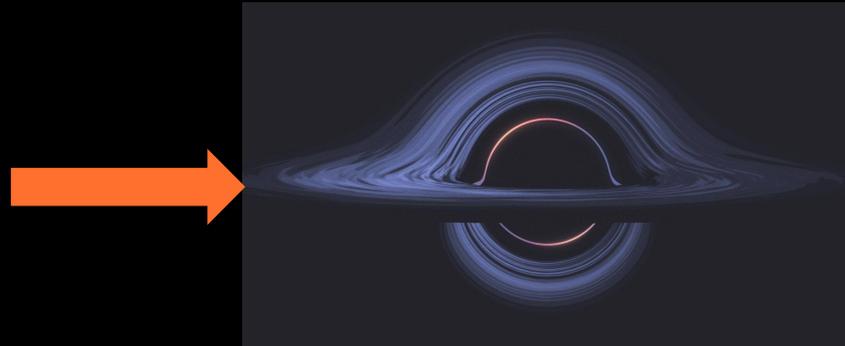
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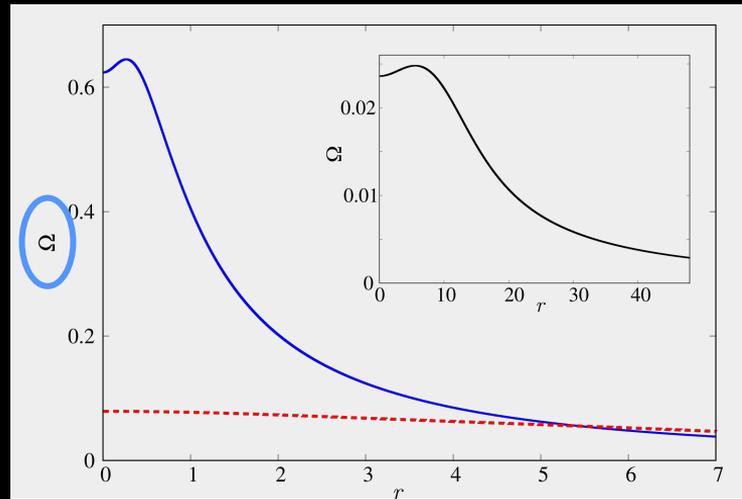
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- However an accretion disk may have an inner edge even around BSs without an ISCO Olivares et al. , MN of the RAS 2020
- This occurs if the angular velocity along TCOs attains a maximum at some radial distance. The corresponding areal radius is denoted R_{Ω}

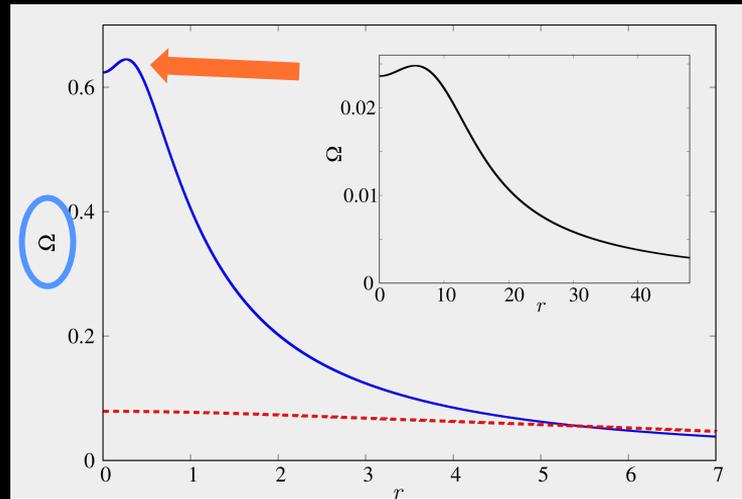
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Geodesic motion: Keeping up

- Does it provide a similar scale, for a BS and a Schw. BH?

$$\xi \equiv \frac{R_{ISCO}}{R_\Omega}$$

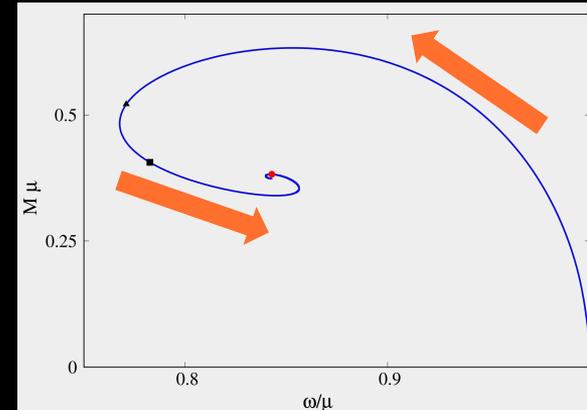
- Moving along the spiral, the ADM mass and frequency undergo oscillations
- The field amplitude at the origin, on the other hand, grows monotonically
- To uniquely label the solutions, let us introduce

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- The field amplitude at the origin, on the other hand, grows monotonically
- To uniquely label the solutions, let us introduce

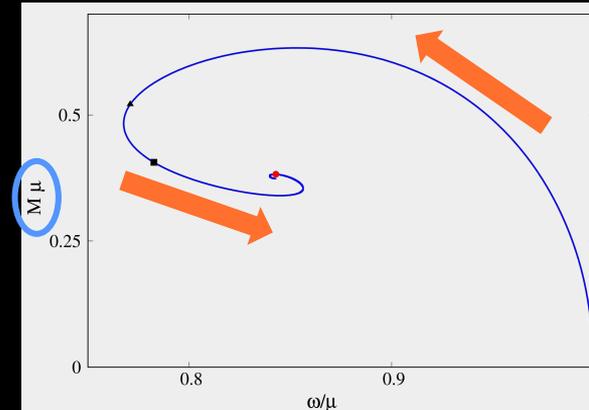


Geodesic motion: Keeping up

- Does it provide a similar scale, for a BS and a Schw. BH?

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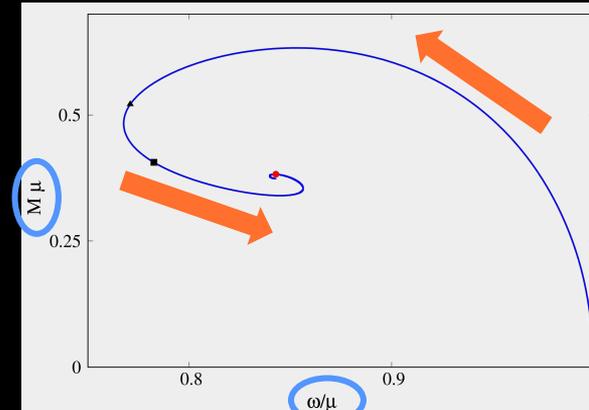


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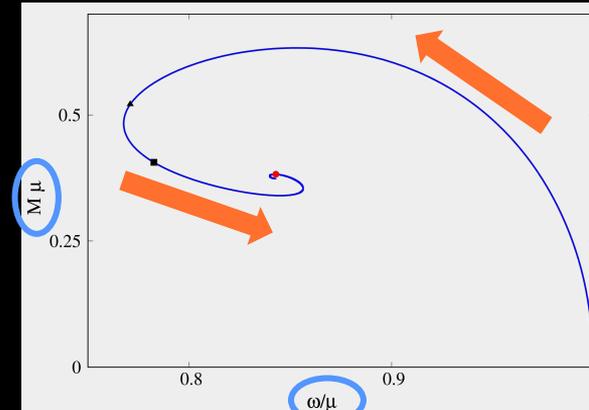


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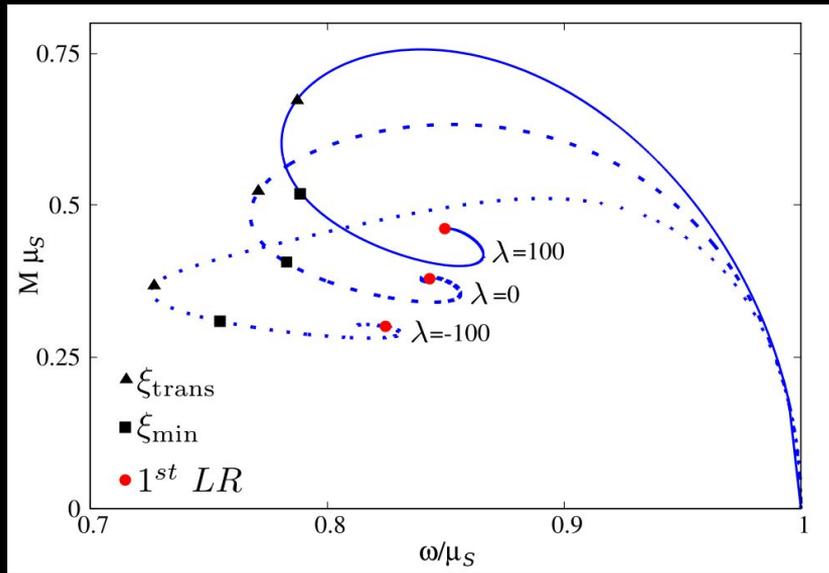
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$$\chi(\boldsymbol{x}) \equiv \frac{\varphi_0(\boldsymbol{x})}{\varphi_0(M_{\max})} \quad [\text{scalar}] \quad \text{or} \quad \chi(\boldsymbol{x}) \equiv \frac{f_0(\boldsymbol{x})}{f_0(M_{\max})} \quad [\text{vector}] ,$$

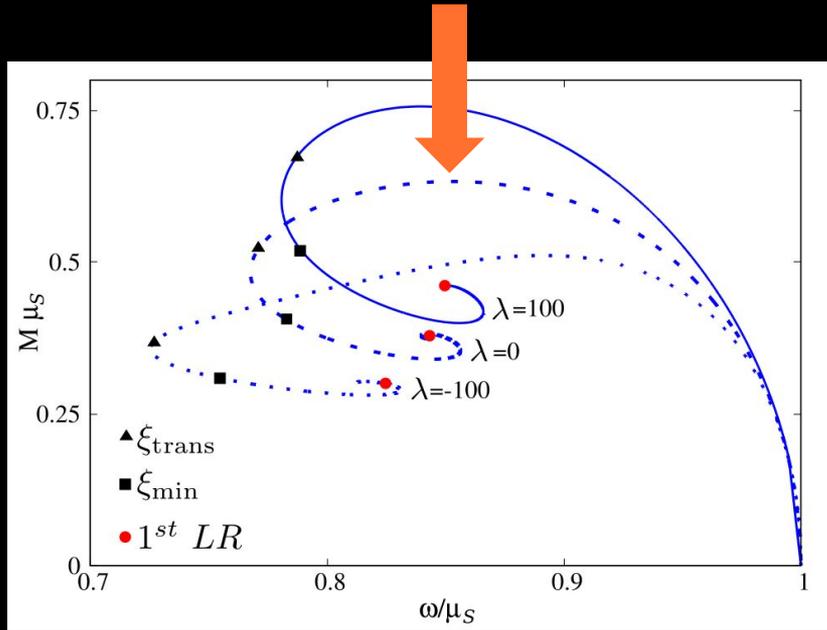
Numerical Results

Boson Stars: Polynomial Scalar



$$U_{\text{poly}} = \mu_S^2 \Phi^2$$

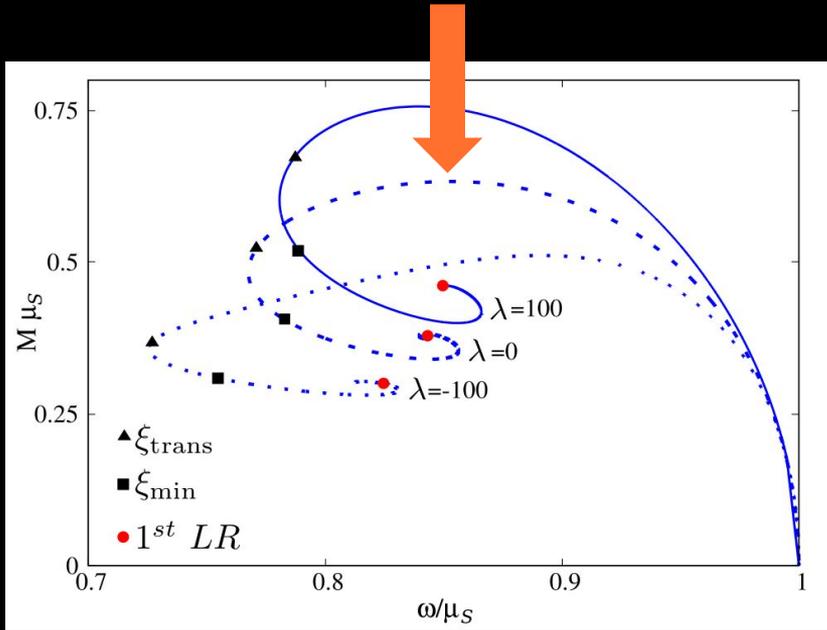
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Solutions are stable up to the first maximum of the mass

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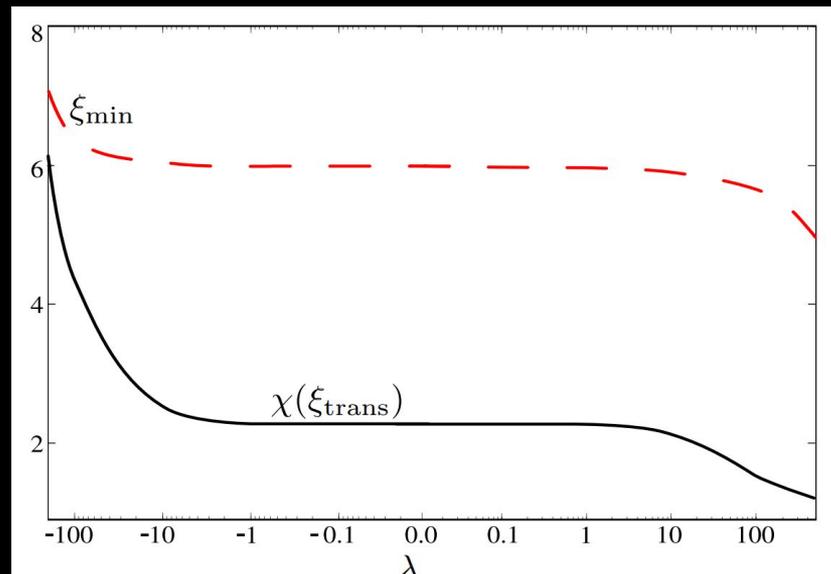
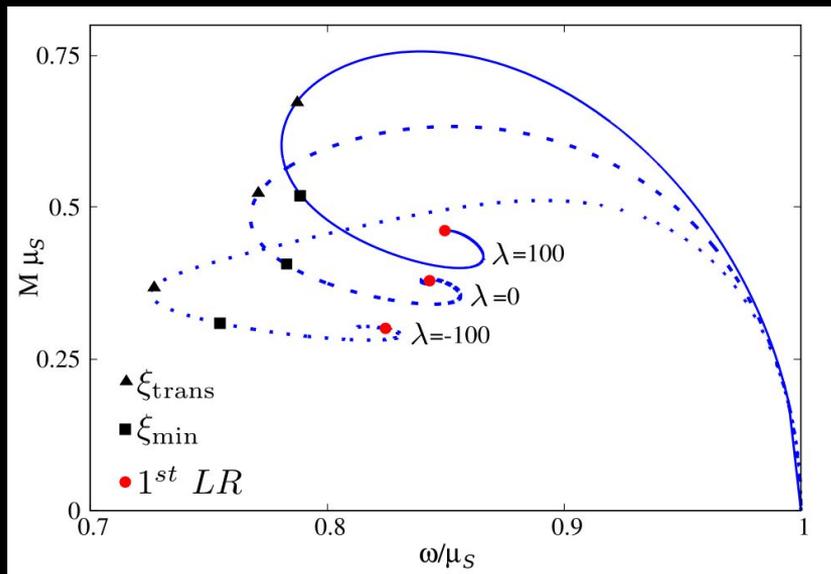


Solutions are stable up to the first maximum of the mass

There is no stable ultracompact Boson Star solution

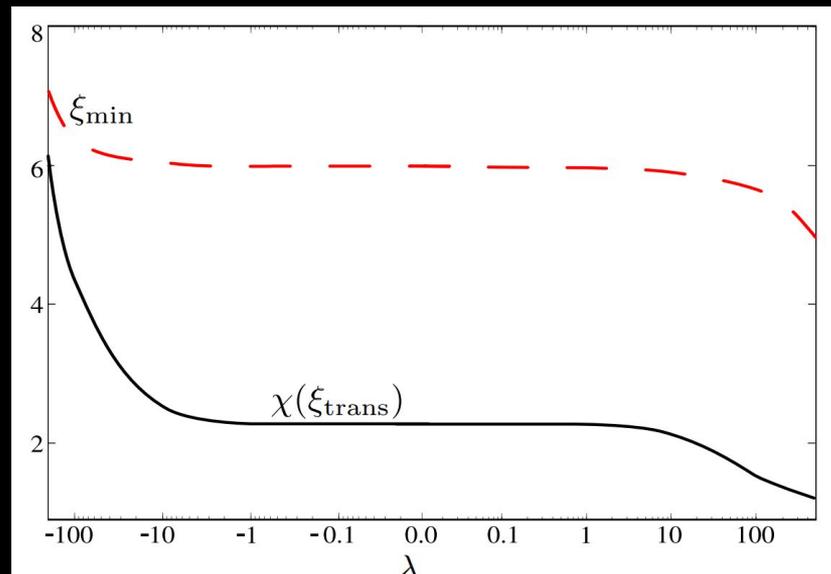
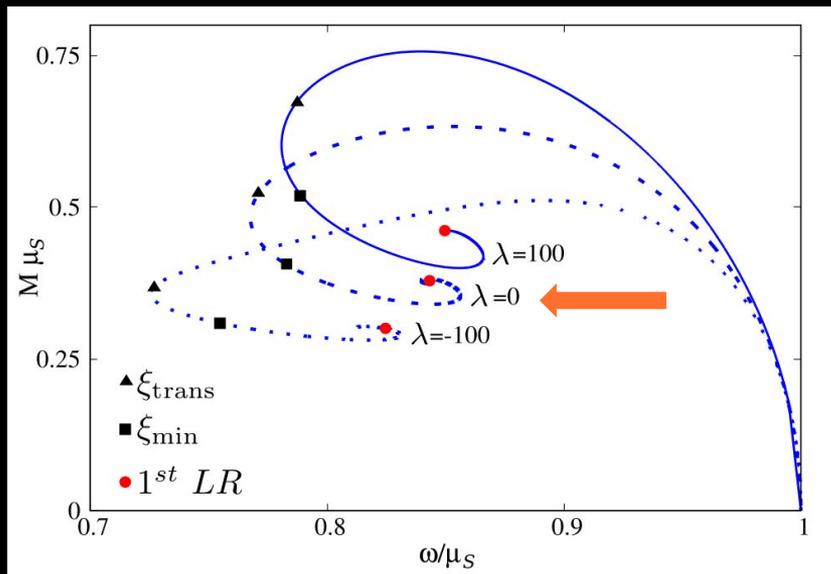
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Boson Stars: Polynomial Scalar



