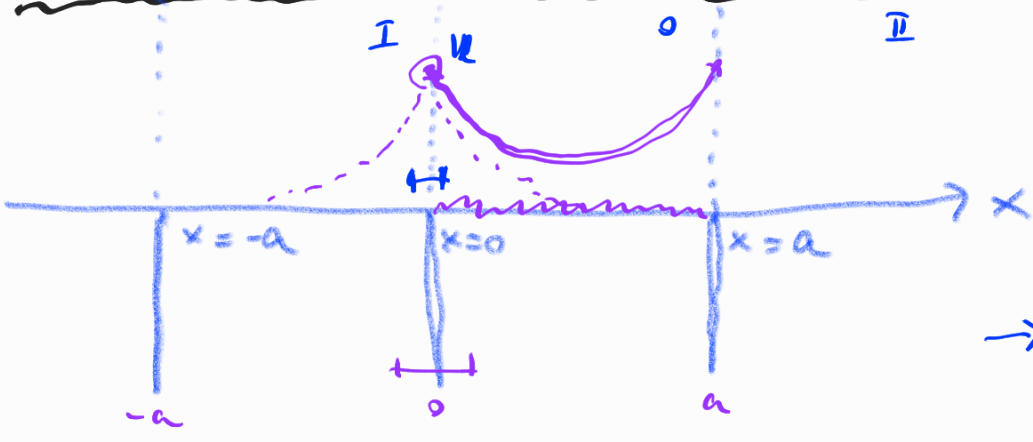


QM II - cv 8

• Dada lanka → δ -krystal



$$H = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} - \lambda \sum_{n \in \mathbb{Z}} \delta_{na}$$

$$H\psi = E\psi$$

Blochův teorem

$$\psi(x) = u_k(x) e^{ikx}$$

$$\rightarrow \boxed{u_k(x) = u_k(x+a)}$$

$$\boxed{k \in \left(-\frac{\pi}{a}, \frac{\pi}{a}\right)}$$

① řešení $x \in (0, a)$ -- $H = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2}$ -- $\psi_0(x) = e^{\kappa x} + A e^{-\kappa x}$

$E < 0$ -- $E = -\frac{\hbar^2 \kappa^2}{2m}$ $E > 0$... $\kappa \rightarrow i\kappa$

② rozšíření na I -- $x \in (-a, 0)$ $\psi(x+a) = u_k(x) e^{i\kappa(x+a)} = e^{i\kappa a} \psi(x)$

$x \in (-a, 0)$ $\psi_I(x) = e^{-i\kappa a} \psi_0(x+a) = e^{-i\kappa a} (e^{\kappa(x+a)} + A e^{-\kappa(x+a)})$

③ kraj. podm. spojitosť ψ v okolí $x=0$

$$\psi(0+) = \psi_0(0) = 1 + A$$

$$\psi(0-) = \psi_I(0) = e^{-i\kappa a} (e^{\kappa a} + A e^{-\kappa a})$$

$$A = \frac{e^{i\kappa a} - e^{-\kappa a}}{e^{-\kappa a} - e^{i\kappa a}}$$

$$\psi_0 = \frac{e^{-\kappa a} - e^{i\kappa a}}{e^{-\kappa a} - e^{i\kappa a}}$$

podmínka na derivaci -- δ -chovalm:

$$\psi(0-) - \psi'(0+) = \lambda \psi_0$$

$$\lambda = \frac{2m\lambda}{\hbar^2}$$

$$\psi_I'(0) - \psi_0'(0) = \kappa e^{-i\kappa a} (e^{\kappa a} - A e^{-\kappa a}) \rightarrow \kappa(1-A) = \lambda \psi_0$$

rovnice pro $\kappa(k)$... $\Rightarrow E_k = -\frac{\hbar^2 \kappa^2}{2m} < 0$... $\left\{ \begin{array}{l} E > 0 \\ \kappa \rightarrow i\kappa \end{array} \right.$

