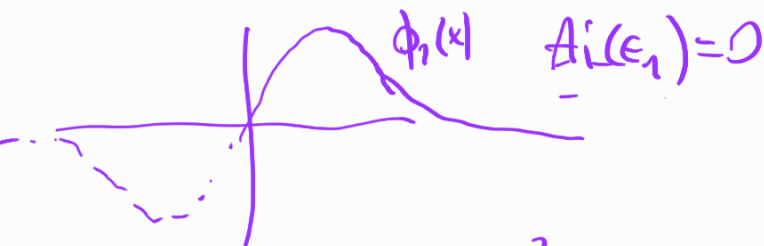
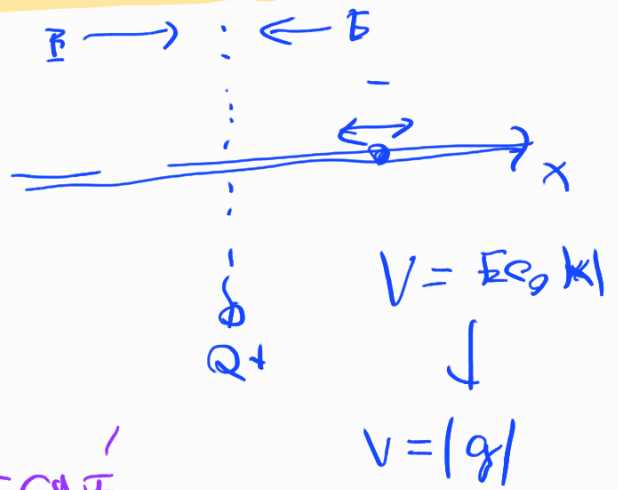
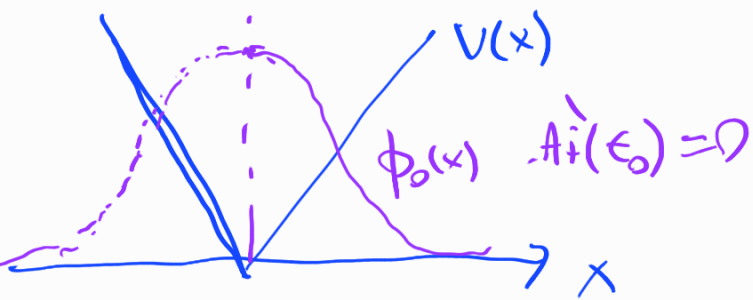


# QM II - cv 3 Variacím metoda a WKB



PŘESNĚ

•  $\psi_0(q) \approx e^{-\alpha q^2}$

$$E[\psi] = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} = f(\alpha)$$



• první excit. stav: ~~zákl.~~ stav  $\mathcal{R} = \mathcal{R}_S \oplus \mathcal{R}_L$

$\psi_1(q) = q e^{-\alpha q^2}$

$$E[\psi_1] = \frac{\langle \psi_1 | H | \psi_1 \rangle}{\langle \psi_1 | \psi_1 \rangle} = \frac{2}{\sqrt{2\pi}} \frac{1}{\sqrt{\alpha}} + 3\alpha$$

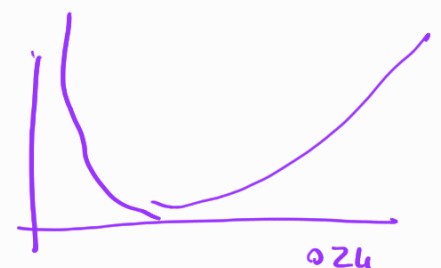
$$\min E'(\alpha) = 3 - \frac{1}{\sqrt{2\pi}} \alpha^{-3/2} = 0$$

$$\alpha_0 = (3\sqrt{2\pi})^{-2/3}$$

$E_1 \approx E[\alpha_0] = 2.3448 \dots$  Přesně: 2.3338

$\Delta E \approx 0.006$

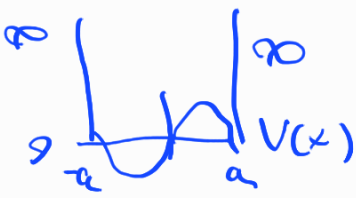
• základní stav  $\psi_\alpha(x) = \frac{1}{x^2 + \alpha^2}$