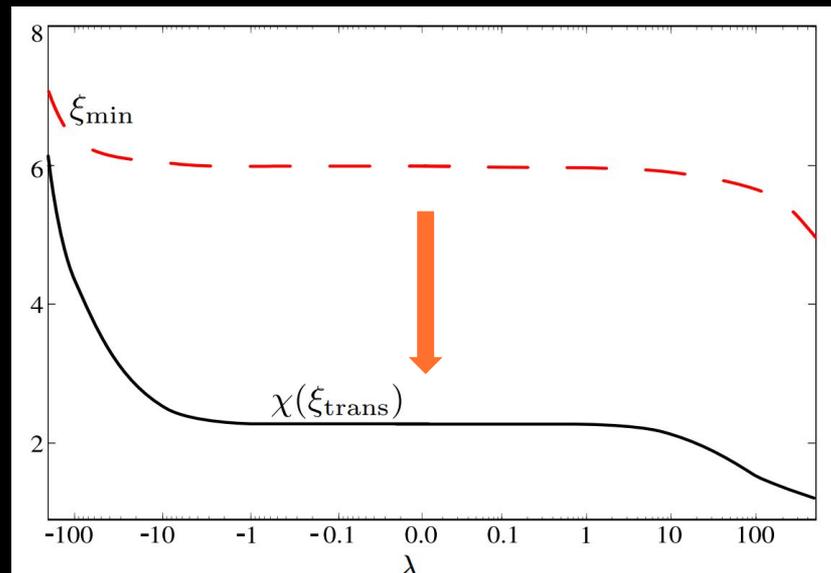
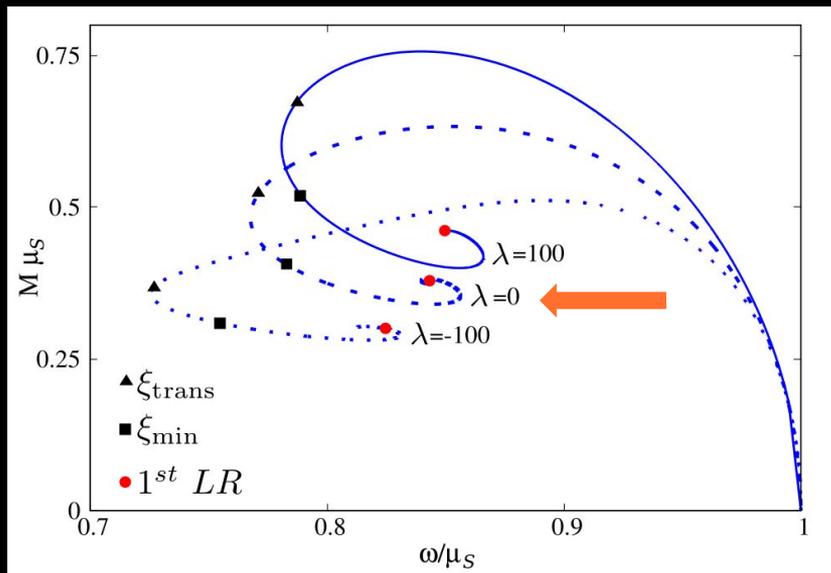
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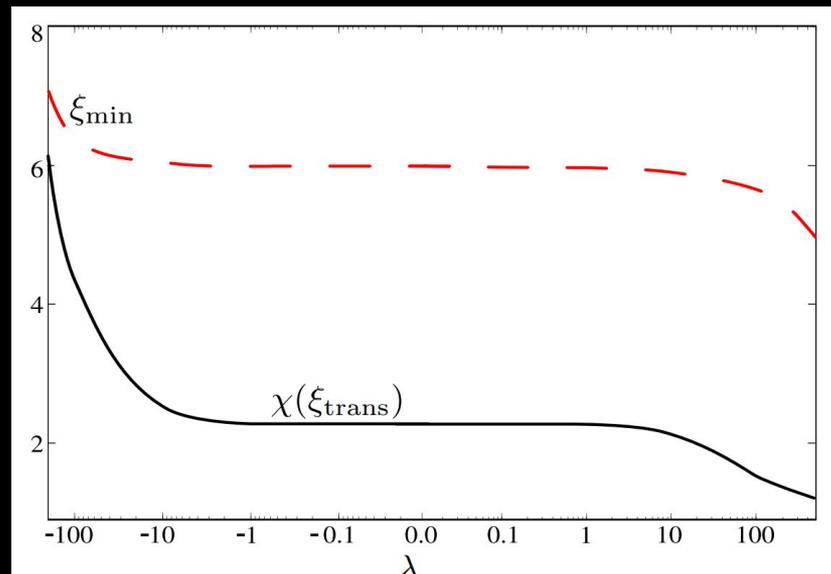
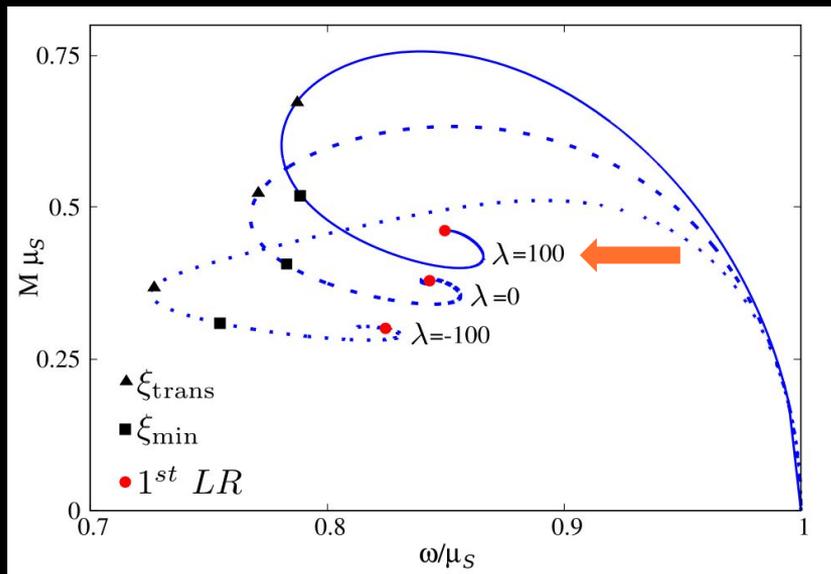
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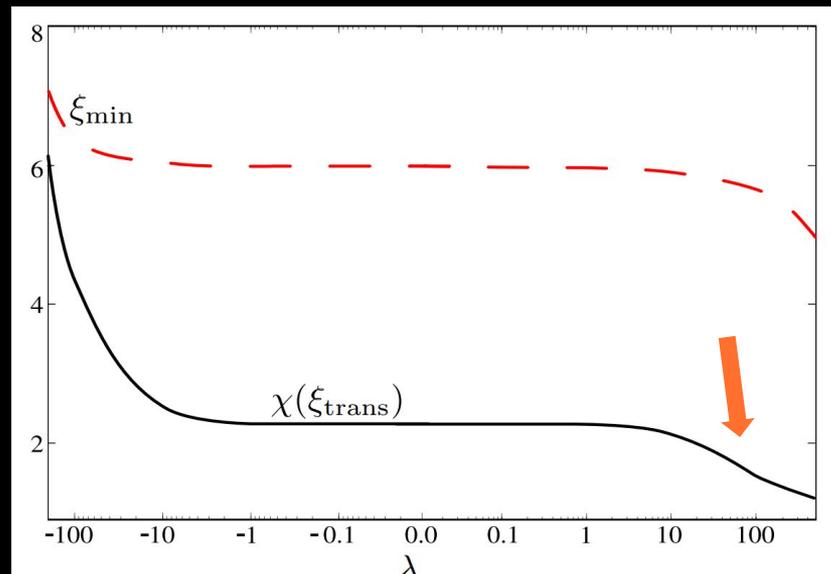
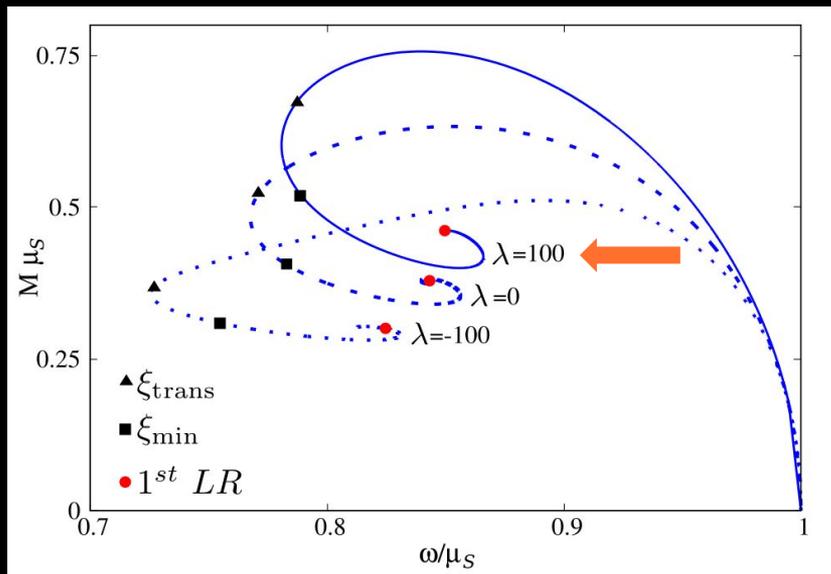
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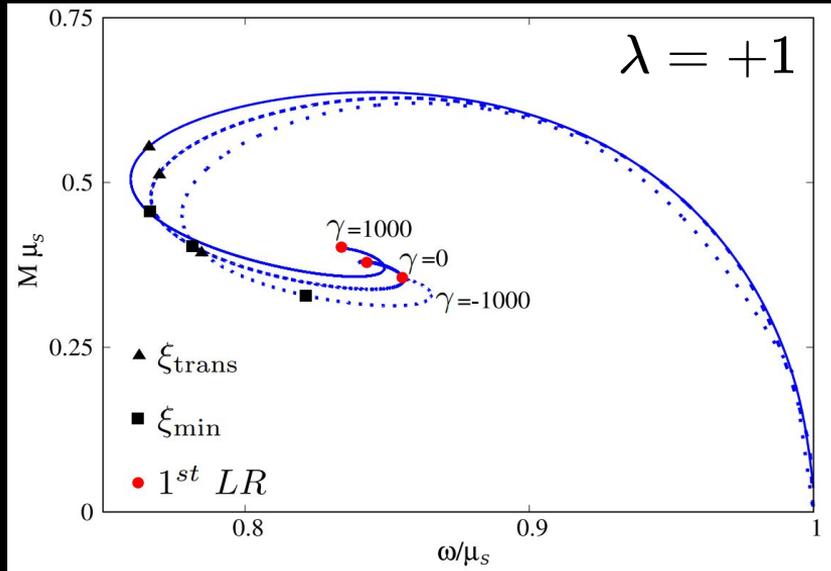
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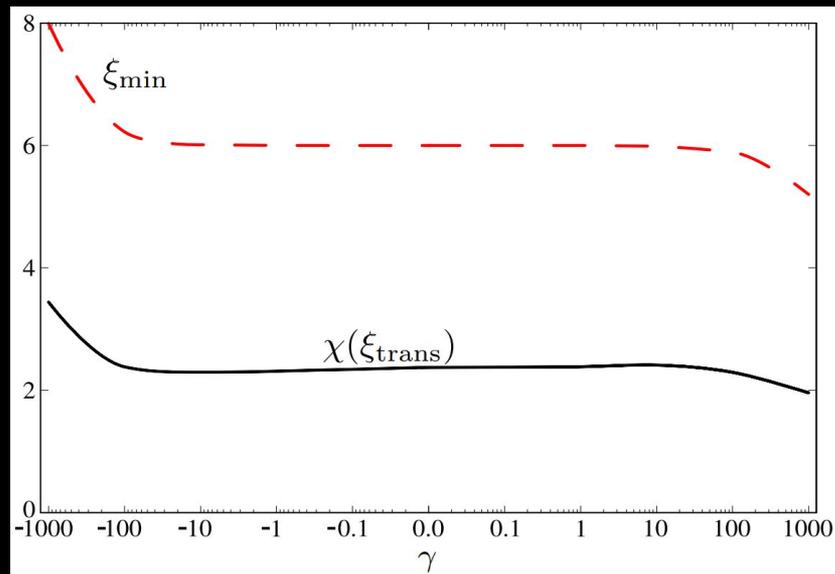
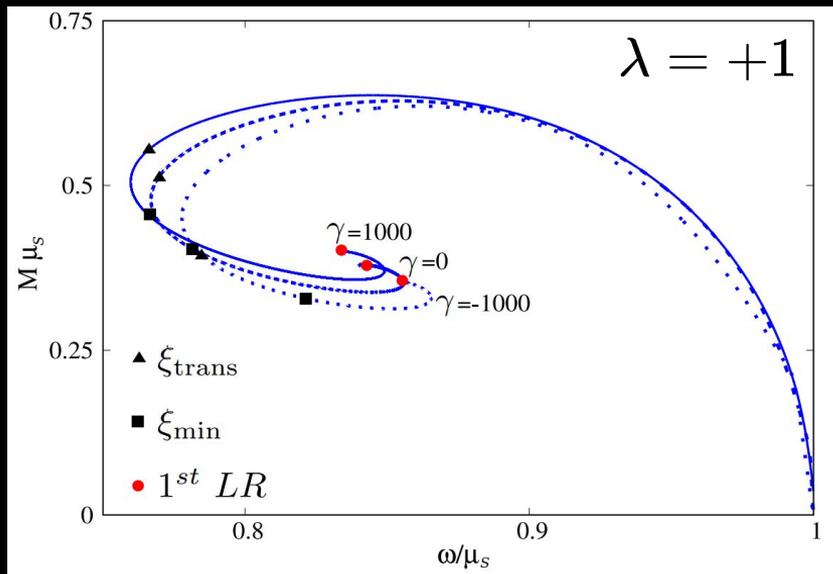
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Boson Stars: Polynomial Scalar



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- $\chi(\xi_{trans}) = 1.51$
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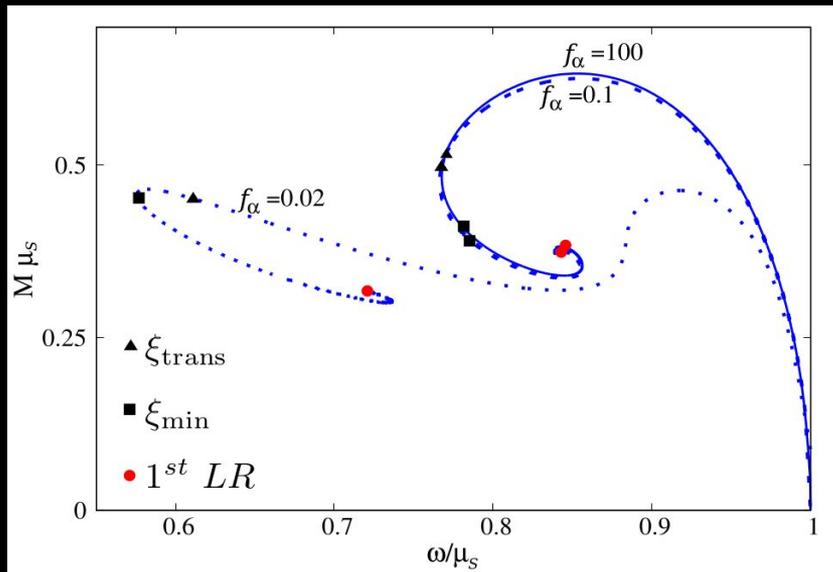
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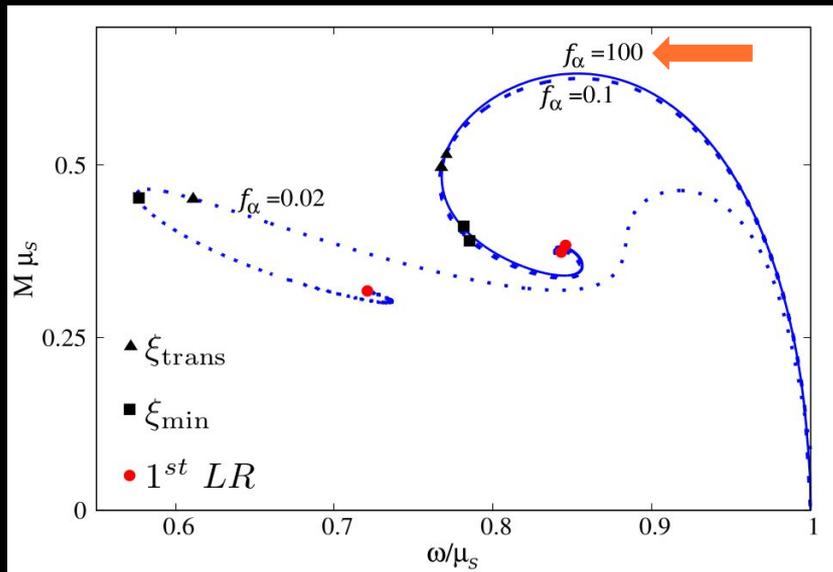
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Boson Stars: Axion Scalar



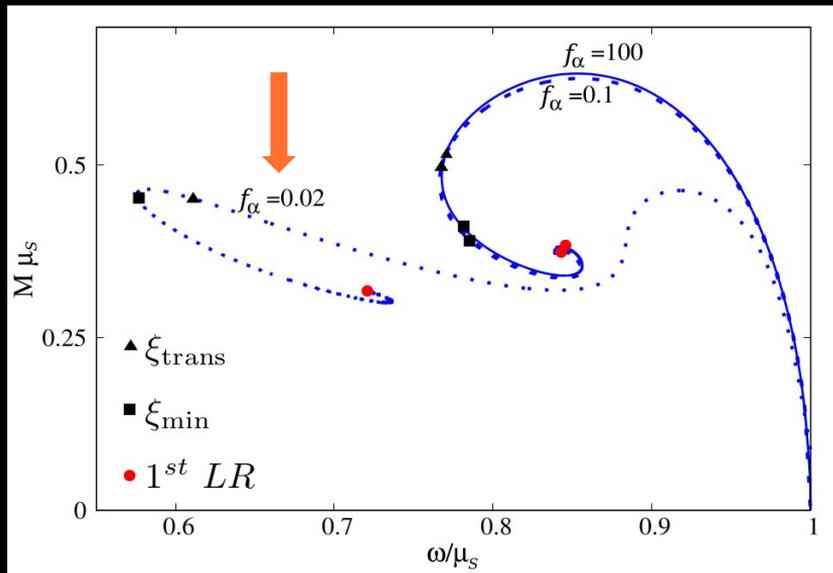
$$U_{\text{axion}} = \frac{2\mu_s^2 f_\alpha^2}{\hbar B} \left[1 - \sqrt{1 - 4B \sin^2 \left(\frac{\Phi \sqrt{\hbar}}{2f_\alpha} \right)} \right]$$

Boson Stars: Axion Scalar



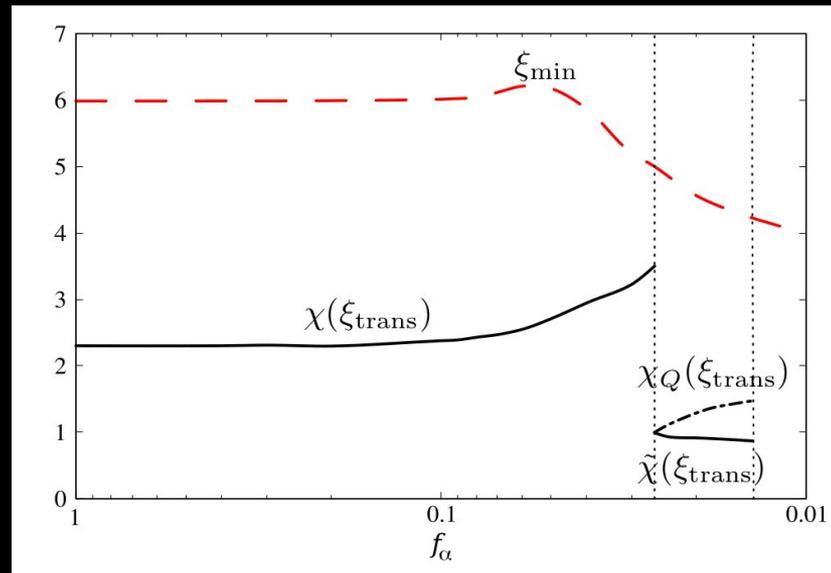
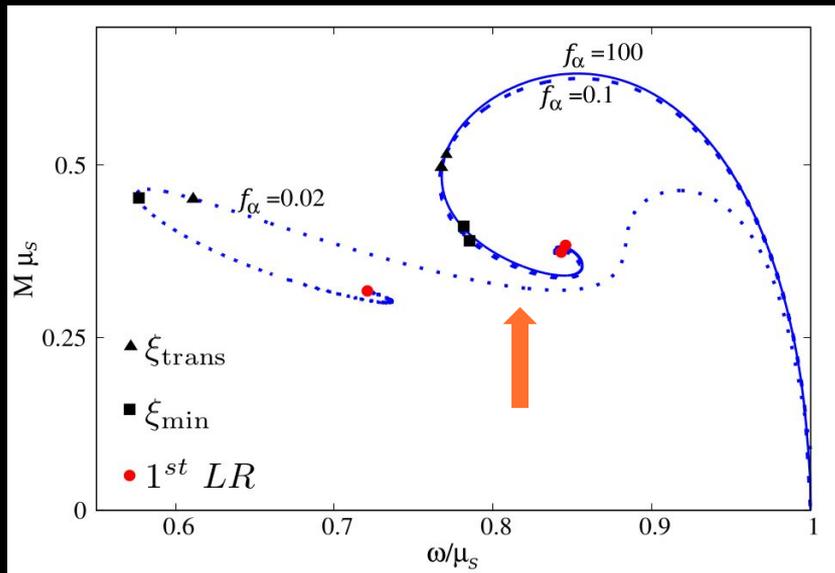
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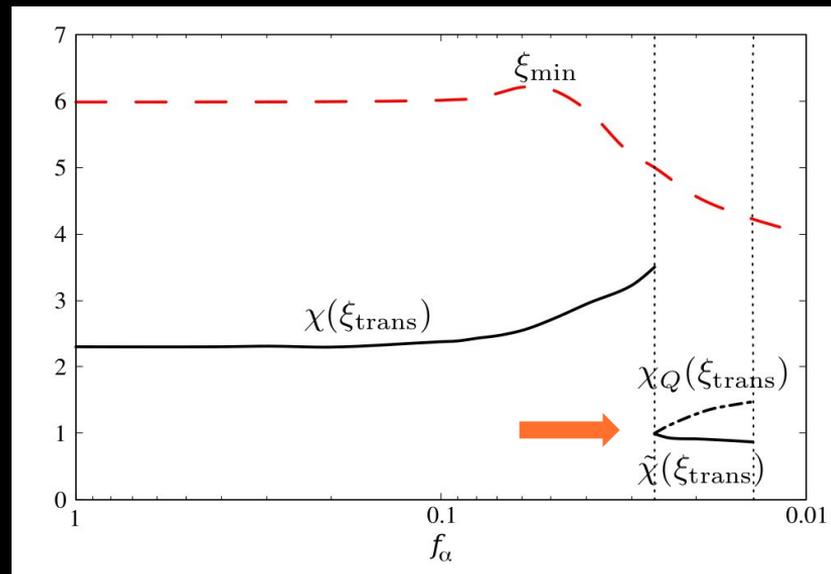
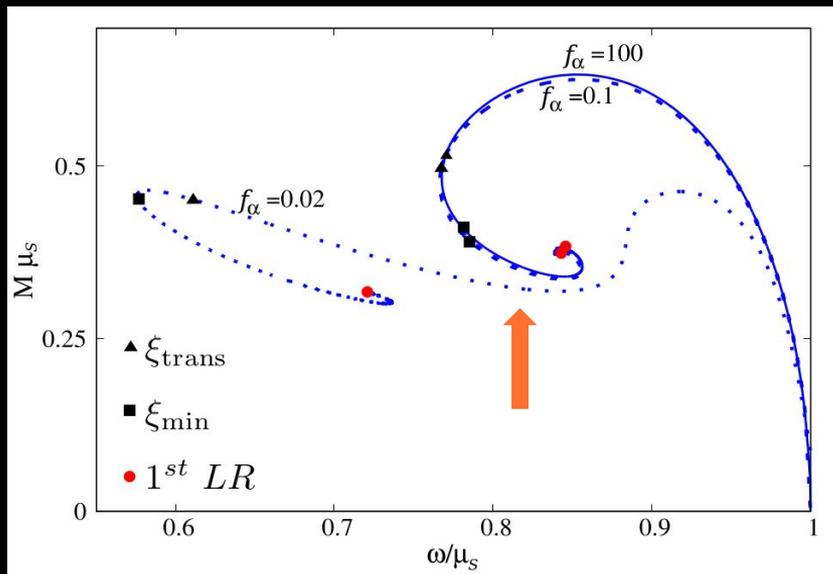
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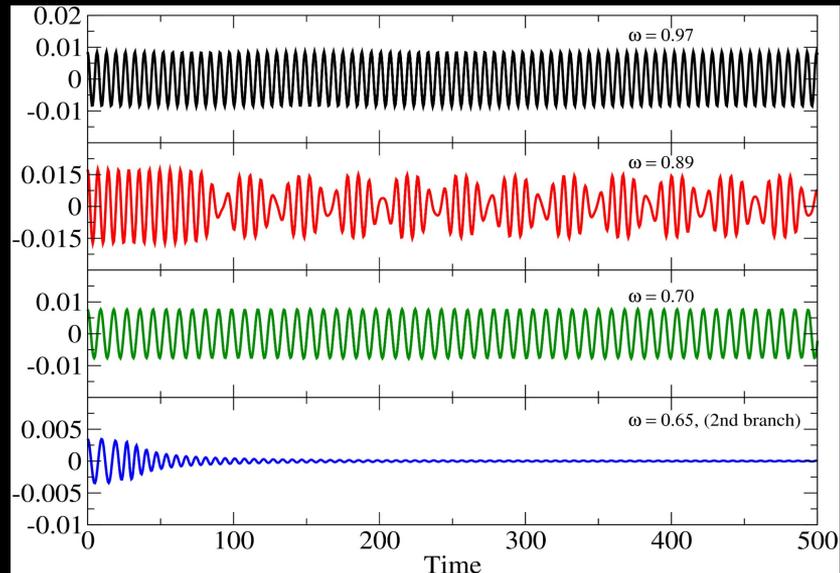
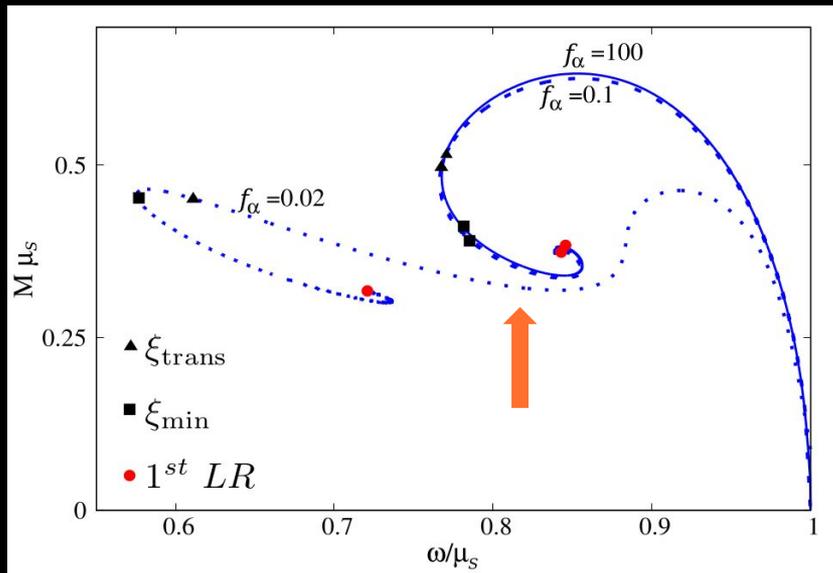
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Boson Stars: Stability