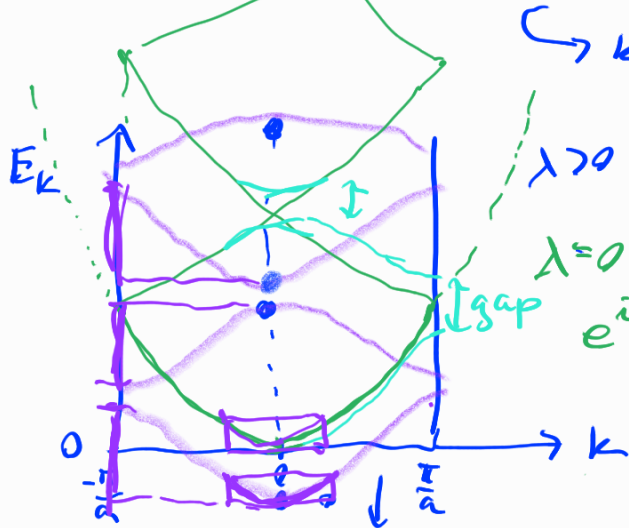
$\cos \kappa a = \cosh \kappa a - \frac{\lambda}{4\kappa} \sinh \kappa a$... pro $E_k = -\frac{\hbar^2 \kappa^2}{2m} < 0$

$\cos \kappa a = \cos \kappa a - \frac{\lambda}{4\kappa} \sin \kappa a$... pro $E_k = +\frac{\hbar^2 \kappa^2}{2m} > 0$

$f(\kappa) = 0$ ↑ řešení $\kappa(k)$

$$\frac{e^{\kappa a} + e^{-\kappa a}}{2} = \cosh \kappa a$$

$\forall k \in (-\frac{\pi}{a}, \frac{\pi}{a}) \rightarrow$ napočítal $f(x)$ $x = 0.01, 0.02, \dots, 10.$



\hookrightarrow kořeny $x_1(k), x_2(k), \dots$

$\lambda = 0$ (volybné častice)
 e^{ikx} $E_k = \frac{\hbar^2 k^2}{2m}$
 $x = k$

dlohouhý limitu
 $k \rightarrow 0$
 n e^{ikx}



Úloha 1 - Simon Païnger

struktura:

$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} = -E_1 \frac{d^2}{dx^2} \dots E_n = E_1 n^2$

	1 častice	rozliš.	$\frac{m_1^2 + m_2^2}{2}$	Bosony	Fermi
	$0 \quad \phi_0 = \frac{1}{\sqrt{2\pi}}$	$\phi_0(\varphi_1)\phi_0(\varphi_2)$	$1 \times$	$1 \times$	$0 \times$
	$E_1 \quad \phi_{\pm 1} = \frac{1}{\sqrt{4\pi}} \cos \varphi, \frac{1}{\sqrt{4\pi}} \sin \varphi = \phi_{1s} \in$ $\phi_{\pm 1} = \frac{1}{\sqrt{2\pi}} e^{im\varphi} \quad m = \pm 1$	$\phi_0(\varphi_1)\phi_{\pm 1}(\varphi_2)$ $(4 \times)$ $\phi_{\pm 1}(\varphi_1)\phi_0(\varphi_2)$	$2 \times$	$2 \times$	$2 \times$
$m=2$	$E_1 4 \quad \frac{1}{\sqrt{4\pi}} \cos/\sin 2\varphi$	$2 E_1 \quad \phi_2(\varphi_1)\phi_2(\varphi_2) \cdot$ $\phi_2(\varphi_1)\phi_0(\varphi_2) \cdot$ $\phi_0(\varphi_1)\phi_2(\varphi_2) \cdot$ $\phi_2(\varphi_1)\phi_0(\varphi_2)$	$4 \times$	$2 \times$	$0 \times$
		$\phi_2(\varphi_1)\phi_2(\varphi_2)$ $\phi_2 \quad \phi_1$	$1 \times$	$1 \times$	$1 \times$

body 1-3

body 4

$H = -\epsilon_1 \left(\frac{\partial^2}{\partial \varphi_1^2} + \frac{\partial^2}{\partial \varphi_2^2} \right) + \lambda \delta(\varphi_1 - \varphi_2)$