$E[\alpha] = \frac{1}{2} \frac{1}{\alpha^2} + \frac{2}{\pi} \alpha \dots \alpha = \sqrt[3]{\frac{\pi}{2}} \dots E_0 = 1.11006$

**ULOHA 2**

Částice v kruhovém potenciálu jámě



$a=1$

$E = \frac{\hbar^2}{2m} \frac{1}{a^2}$

$\hat{H} \psi = E \psi$

... uvně je  $\psi = 0$  mě jámy

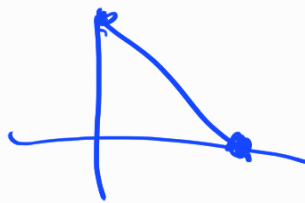
$\psi \in L^2$  (úskla polární 1) } třída řešení  
 $\psi(r=1) = 0$

①  $\psi(x, y) = \sin\left(\frac{n}{\pi} \theta\right)$

②  $\psi(x, y) = \cos\left(\frac{2m}{\pi} \theta\right) \rightarrow E[\psi]$

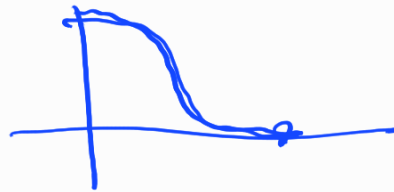
$f(r) \nabla f = f'(r) \vec{e}_r$

$\sin^2 \cdot r \, dr$   
 $\cos^2 \cdot r \, dr$



③  $\psi(x, y) = (1-r)$

$(1-r)^2$



$m=0$

$e^{im\theta} = 1$

= polynom v  $(1-r)$

Excit. stavy:  $\psi = a\phi_1 + b\phi_2$   $\begin{pmatrix} \langle \phi_1 | \phi_1 \rangle & \langle \phi_1 | \phi_2 \rangle \\ \langle \phi_2 | \phi_1 \rangle & \langle \phi_2 | \phi_2 \rangle \end{pmatrix}$

Hyleraas-unděin rel. č.  $E_0, E_1$  approx. zákl. a it sta.

symetrie: rotační symetrie

$L = i\hbar \partial_\theta \quad [\hat{H}_1, L] = 0 \quad \dots e^{im\theta}$

$$\{ \hat{H}, \hat{L} \} \quad \psi(r, \varphi) = f(r) e^{im\varphi}$$

různé hodnoty  $m \rightarrow \mathcal{H} = \mathcal{H}_{m=0} \oplus \mathcal{H}_{m=1} \oplus \mathcal{H}_{m=2}$

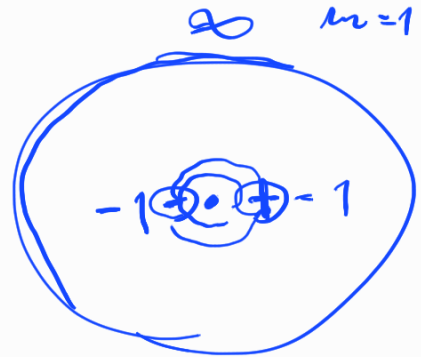
$$E_{m=0,1,2,\dots}^{(m)}$$

Ekv. stav  $m = \pm 1$  ---  $E_{m=0}^{(m)}$

$$\psi(r, \varphi) = f(r) e^{im\varphi}$$

$f(r)|_{r=0} = 0$  ---  $f(0) = 0$

$$f(r) = r(1-r) \quad \text{--- } p\{r(1-r)\}$$



Přesná řešení:  $-\Delta \psi(x, y) = E \psi(x, y)$  v obl.  $\mathbb{R}^2$   
 $\psi(r=1) = 0$

$$\psi(r, \varphi) = f(r) e^{im\varphi} \quad m \in \mathbb{Z}$$

Laplace v polárních souř.