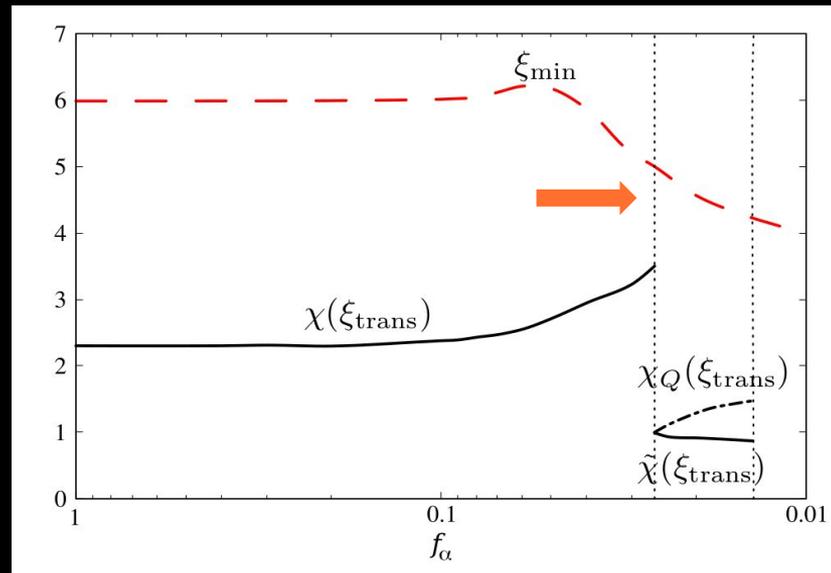
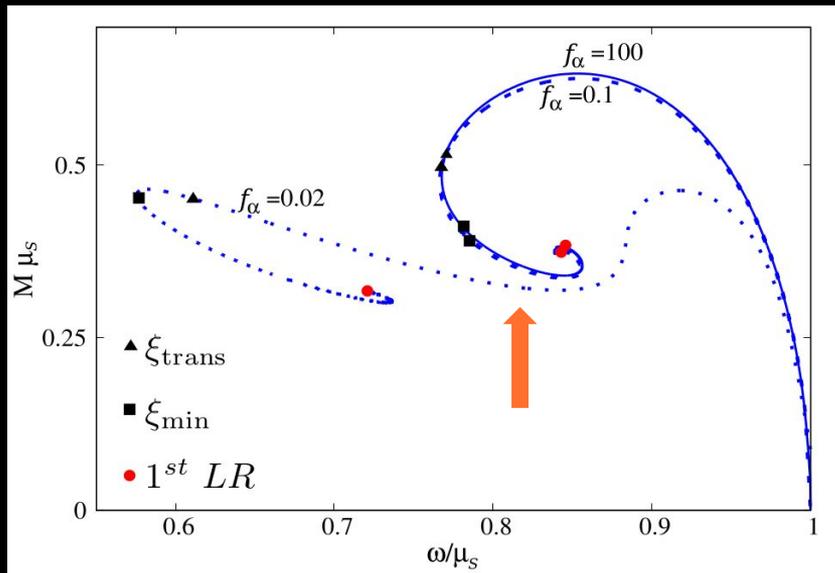
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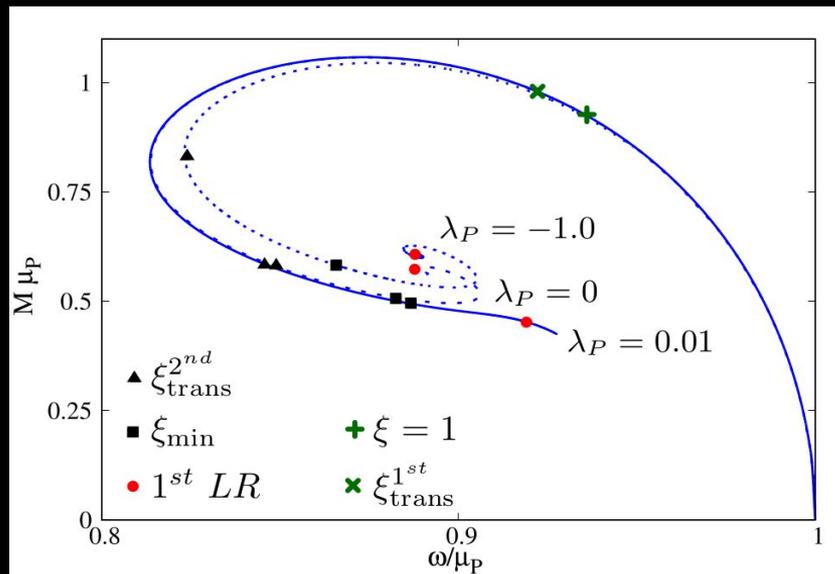
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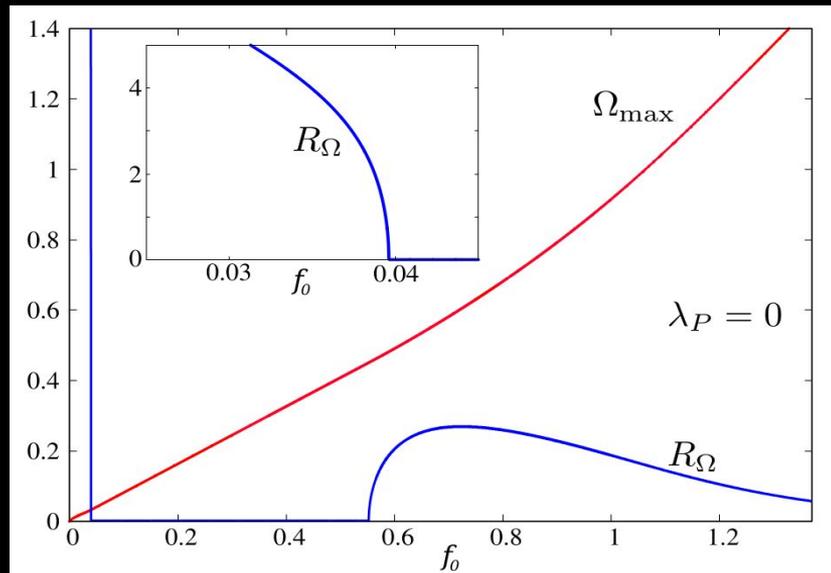
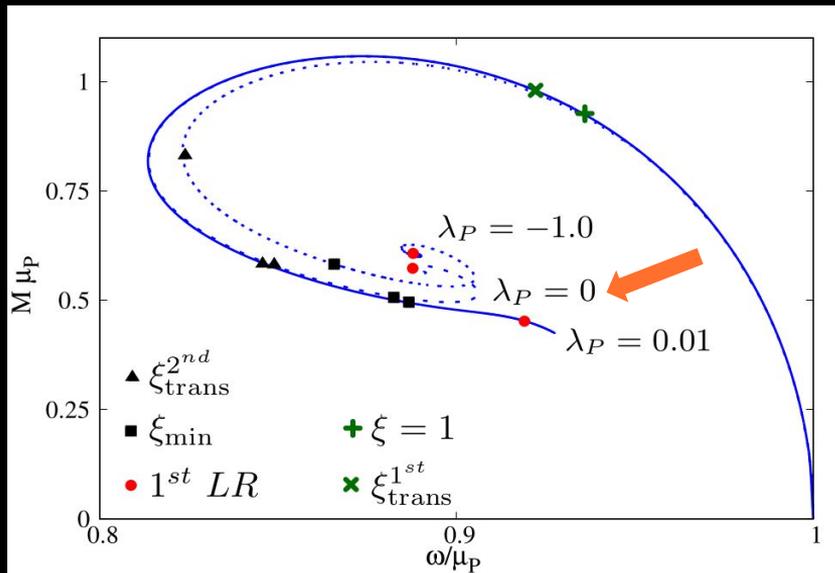
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Boson Stars: Proca



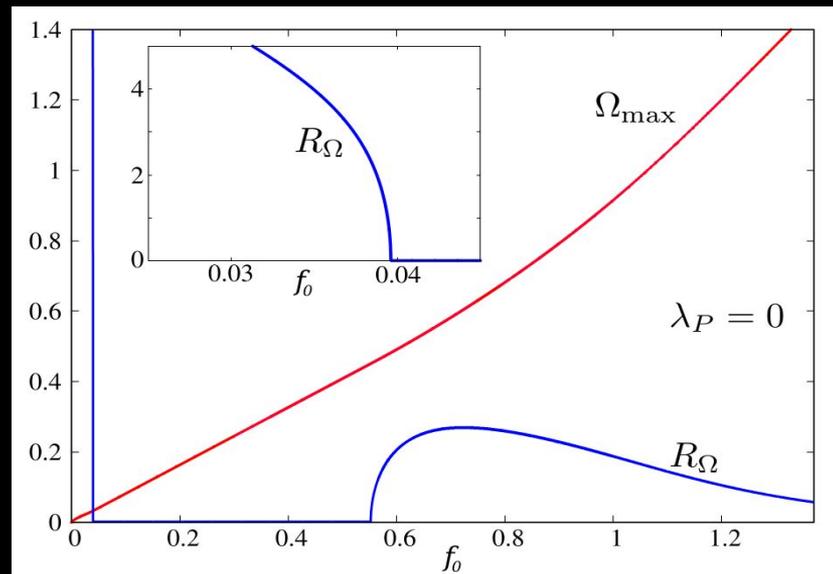
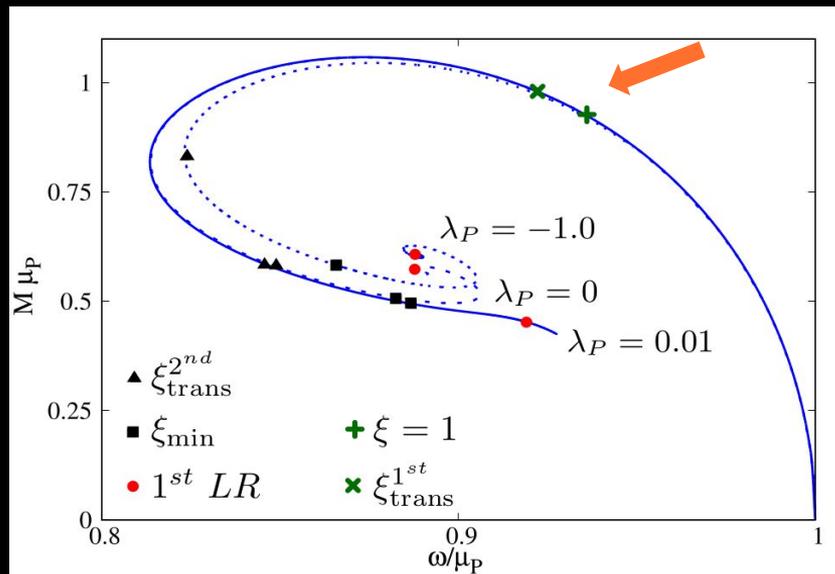
$$V = \frac{\mu_P^2}{2} \mathbf{A}^2$$

Boson Stars: Proca



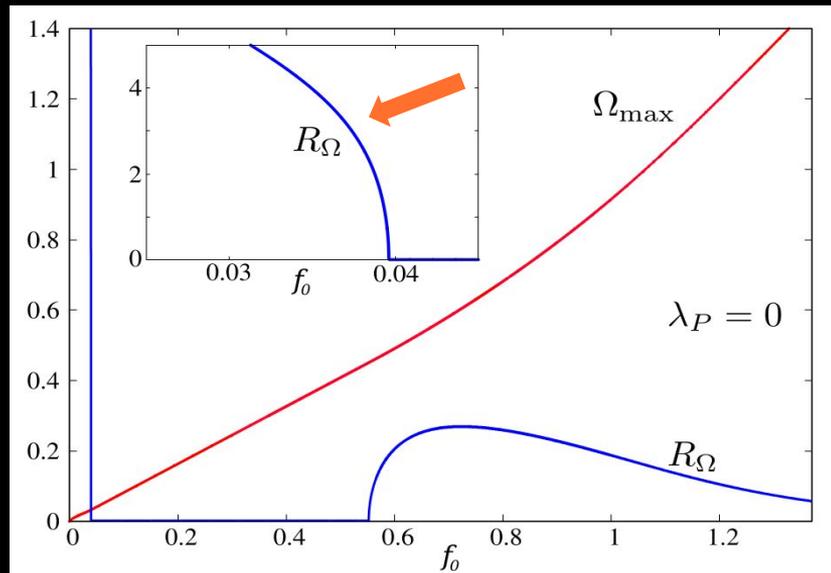
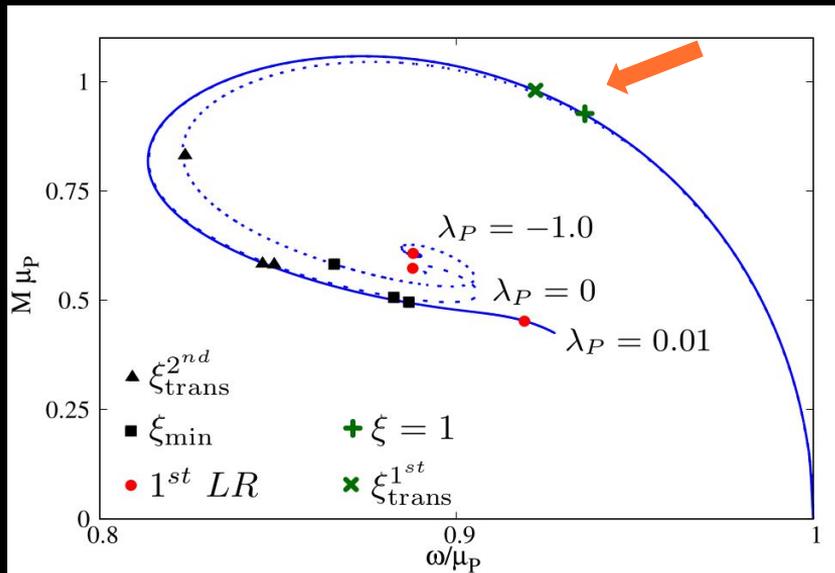
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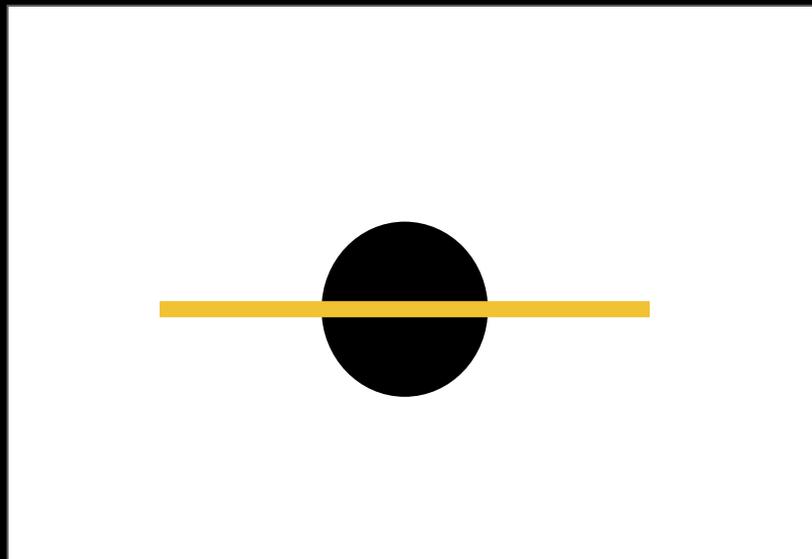
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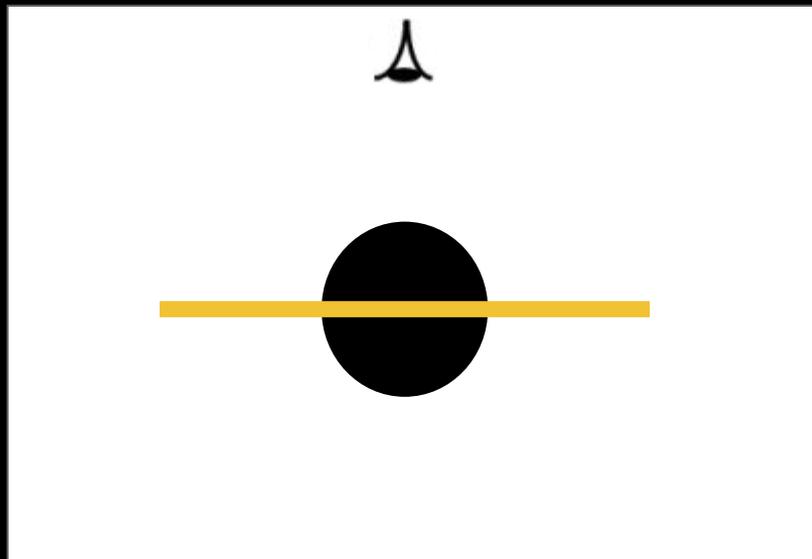
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Shadow