$\varphi_1, \varphi_2 \in (0, 2\pi)$
 $\frac{\varphi_1 + \varphi_2}{2} = \phi$
 $\varphi_2 - \varphi_1 = \varphi$

přesné řešení sde uděla

My: poruchové řešení - nutno zvlášť jednotlivé kladiny a Bosony/Fermi

základní stav Bosony: $\Psi_0 = \phi_0(\varphi_1)\phi_0(\varphi_2) = \frac{1}{2\pi}$

$E = E_0^{(0)} + E_0^{(n)}$

$E_0^{(n)} = \langle \Psi_0 | H_I | \Psi_0 \rangle = \frac{\lambda}{(2\pi)^2} \int_0^{2\pi} \int_0^{2\pi} \delta(\varphi_1 - \varphi_2) d\varphi_1 d\varphi_2 = \frac{\lambda}{2\pi}$

$E = 0 + \frac{\lambda}{2\pi}$

excitovaný stav bosonů: $\psi_{\pm} = [\phi_0(\varphi_1) \phi_{\pm}(\varphi_2) + \phi_0(\varphi_2) \phi_{\pm}(\varphi_1)] \frac{1}{\sqrt{2}}$

$E = E^{(0)} + E^{(1)}$ 0. řád: $\psi_{\pm} = \frac{1}{\sqrt{2}} \frac{1}{2\pi} (e^{\pm i\varphi_2} + e^{\pm i\varphi_1})$ $E^{(0)} = \epsilon_1$

matriční prvky $H_I \rightarrow$ $\begin{matrix} \psi_+ \\ \psi_- \end{matrix} \begin{pmatrix} H_{++} & H_{+-} \\ H_{-+} & H_{--} \end{pmatrix}$ H_{mm}
 $m = \pm 1$

$H_{++} = \langle \psi_+ | H_I | \psi_+ \rangle = \frac{\lambda}{2} \frac{1}{(2\pi)^2} \iint (e^{-i\varphi_2} + e^{-i\varphi_1}) \delta(\varphi_1 - \varphi_2) (e^{i\varphi_2} + e^{i\varphi_1}) d\varphi_1 d\varphi_2$
 $= \frac{\lambda}{2} \frac{4}{(2\pi)^2} \int_0^{2\pi} d\varphi_1 = \frac{2\lambda}{2\pi} = \frac{\lambda}{\pi}$ $H_{--} = \frac{\lambda}{\pi}$

$H_{+-} = \frac{\lambda}{2} \frac{1}{(2\pi)^2} \iint (e^{-i\varphi_2} + e^{-i\varphi_1}) \delta(\varphi_1 - \varphi_2) (e^{-i\varphi_2} - e^{-i\varphi_1}) d\varphi_1 d\varphi_2$
 $= \frac{\lambda}{2} \frac{4}{(2\pi)^2} \int_0^{2\pi} e^{-2i\varphi_1} d\varphi_1 = 0$... rotační symetrie



M... moment hybn. integrál p...
 $m_1 + m_2$ $0 + 1$ $0 - 1$ $M \begin{matrix} \nearrow +1 \\ \searrow -1 \end{matrix}$

divš. rot. sym $\hat{R}_d \rightarrow \varphi_1 \rightarrow \varphi_1 + d$
 $\varphi_2 \rightarrow \varphi_2 + d$
 $L = -i\hbar \left(\frac{\partial}{\partial \varphi_1} + \frac{\partial}{\partial \varphi_2} \right)$

$[L, H_I] = [\hat{R}_d, H_I] = 0$

energie $E_{\pm} = E_0 + \frac{\lambda}{\pi}$ (zůstává degenerová v divš. zach. unim. hybn.)