$$\frac{\partial^2}{\partial r^2} \psi + \frac{1}{r} \frac{\partial}{\partial r} \psi + \frac{1}{r^2} \frac{\partial^2}{\partial \varphi^2} \psi + E \psi = 0$$

$$\psi(x) = f(kr) e^{im\varphi} \quad \text{--- } E = k^2 \quad z = kr$$

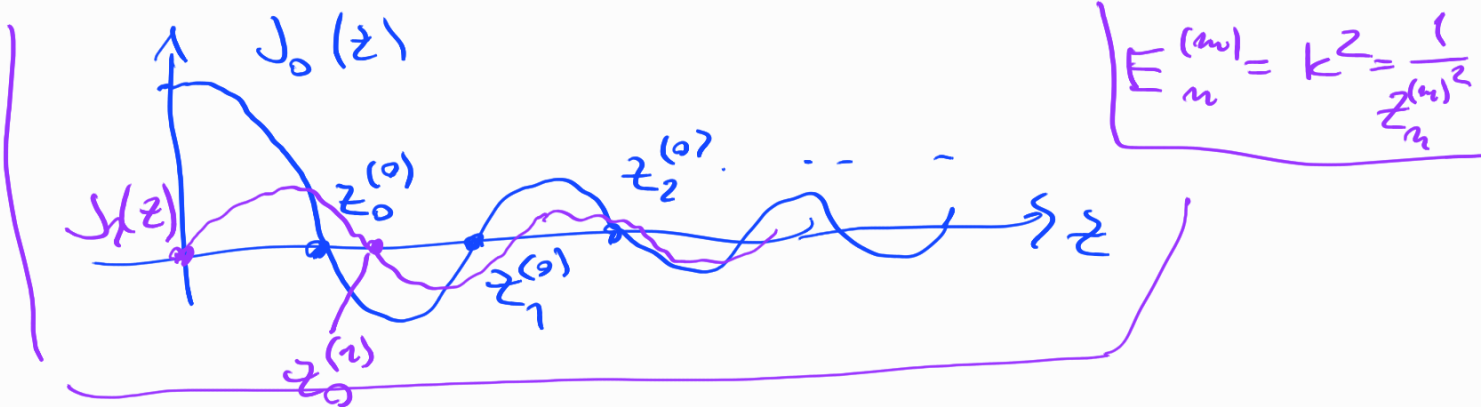
$$f''(z) + \frac{1}{z} f'(z) + \left[ 1 - \frac{m^2}{z^2} \right] f = 0$$

Besselova rovnice  $J_m(z)$

$$\Rightarrow \psi(x, y) = J_m(kr) e^{im\varphi} \quad J_m(kr) = 0$$

kořeny Bessel. fce  $J_m(z_n^{(m)}) = 0$   $n=1, 2, \dots$   
 $n=1 \Rightarrow k = \frac{1}{z_n^{(m)}}$





① Odrůzení (WKB) asymptot Airyho funkce

↑ Duhy

Def:  $Ai''(x) = x Ai(x)$

$\rightarrow Ai = \frac{1}{\pi} \int_0^{\infty} \cos\left(\frac{t^3}{3} + xt\right) dt = \frac{1}{\pi} \int_{-\infty}^{\infty} e^{i\left(\frac{t^3}{3} + xt\right)} dt$

$e^{if(t)}$



$f'(t) = 0 = t^2 + x$   
 $\rightarrow t_0 = \sqrt{-x}$

Metoda stacionární fáze

$f''(t) = 2t_0$

Taylor  $f(t) \approx t_0 \left( x + \frac{1}{3} (-x) \right) + t_0 (t-t_0)^2 = \frac{2}{3} x t_0 + t_0 (t-t_0)^2$

①  $x > 0 \quad t_0 = \pm i\sqrt{|x|}$

$Ai(x) \approx \frac{1}{\sqrt{\pi|x|}} e^{i\sqrt{x} \frac{2}{3}x} \int_{-\infty}^{\infty} e^{(-\sqrt{x})(t-i\sqrt{x})^2} dt = \frac{1}{\sqrt{\pi|x|}} e^{-\frac{2}{3}x^{3/2}}$

Gauss

$Ai(x) \approx \frac{1}{\sqrt{\pi|x|}} e^{-\frac{2}{3}x^{3/2}}$

$x \rightarrow \infty$

(2)  $x < 0$

$\dots t_0 = \sqrt{x}$   
 $\frac{2i}{3} x^{3/2}$

$$Ai(x) = \frac{1}{\sqrt{i\pi|x|}} e^{\frac{2i}{3} x^{3/2}} \cos\left(\frac{2}{3} x^{3/2} + \frac{\pi}{4}\right)$$