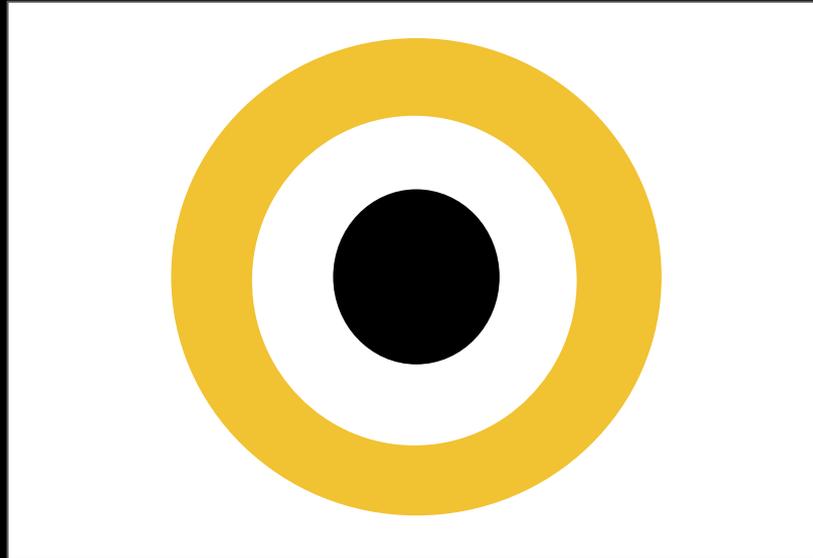
Shadow: BH vs PS



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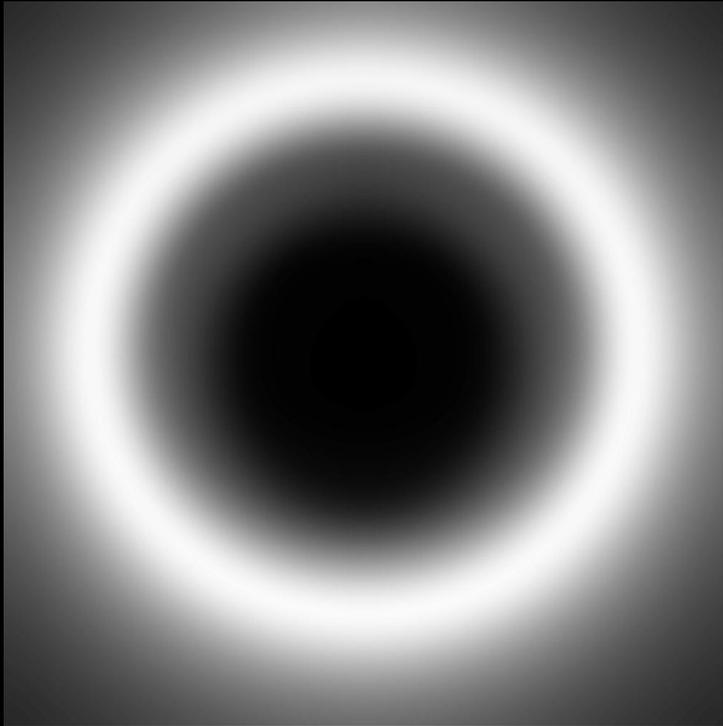


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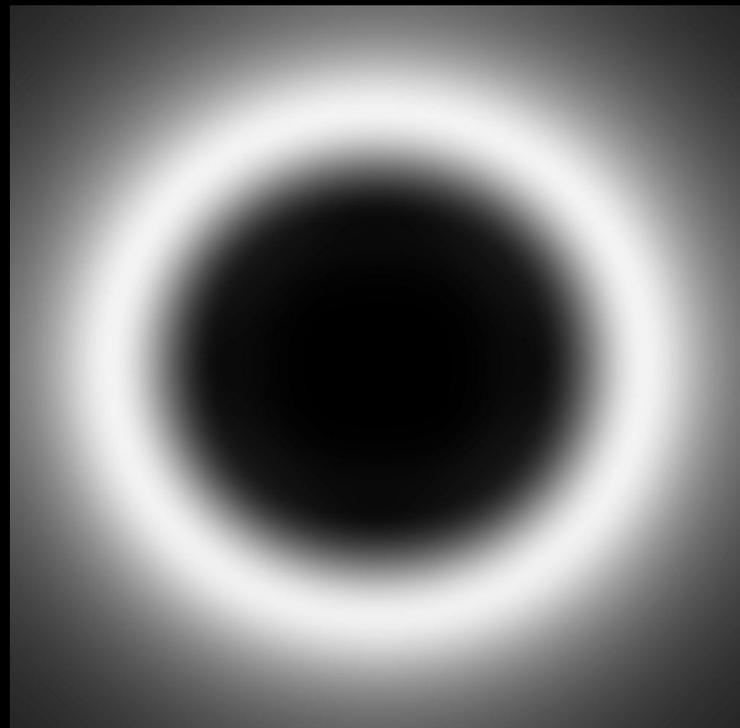


Shadow: BH vs PS

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PS

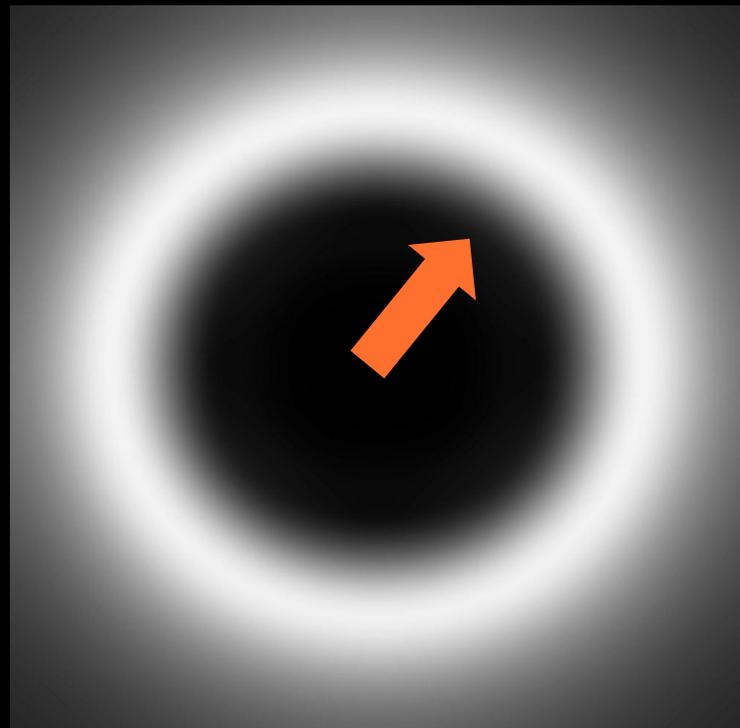


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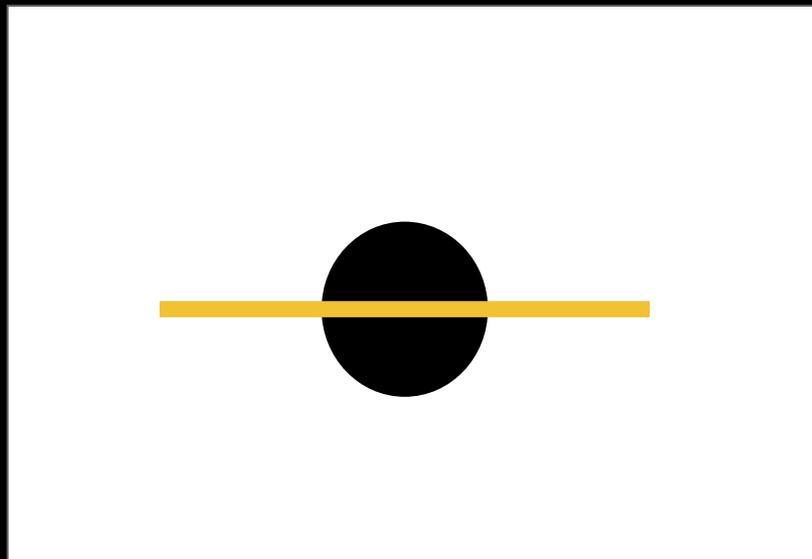
BH



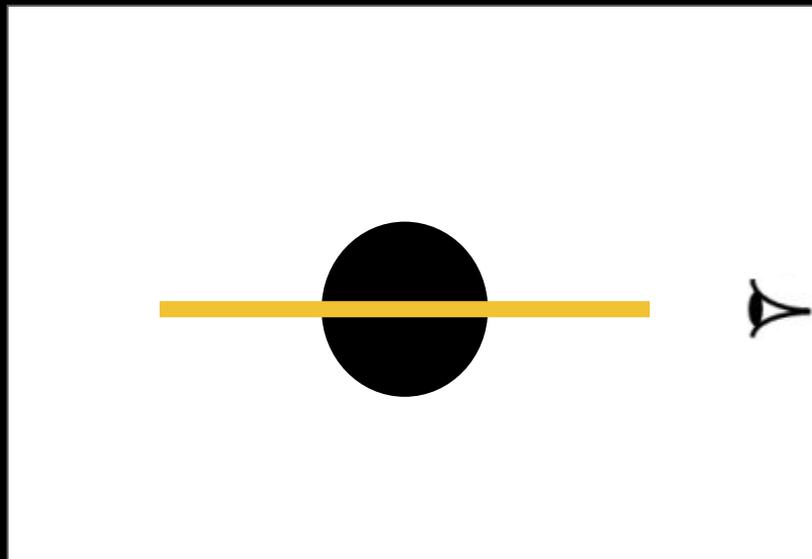
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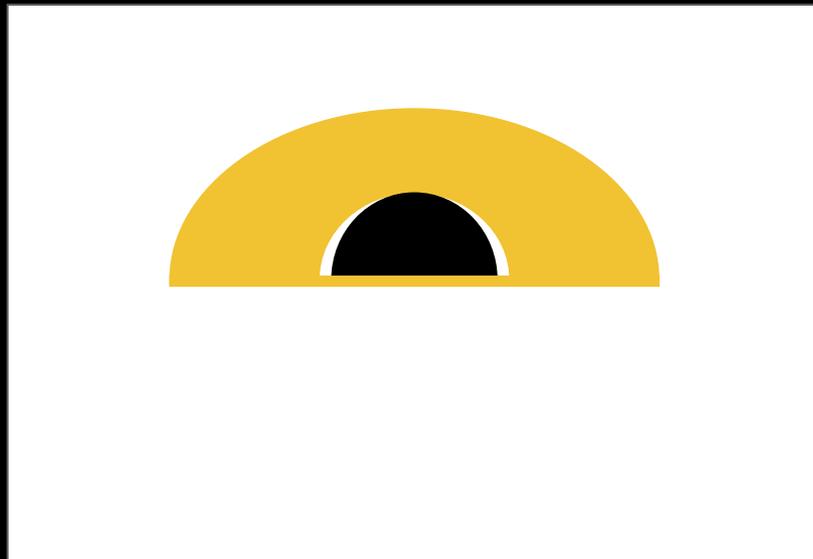
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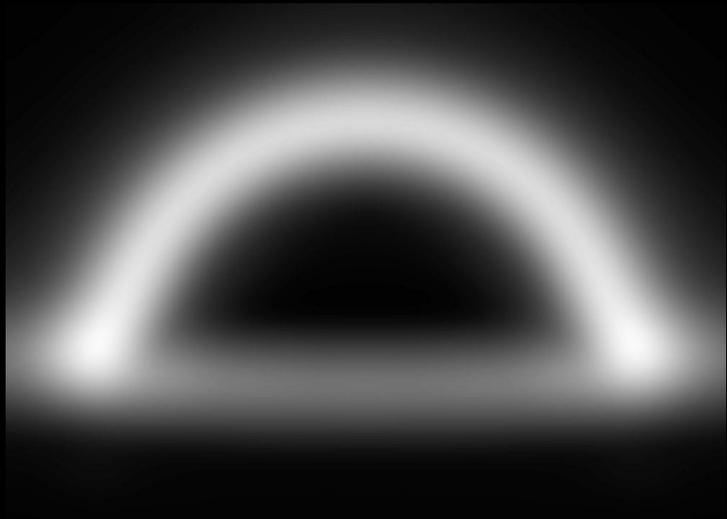


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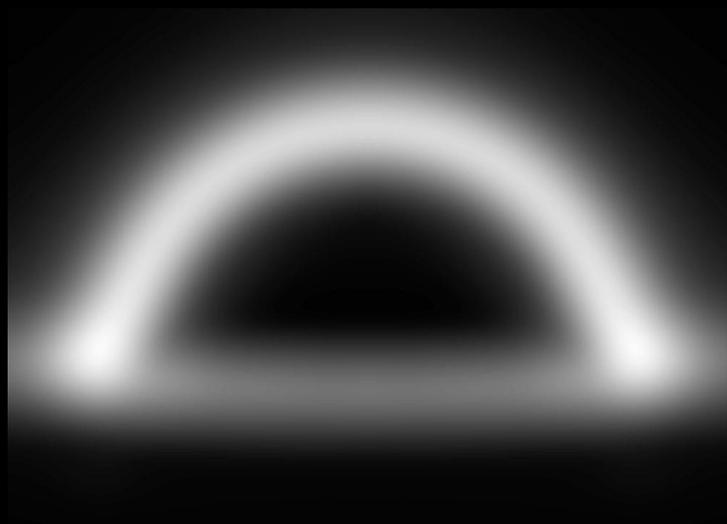
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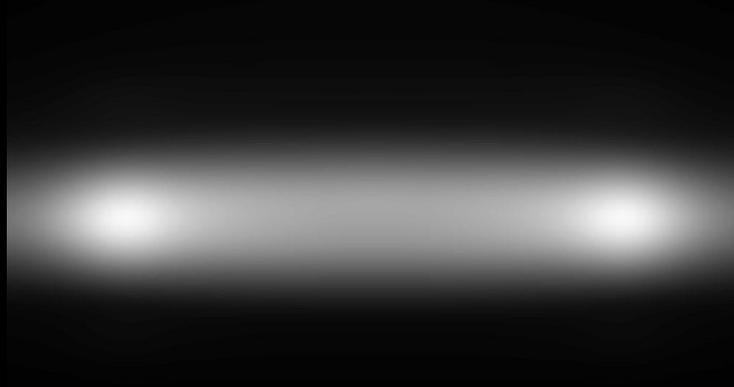


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BH



PS

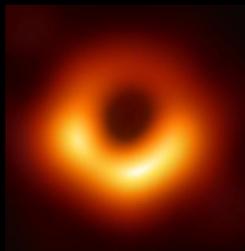


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BH



PS



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- Models with dynamically robust spherical BSs, can mimic the shadow of a Schwarzschild BH
- In the case of spherically stable scalar BSs:
- While polynomial self-interaction cannot easily solve this issue;
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- On the other hand, for spherical PSs
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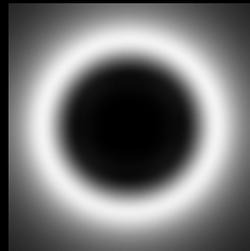
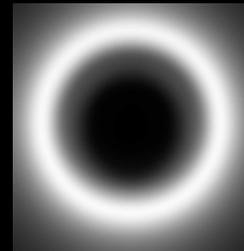
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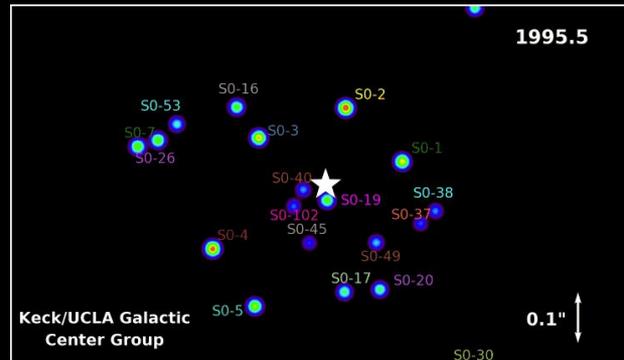
Kinematics

Black Hole mimicker

- In the absence of an illumination, direct observations are impossible
- However, the gravitational potential still exist allowing the orbiting luminous objects
- The dark object nature can be inferred by the orbit of the luminous
- A distinguishable characteristic of a BS is a spatial matter distribution

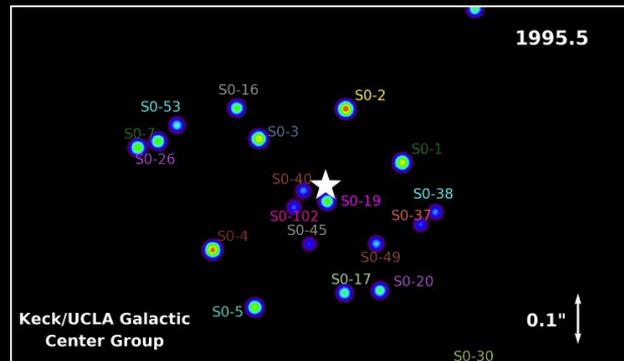
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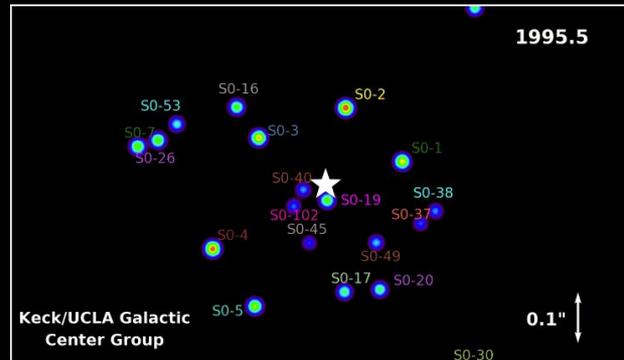
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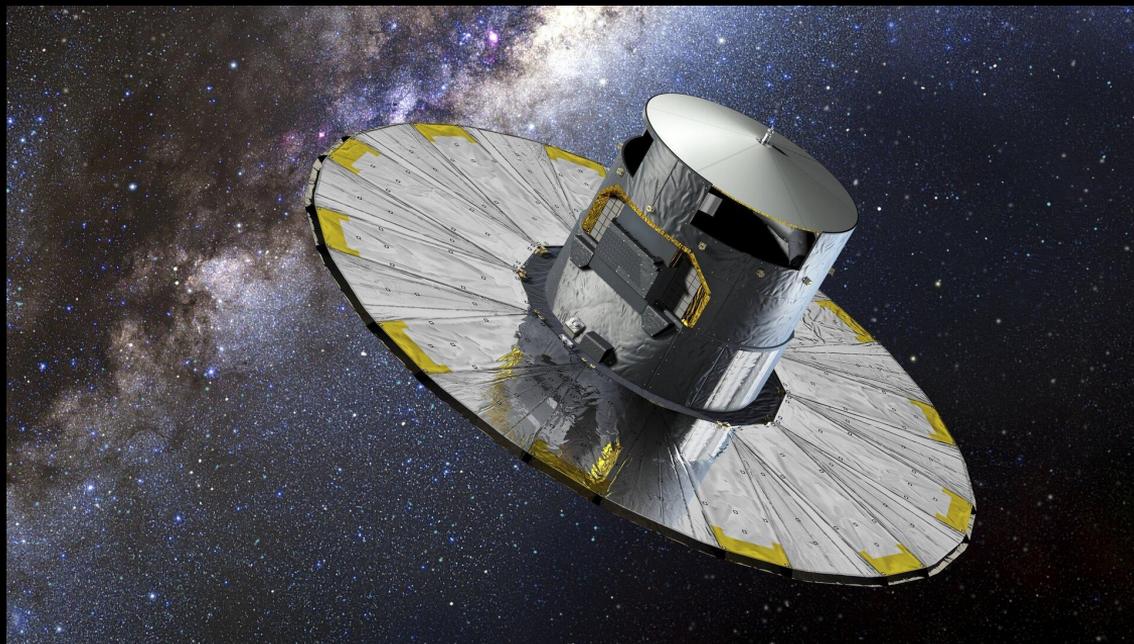
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Black Hole mimicker

- In the absence of an illumination, direct observations are impossible
 - However, the gravitational potential still exist allowing the orbiting luminous objects
 - The dark object nature can be inferred by the orbit of the luminous
 - A distinguishable characteristic of a BS is a spatial matter distribution
-
- Can a luminous star probe the nature of the central dark object?

Observation: GAIA



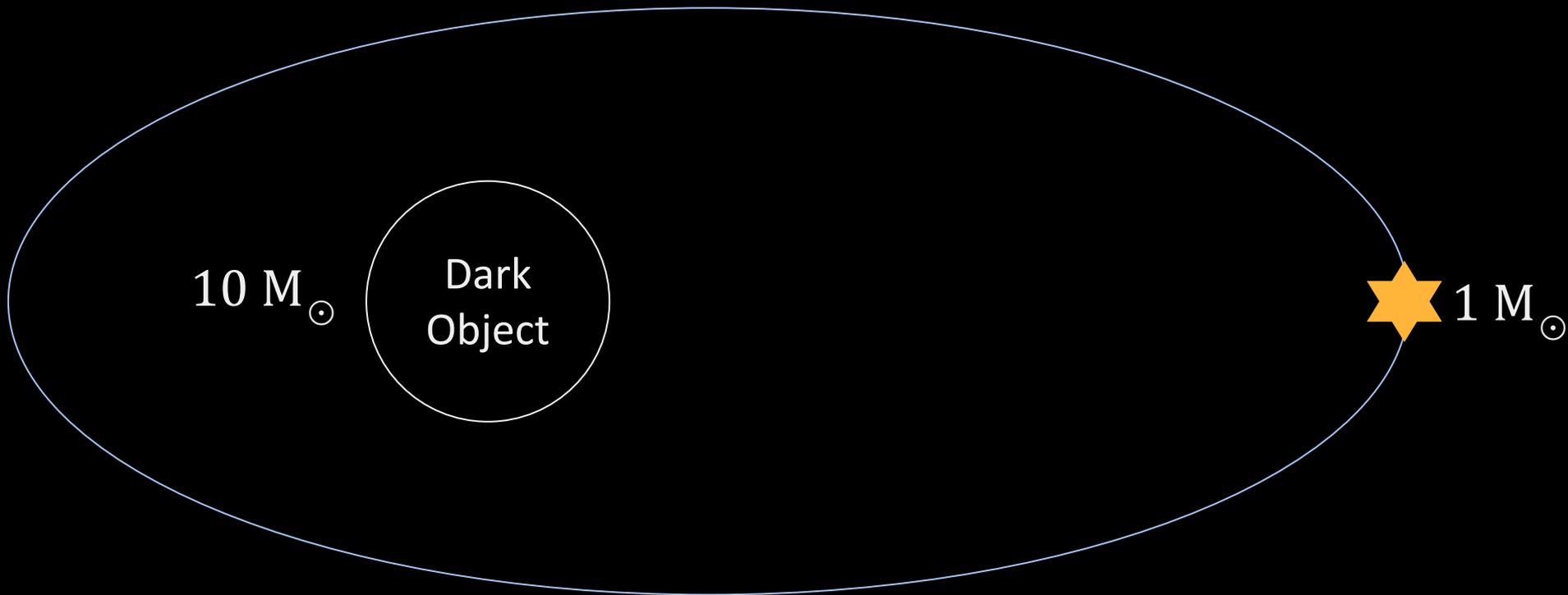
Observation: GAIA



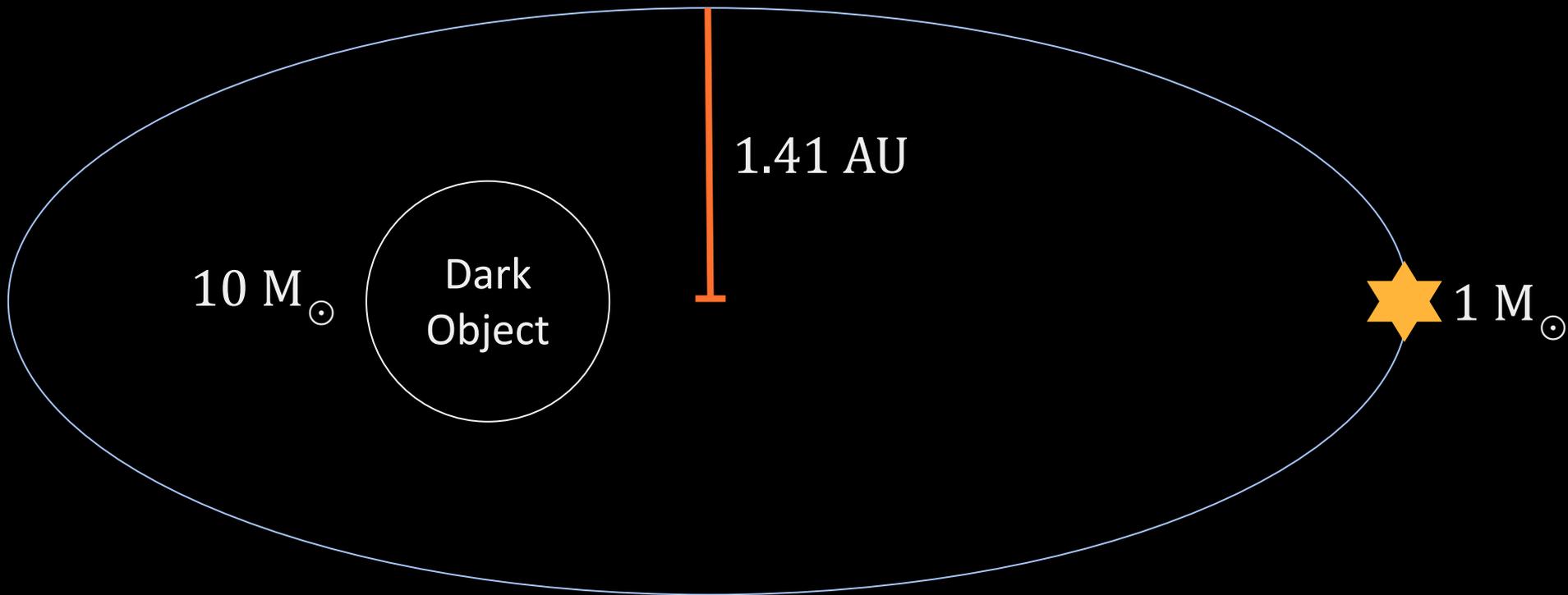
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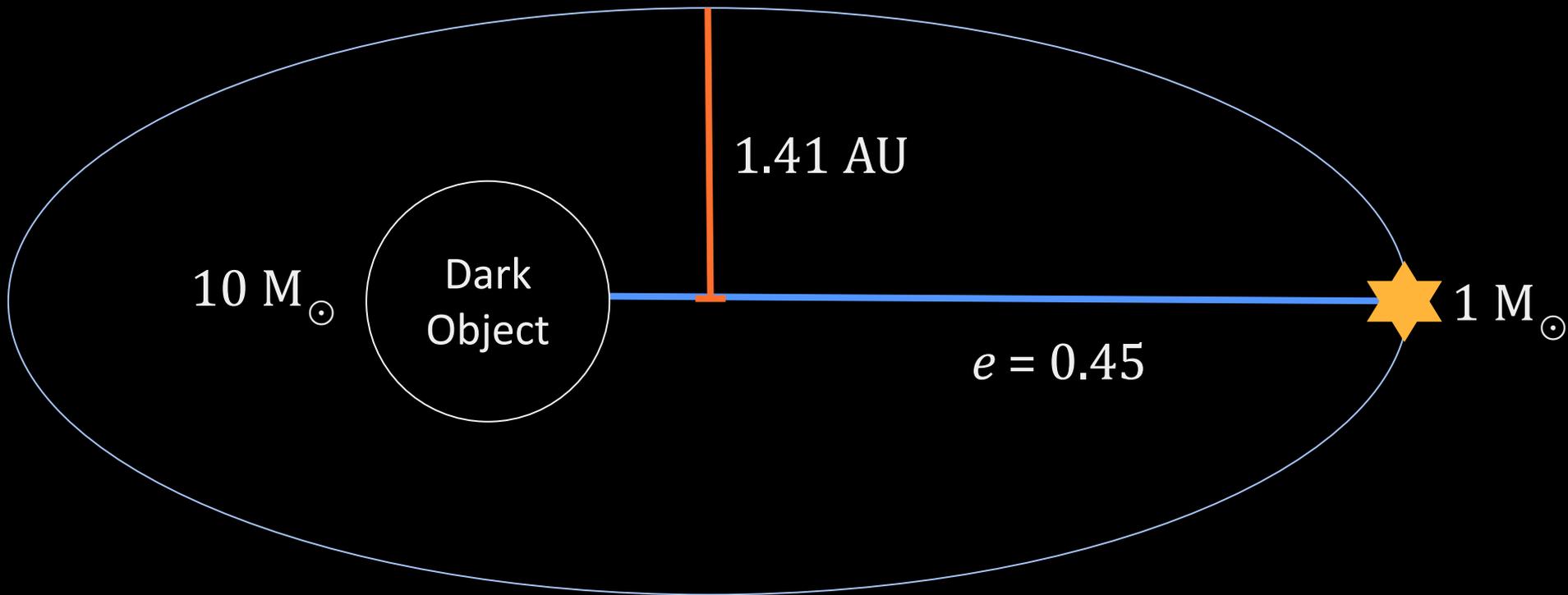
Observation: GAIA



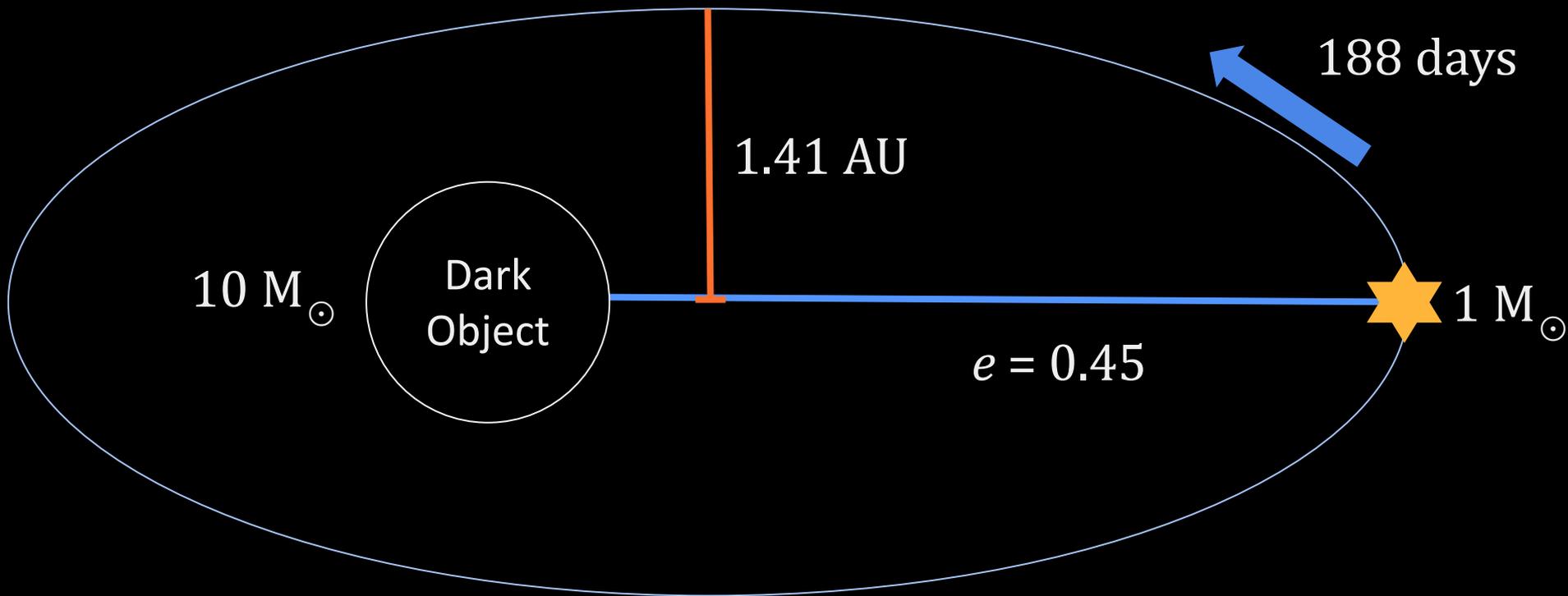
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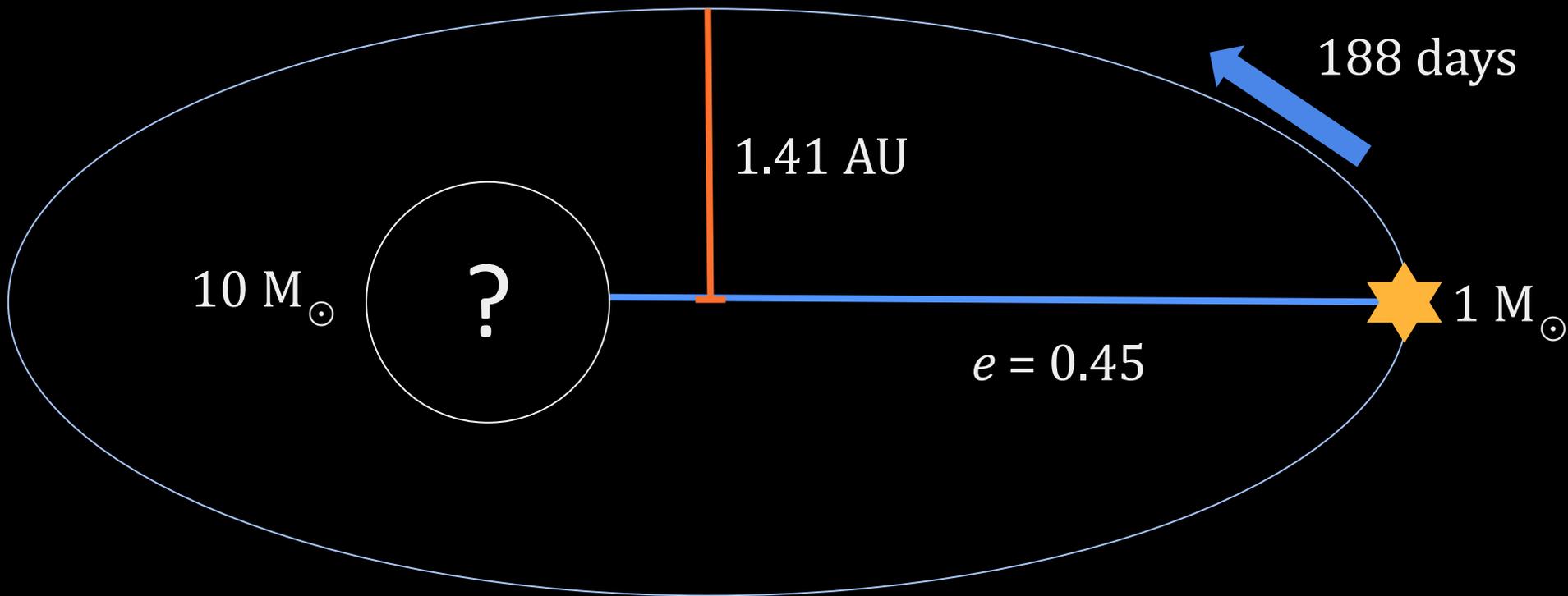
Observation: GAIA



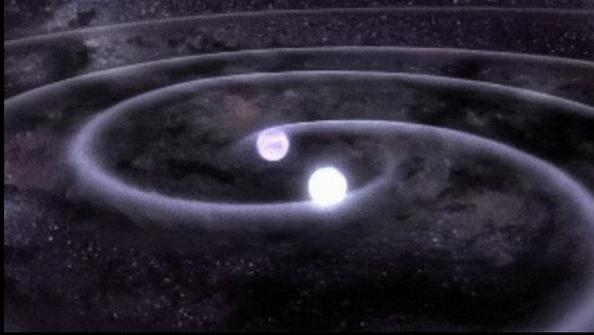
Observation: GAIA



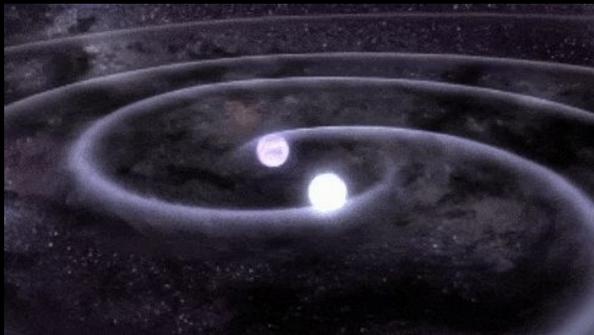
Observation: GAIA



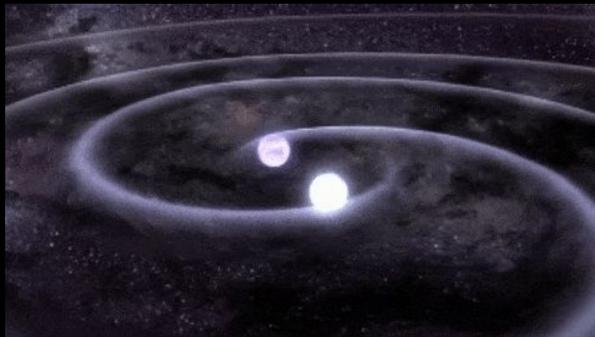
System



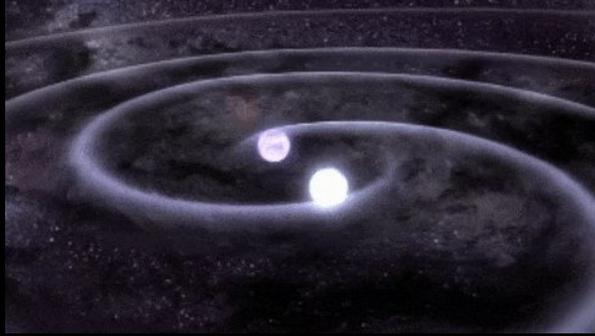
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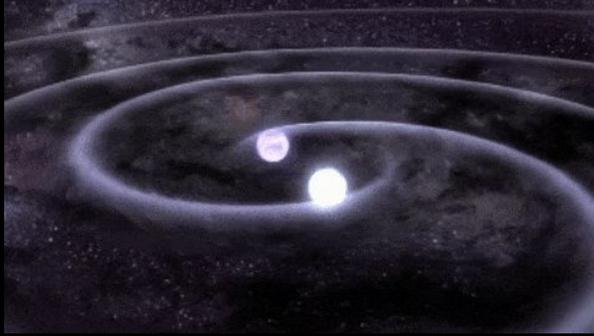
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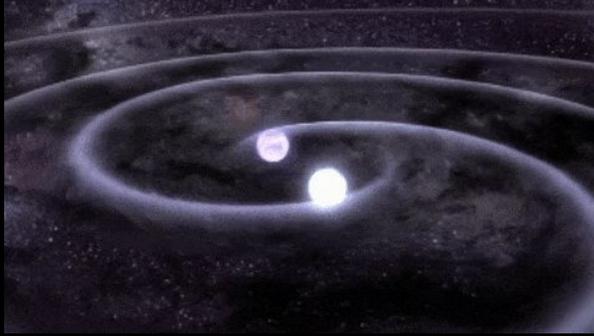


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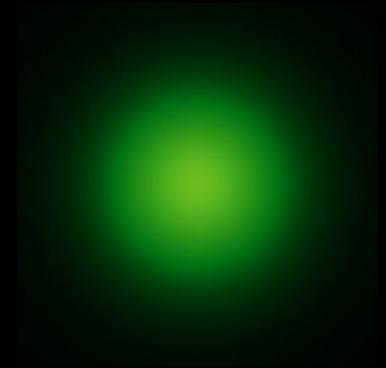
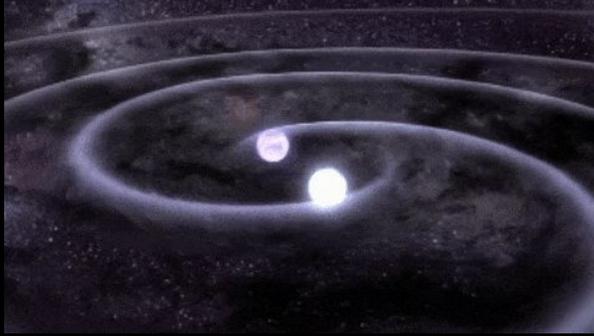
KNOWLEDGETHROUGHSCIENCE

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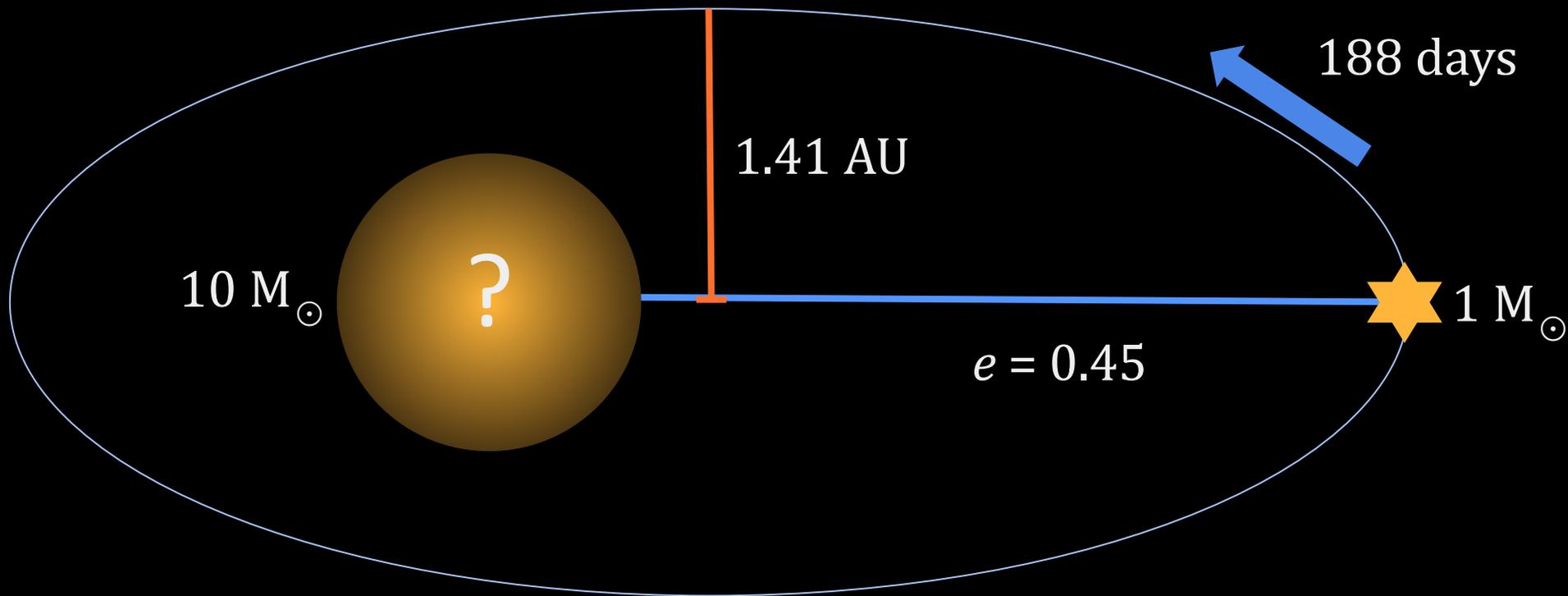
Mass range
Formation
Stability

System

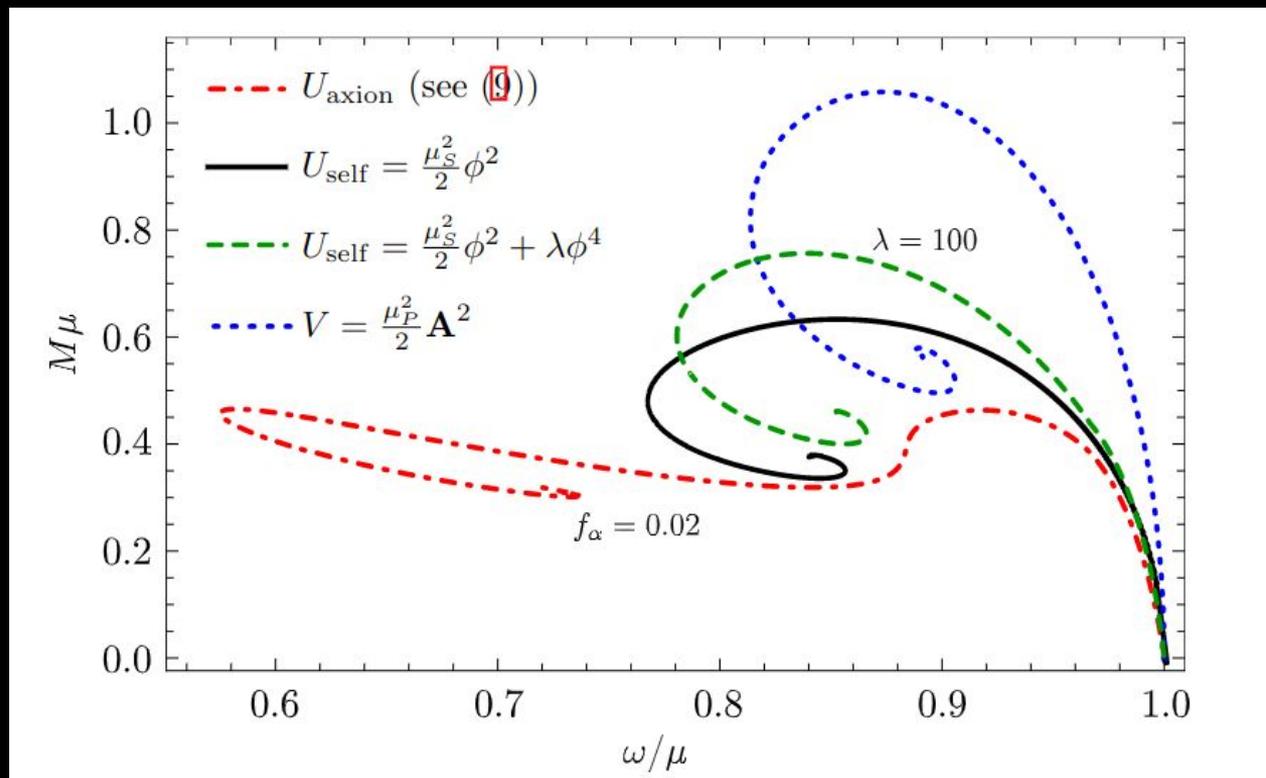


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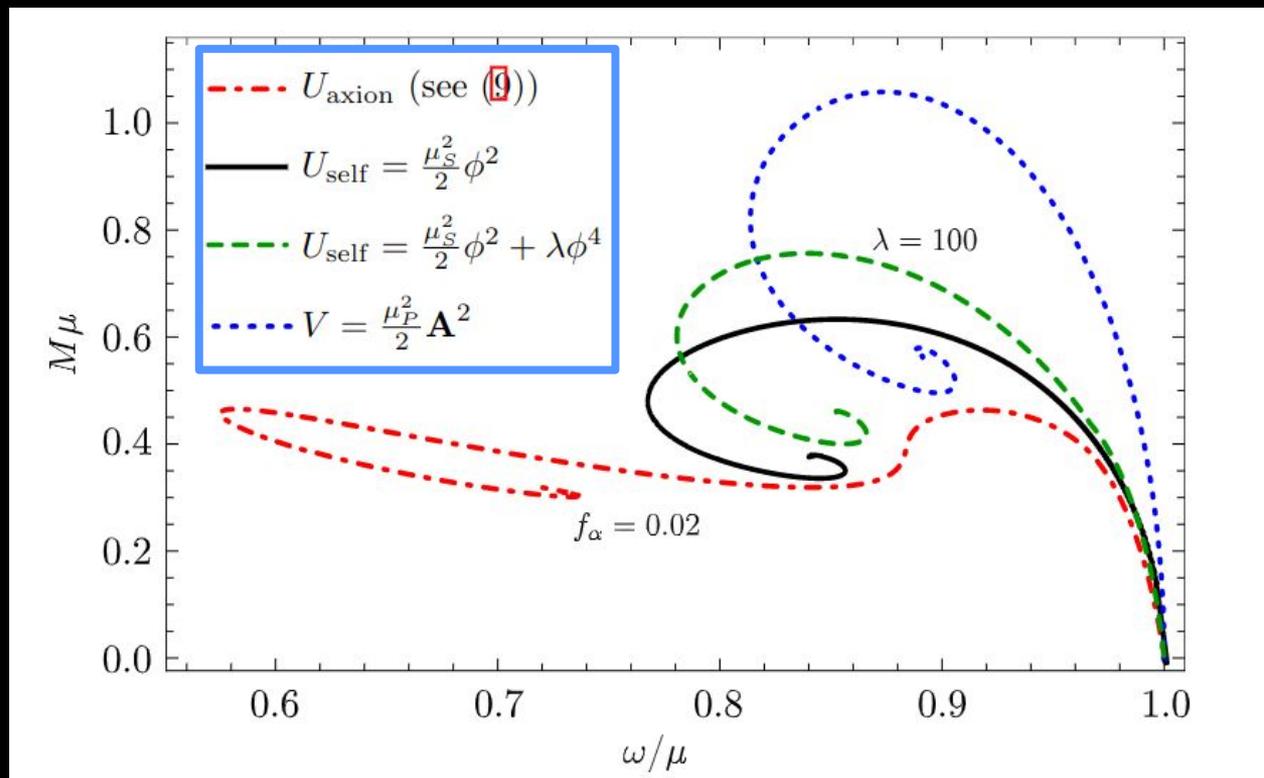
Observation: GAIA



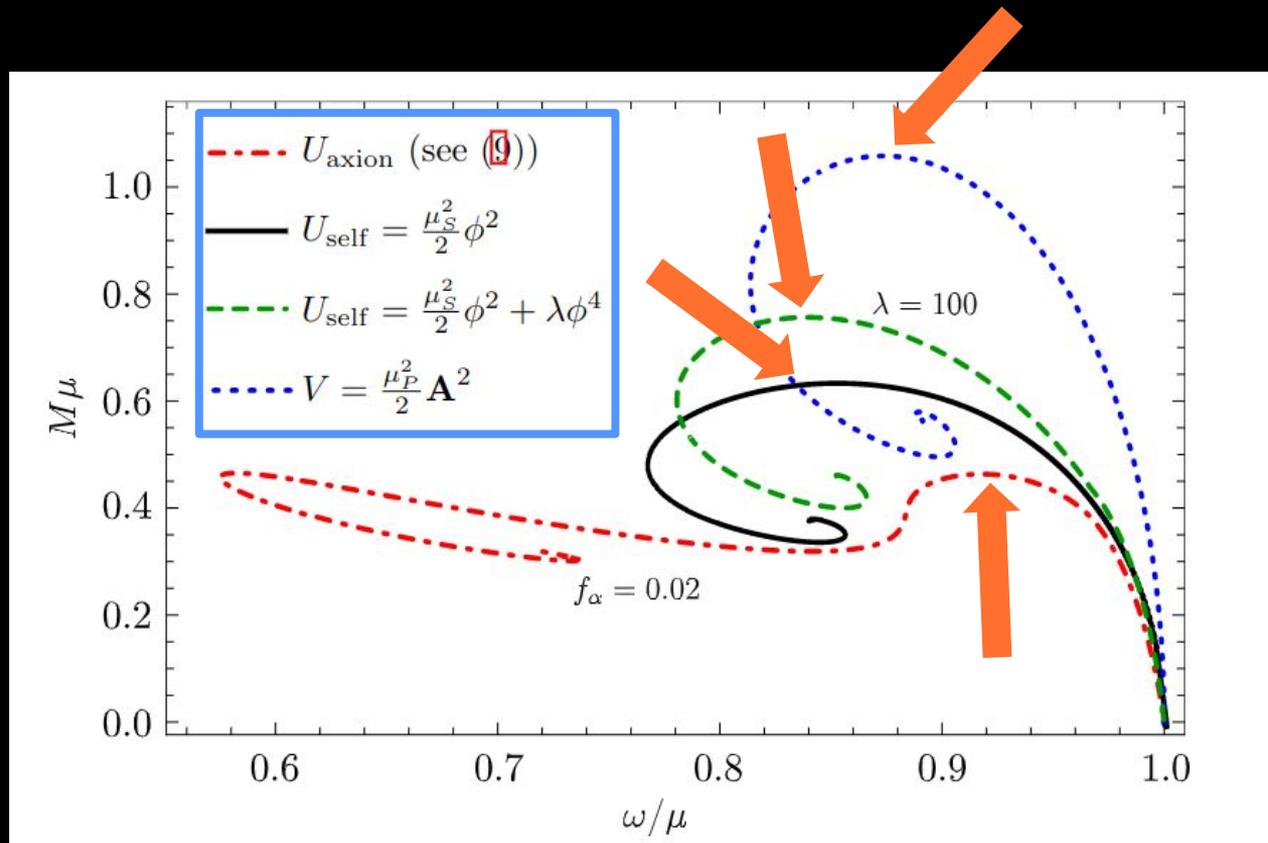
Boson Star



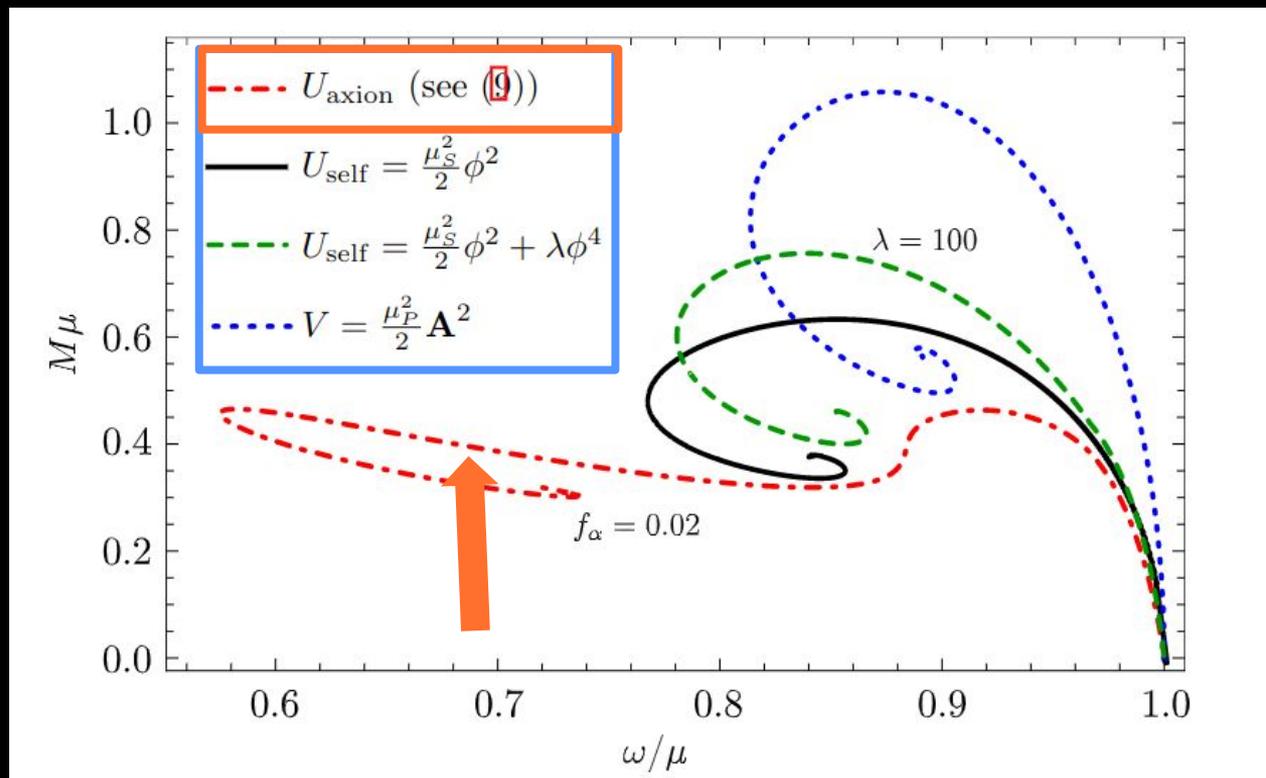
Boson Star



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Boson Star



Orbits

Geodesic motion: Timelike Circular Orbits $k = -1$

- Let us now consider timelike geodesics ($k = -1$)
- Before we wanted the last circular stable orbits
- Now we want an elliptic orbit
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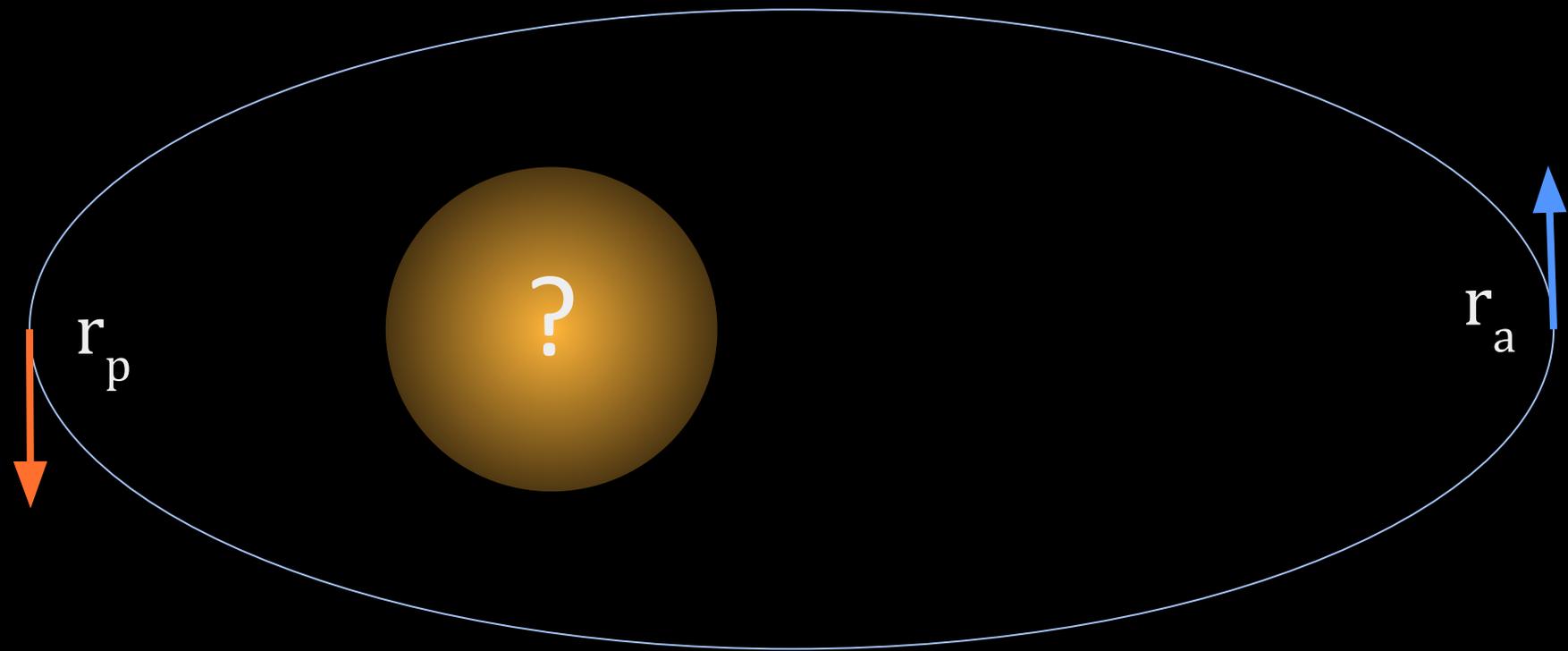
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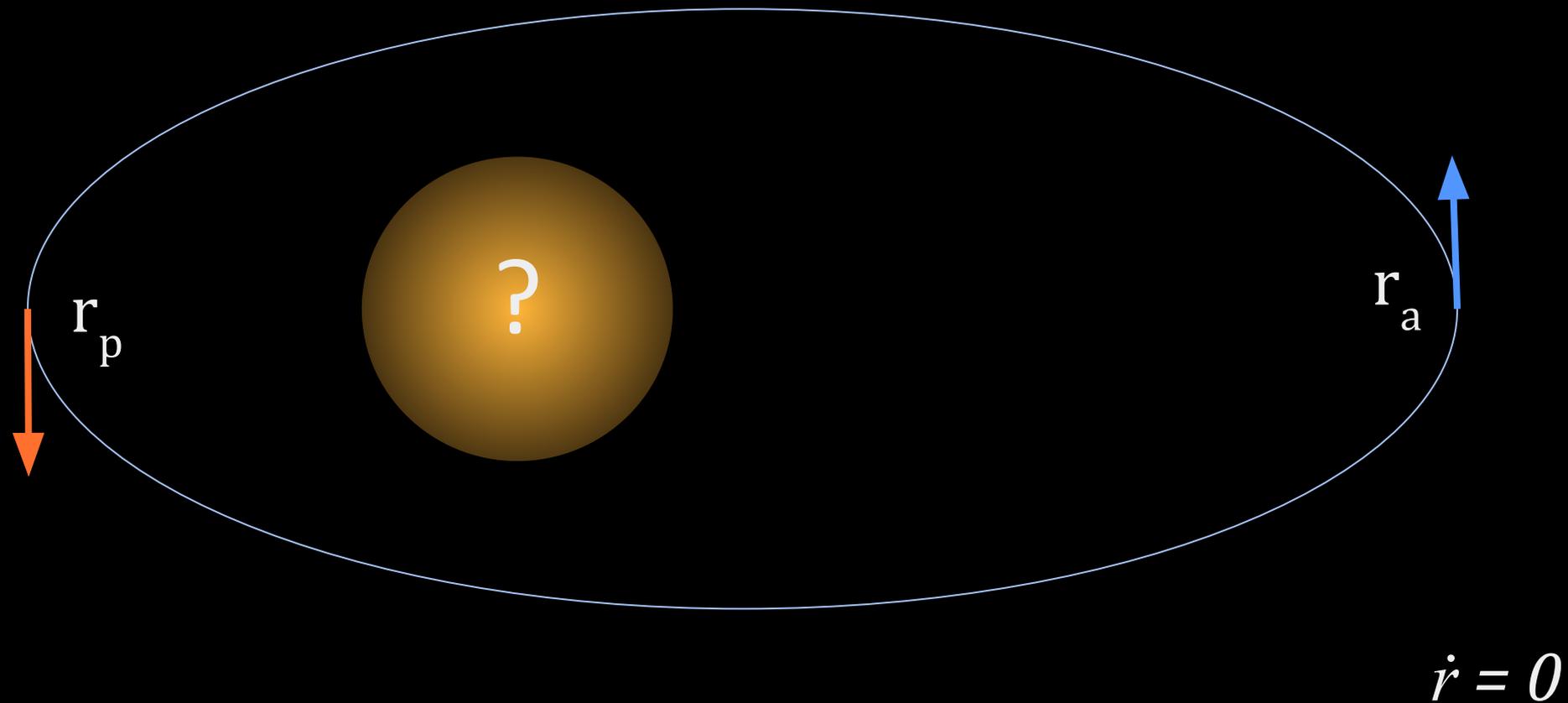
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$$c = 63241 \text{ AU} \cdot \text{yr}^{-1}, \quad G = 39.748 \text{ AU}^3 \cdot M_\odot \cdot \text{yr}^{-2}, \quad M_{\text{Pl}} = 1.094 \times 10^{-38} M_\odot$$

Orbits: Class I

Scalar

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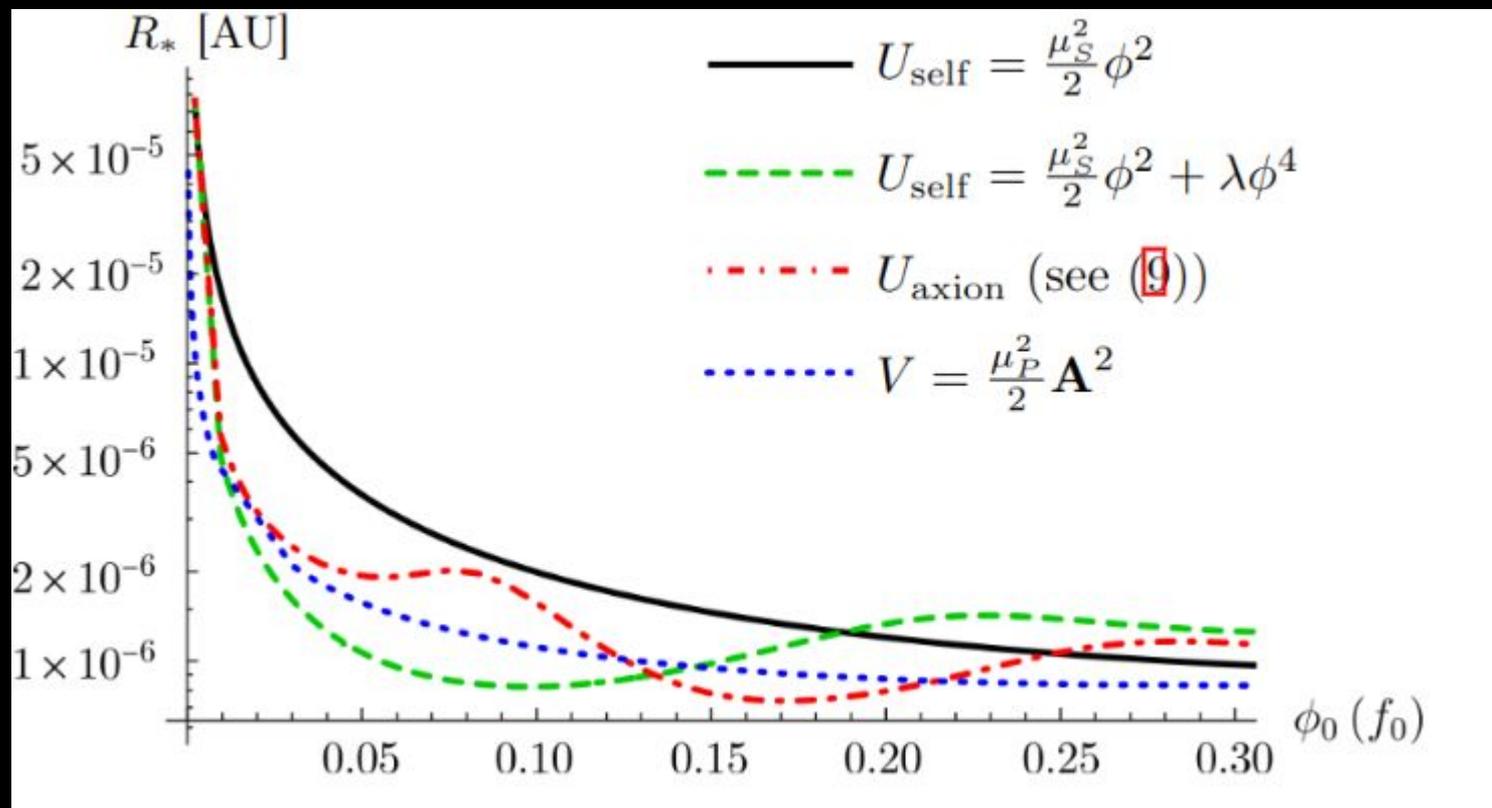
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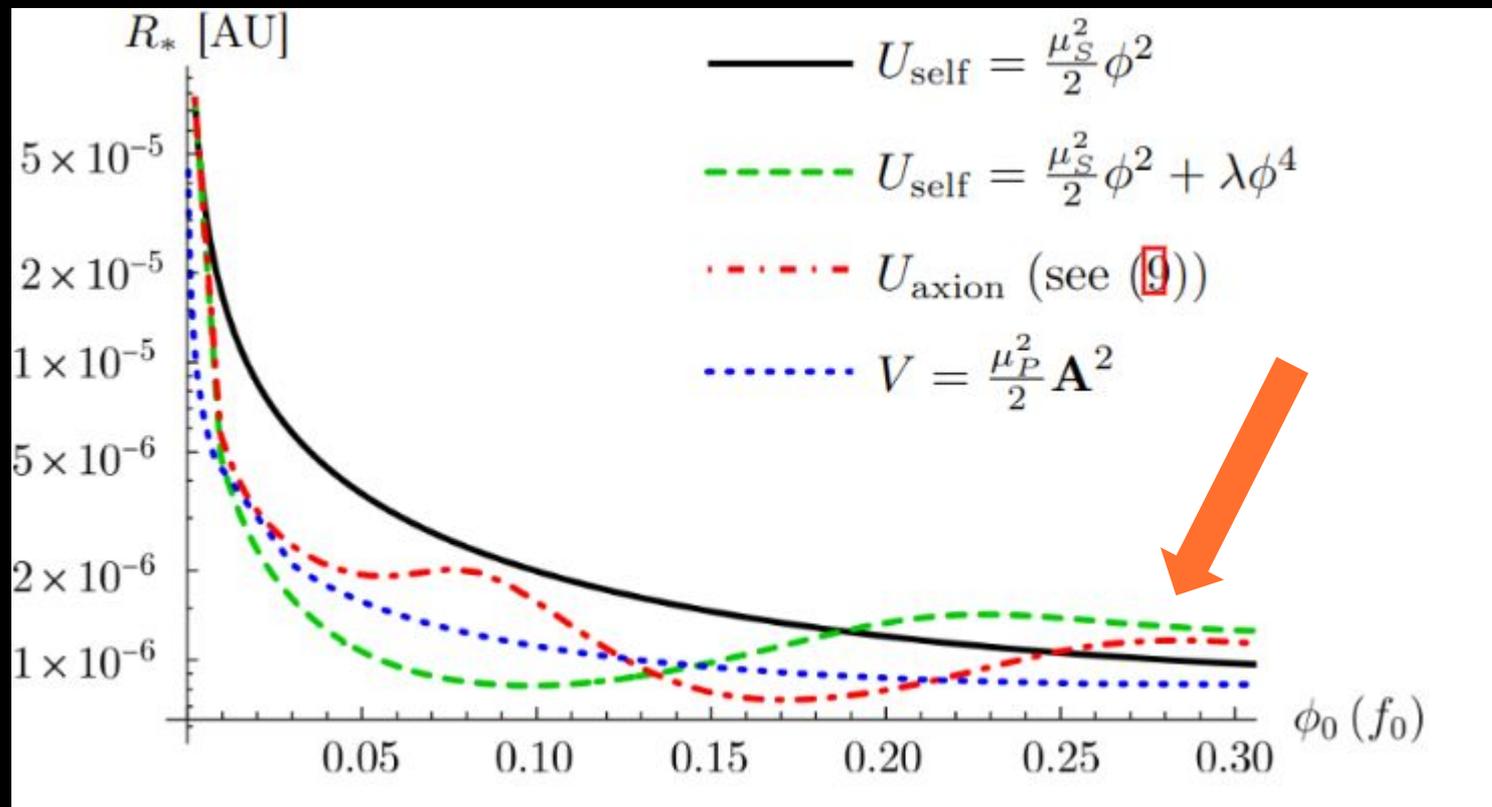
Vector



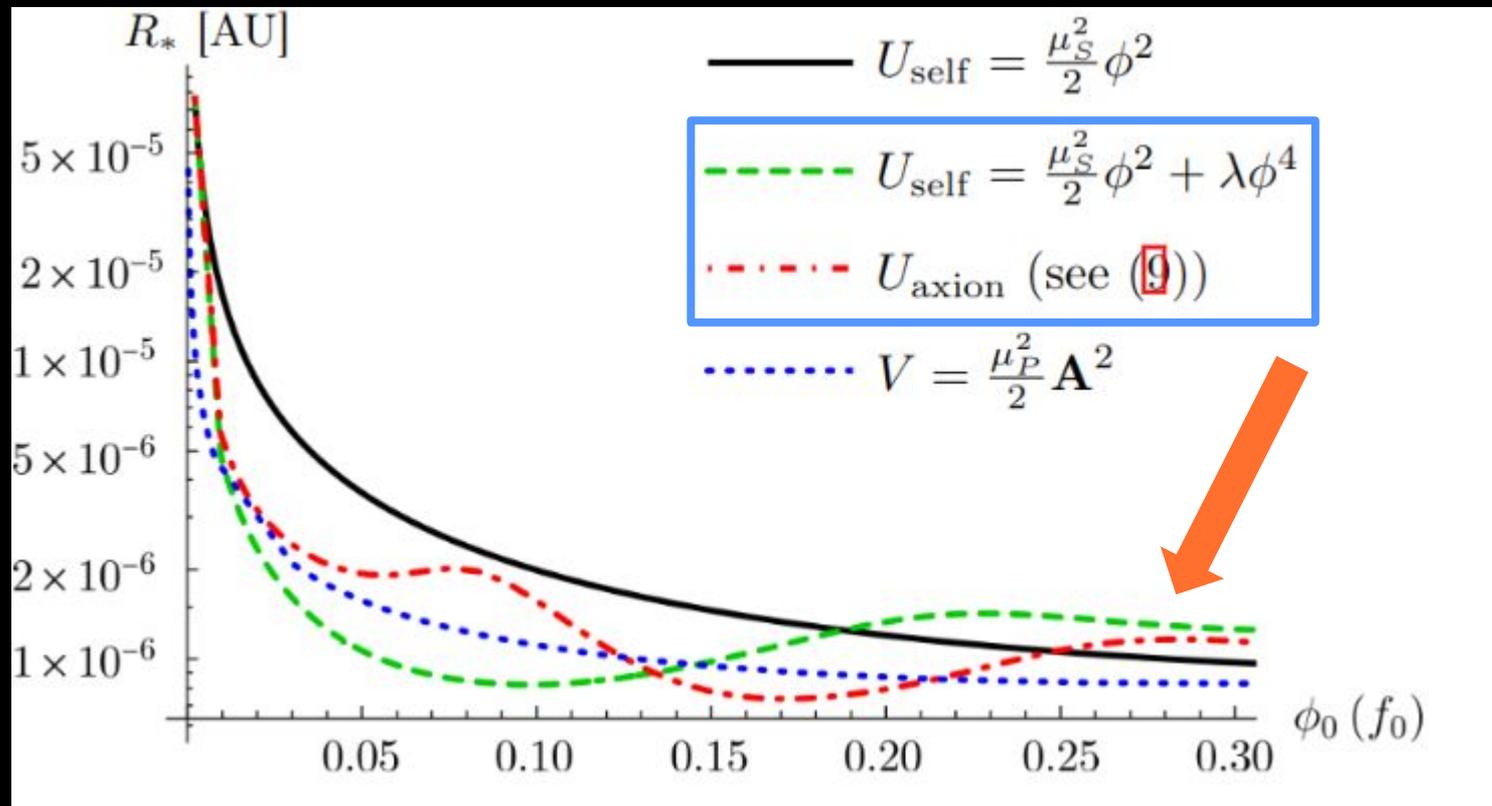
Orbits: Class I



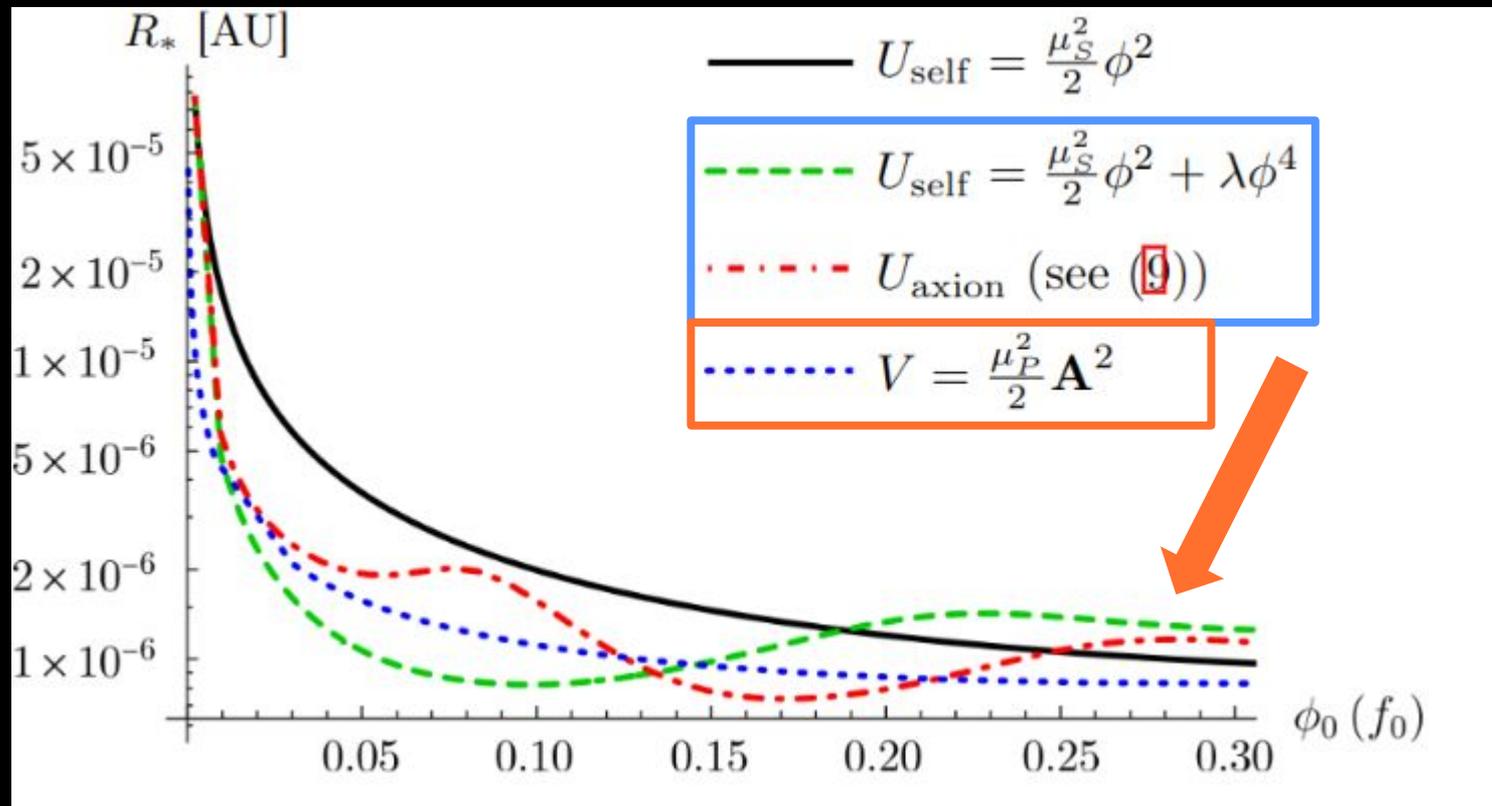
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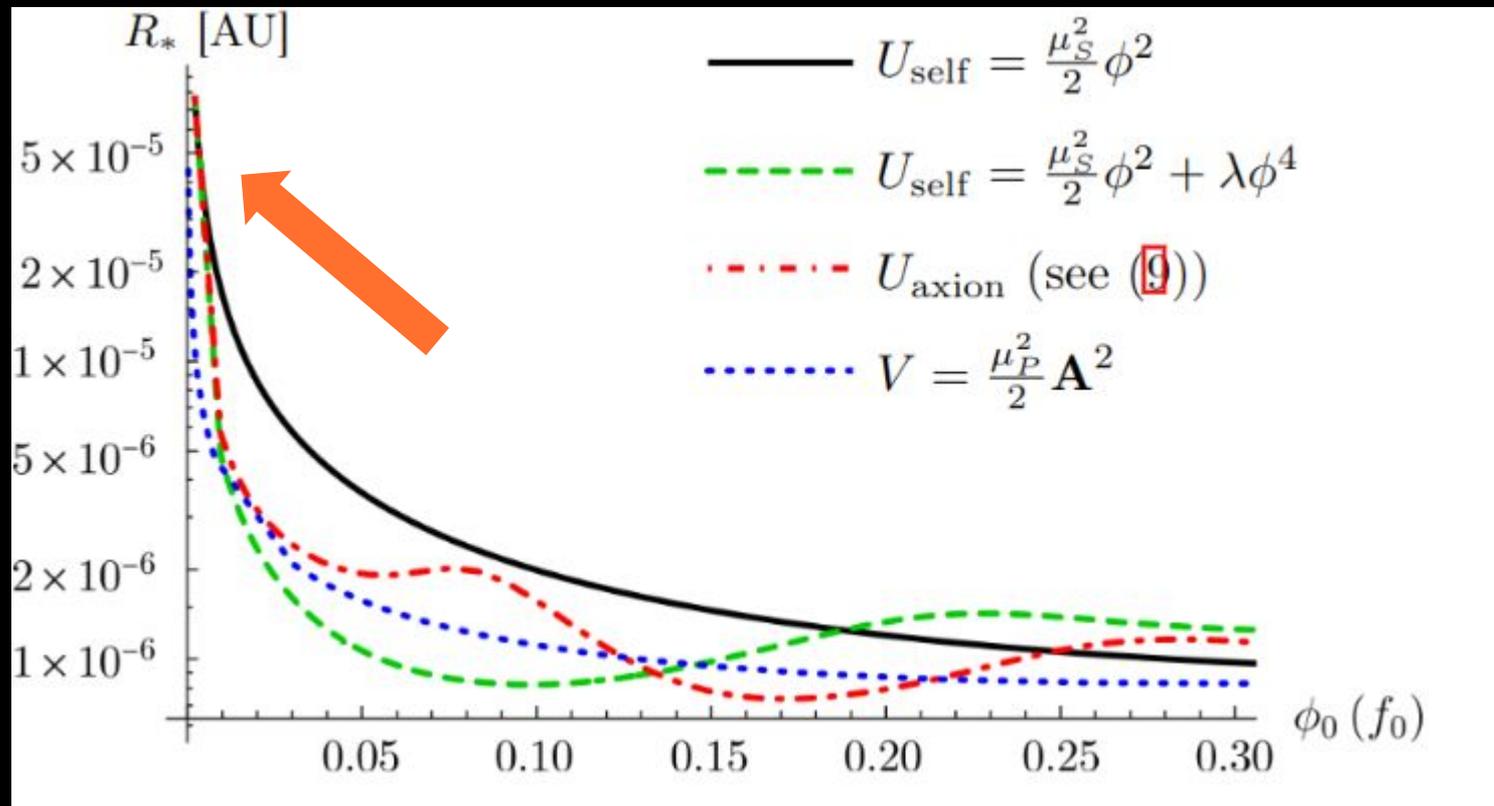
Orbits: Class I



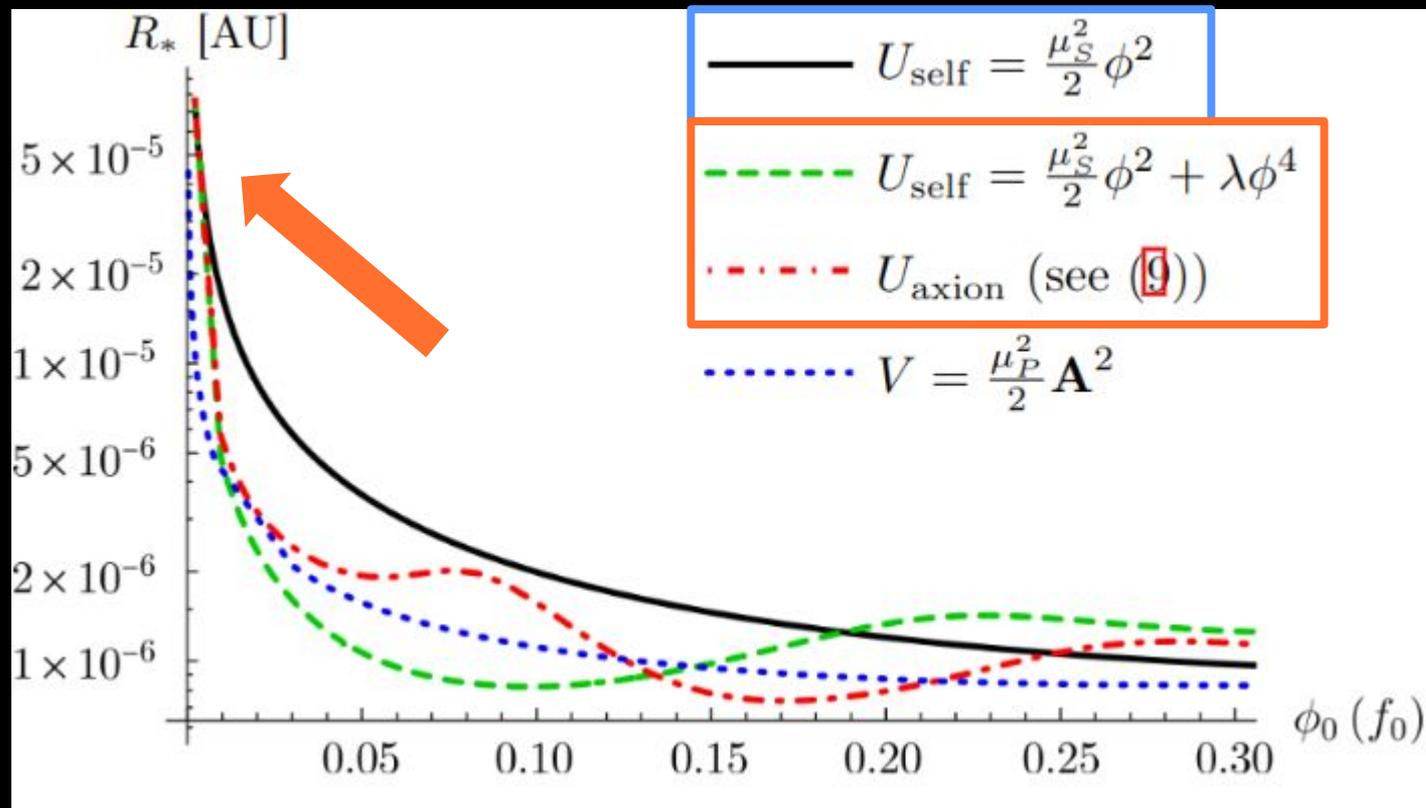
Orbits: Class I



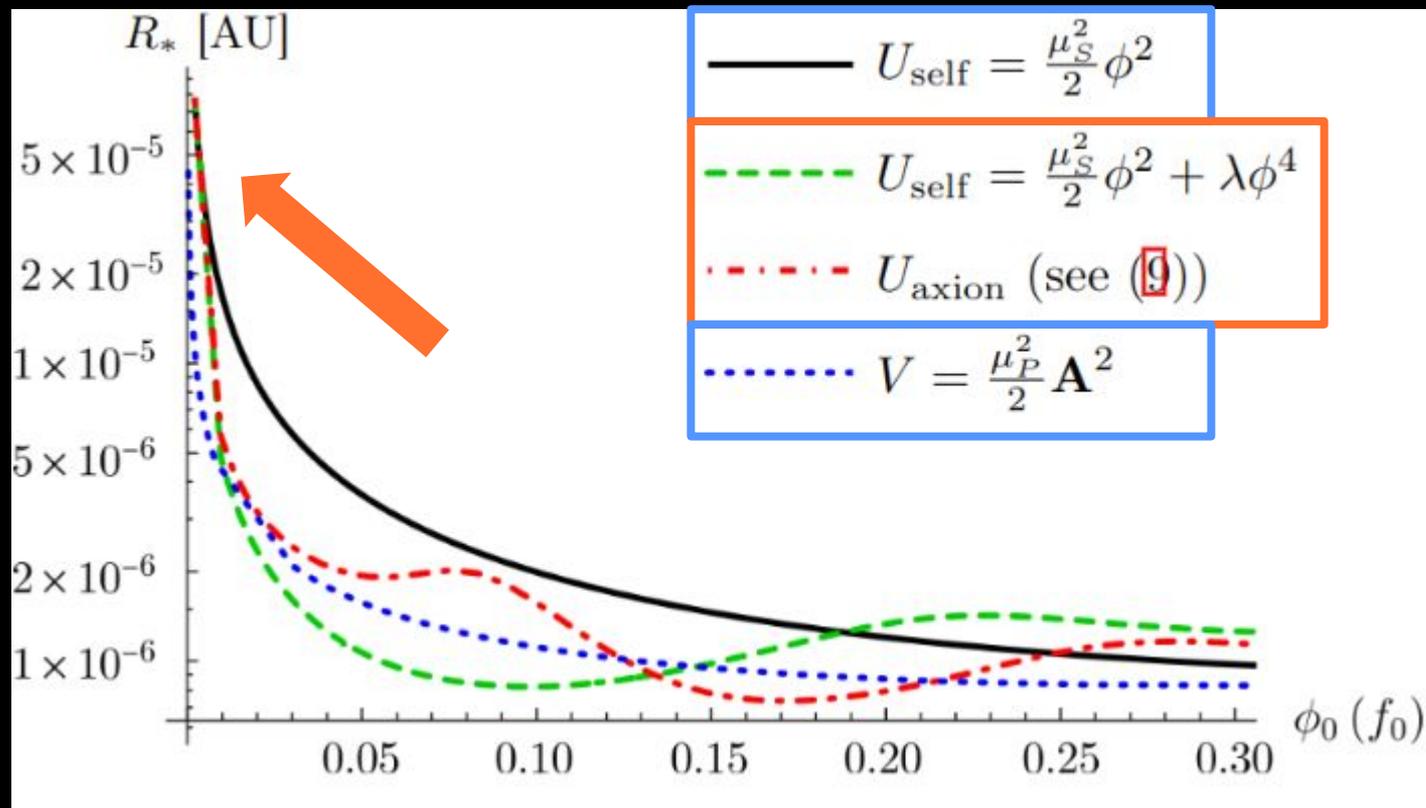
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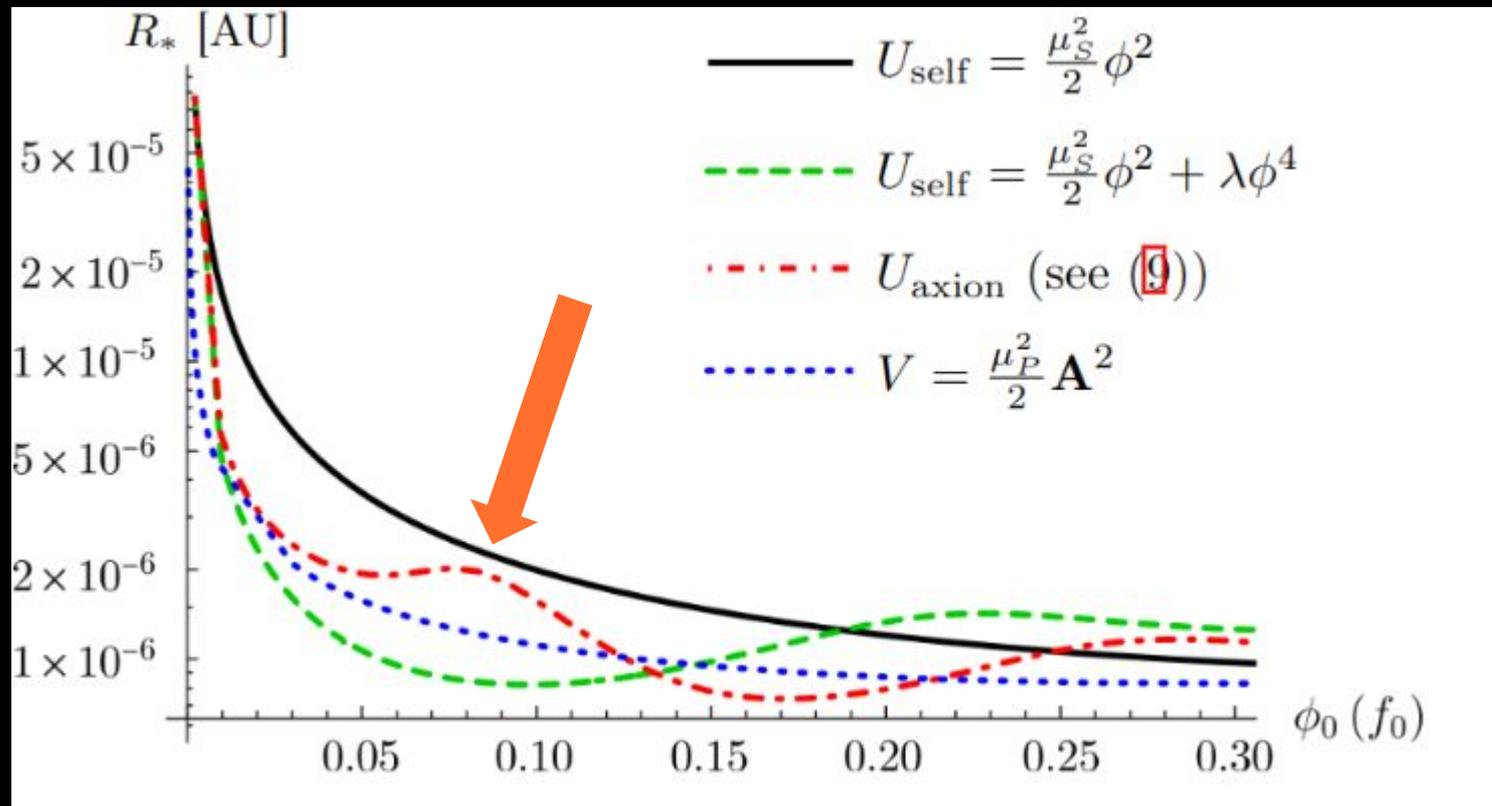
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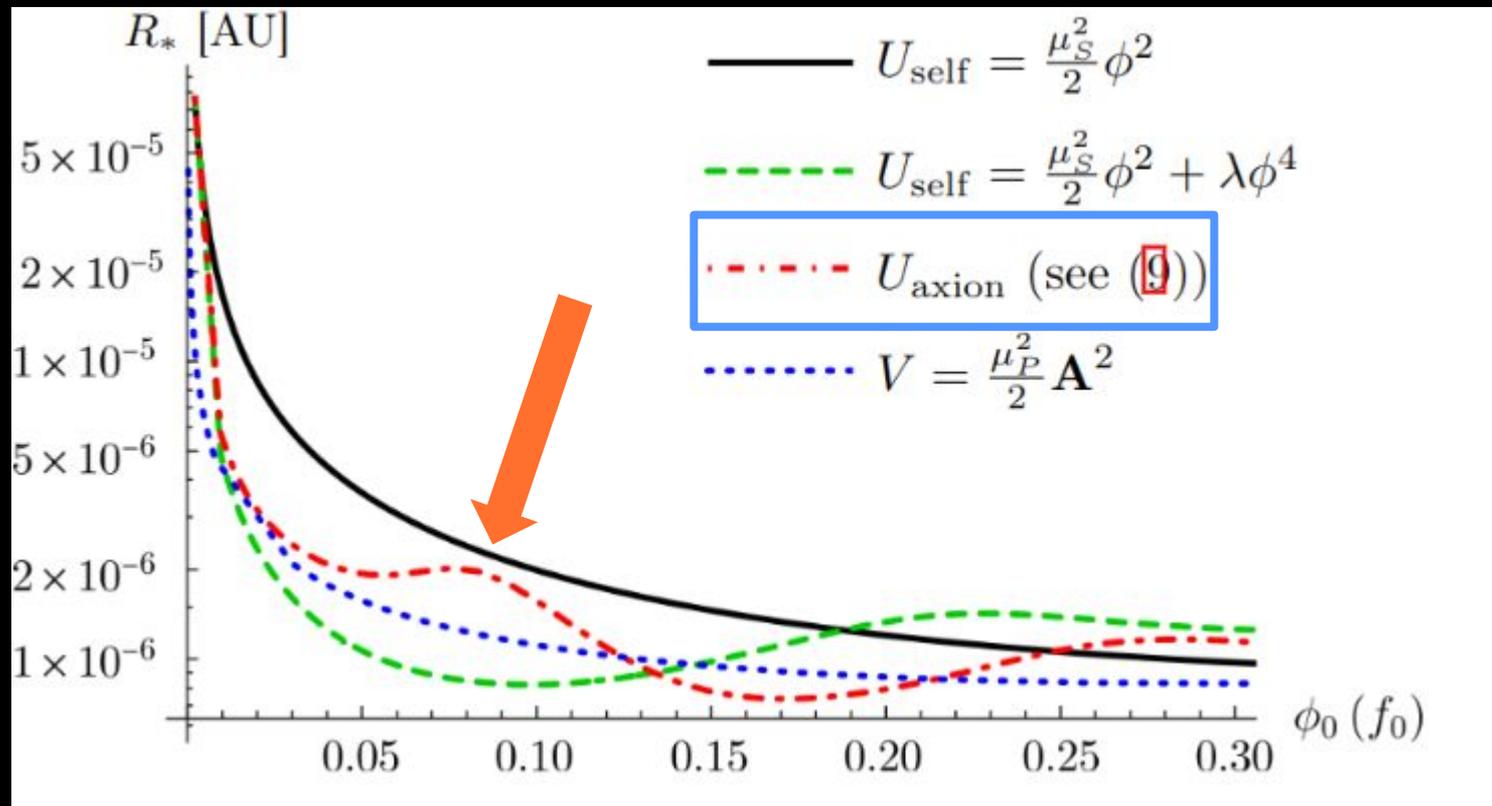
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Boson star characteristics - Case I

Model	Potential	Boson mass (μ) at M_{Max} [GeV/ c^2]	Radius R_* at M_{Max} [km]	Interaction coupling (λ)
Scalar, non-interacting	$U_{\text{self}} = \frac{\mu_S^2}{2} \phi^2$	10^{-20}	134	$\lambda = 0$
Scalar, self-interacting	$U_{\text{self}} = \frac{\mu_S^2}{2} \phi^2 + \lambda \phi^4$	10^{-20}	87	$\lambda = 100$
Scalar, axion-like	U_{axion} (see (9))	10^{-19}	317	$f_\alpha = 0.02$
Vector	$V = \frac{\mu_P^2}{2} \mathbf{A}^2$	10^{-20}	124	0

Orbits: Class I

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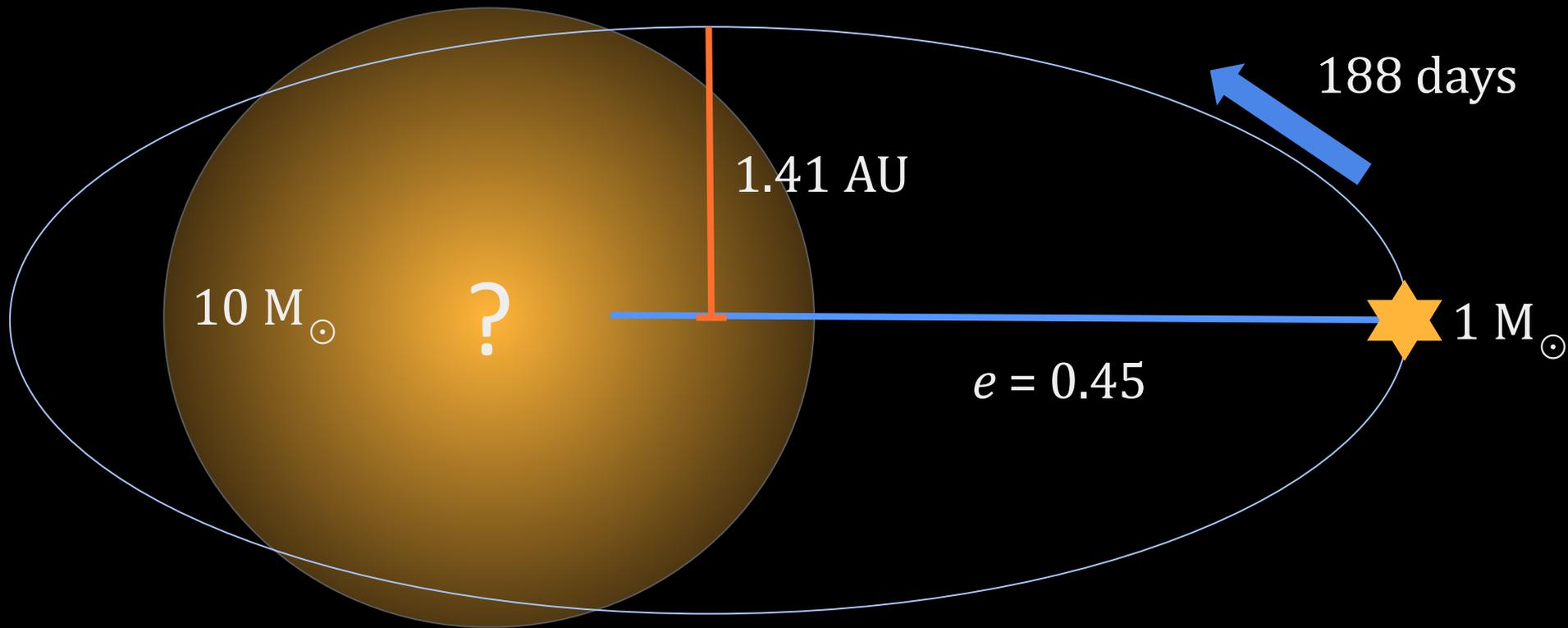
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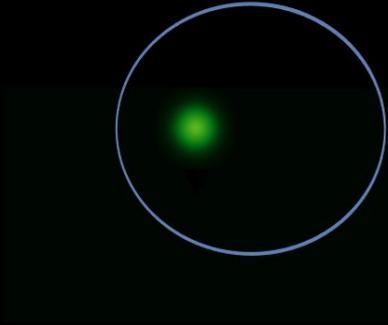
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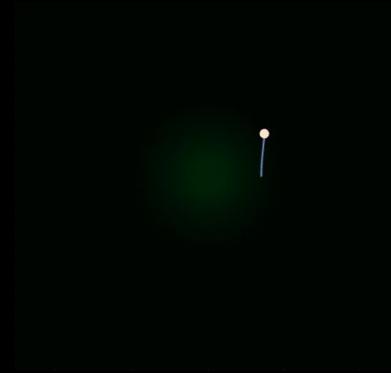
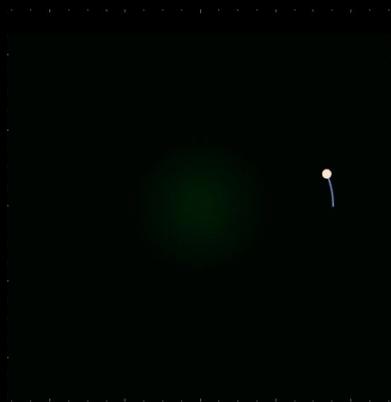
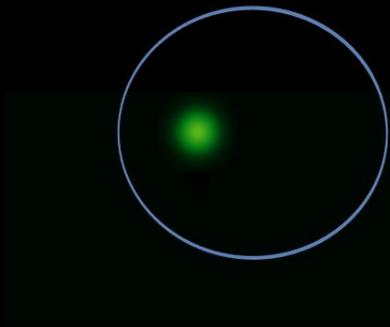
Orbits: Class II



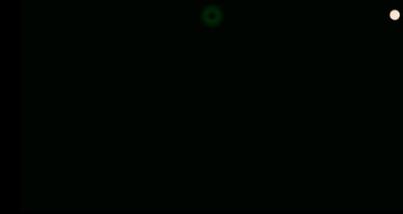
Class II: Scalar



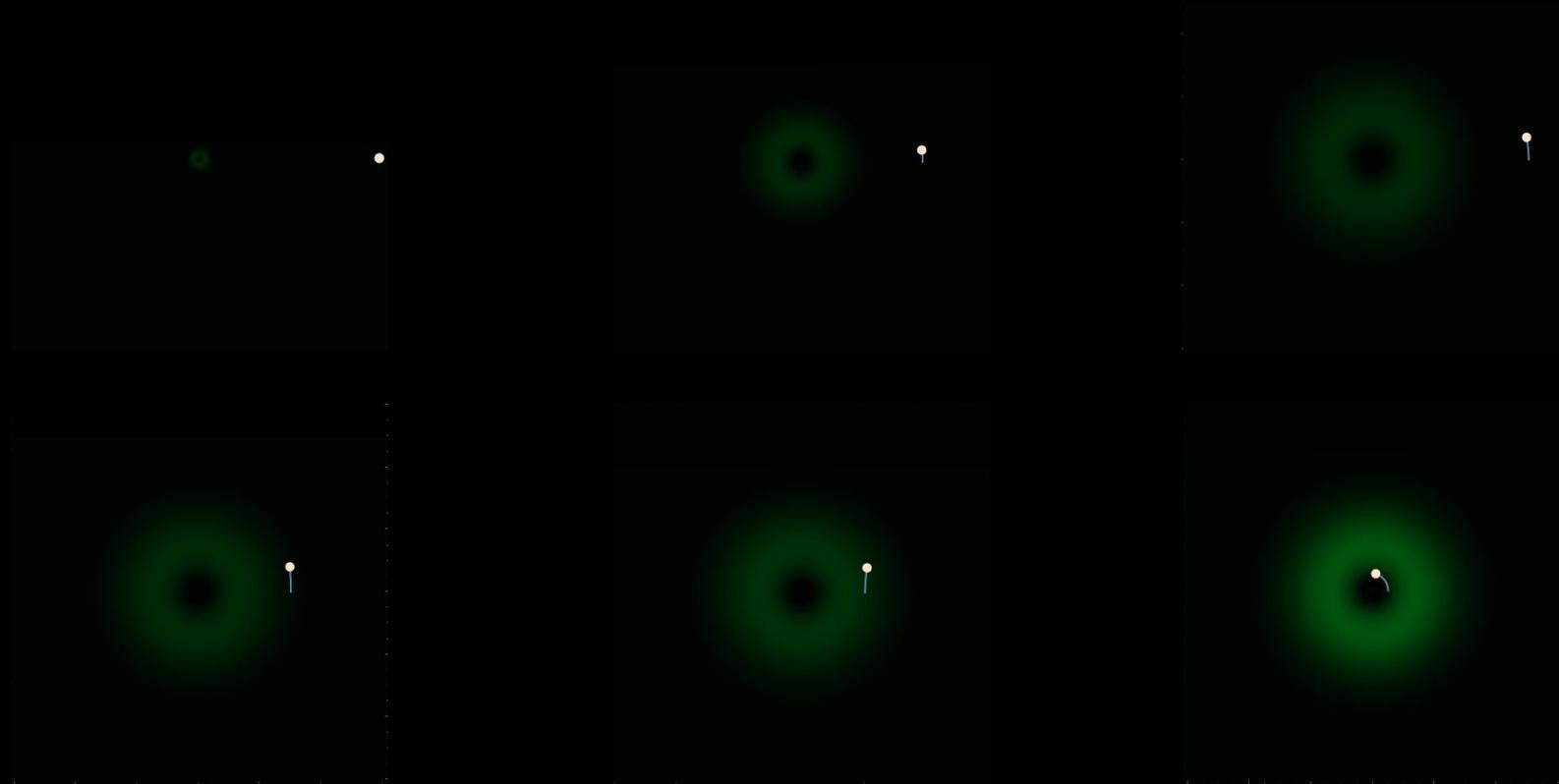
Class II: Scalar



Class II: Vector



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Summary

- Models with dynamically robust spherical BSs, can reproduce the observations made by GAIA
- The largest boson star radii occurs in the stable region
- In this part, the mass is the leading term

- Axionic stars may have the orbiting star inside the matter distribution
- In this case, orbits become highly eccentric and precessing
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Conclusion

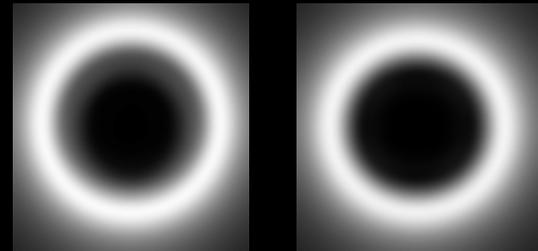
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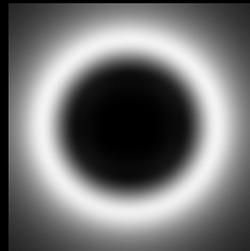
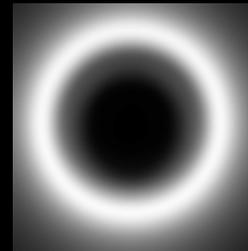


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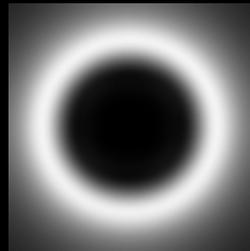
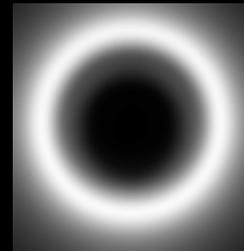
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But only under certain observational conditions



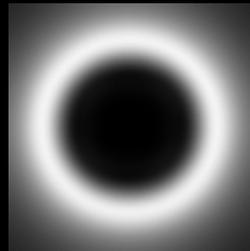
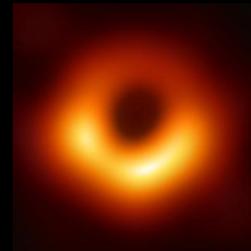
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A Boson star can mimic a BH's phenomenology



Thank You!

Děkuji!



Boson Stars that mimic Black Holes

doi.org/10.1088/1475-7516/2021/04/051

arXiv:2304.09140

pombo@fzu.cz

