

Markov chains

- in general, it is a sequence of possible events

or random quantities $S^{(k)}$, $k=1, \dots, \infty$

which are chosen from a certain set of states

(configurations) $\{A_i\}$, $i=1, \dots, M$ (could be infinite, but we will consider M to be finite for simplicity)

but occurrences of states in the chain are not

independent, they depend on the previous state:

if in "time" k state A_i appears with probability $\pi_i^{(k)}$
then in "time" $k+1$ state A_j appears with prob. $\pi_j^{(k+1)}$

given by

$$\pi_j^{(k+1)} = \sum_{i=1}^M \pi_i^{(k)} W_{i \rightarrow j}$$

or written as vectors and using the transition matrix

W with elements $W_{i \rightarrow j} = P(A_i \rightarrow A_j) =$ probability
of transition from A_i to A_j

$$\pi^{(k+1)} = \pi^{(k)} \cdot W$$

W must satisfy condition

$$\sum_{j=1}^M W_{i \rightarrow j} = 1 \quad \text{for all } i \quad \left(\begin{array}{l} \text{sum of probabilities} \\ \text{of going from } A_i \text{ to} \\ \text{anywhere} \end{array} \right)$$

- examples in Mathematics

1) Problematic network - two-state system

see Nezbeda, Kolata, Kotrla, page 46

2) Random flipping of a die over an edge

- in both examples above, we can see that after a while the probability π_j of A_j is independent of the previous states

$$\pi^{(k+1)} \approx \pi^{(k)} \text{ for } k \text{ large} \quad \left(\begin{array}{l} \text{"the system"} \\ \text{is losing memory"} \end{array} \right)$$

or

$$\lim_{n \rightarrow \infty} \pi^{(n)} = \pi$$

where π satisfies the equation for left eigenvalues and eigenvectors of the transition matrix

$$\lambda \pi = \pi \cdot W$$

and π is the eigenvector corresponding to $\lambda = 1$:

$$\pi = \pi \cdot W$$

[Note: It can be shown that for transition matrices the other eigenvalues are less than 1 (in absolute value)]

- under which conditions this happens?

It can be shown that (see e.g. W. Feller: An Introduction to Probability Theory and Its Applications) if 1) all states (configurations) can be obtained (reached) from an arbitrary state in a finite time with non-zero probability

2) and no state is periodic

(A_i is periodic, if there is a period m such

that for $\pi_i^{(k)} = 0$ will be $\pi_i^{(k+m)} = 0$

and for $\pi_i^{(k)} \neq 0$ will be $\pi_i^{(k+m)} \neq 0$)

then for arbitrary initial distribution of probabilities

$\pi^{(n)}$ there exists a limit

$$\pi = \lim_{k \rightarrow \infty} \pi^{(k)}$$

and π satisfies the equation

$$\pi = \pi \cdot W$$

Such a system and a set of states are called ergodic.

- in Monte Carlo simulations we do not have the transition matrix, but only the limiting distribution of probabilities give usually as

$$\pi_j = g(A_j) = \frac{e^{-\beta H(A_j)}}{Z}, \quad Z = \sum_j e^{-\beta H(A_j)}$$

but it can be more general $g(A_j)$, of course.

How to construct the transition matrix W ?

- we have 3 conditions:

$$W_{i \rightarrow j} \geq 0 \text{ for all } i, j = 1, \dots, M$$

$$\sum_{j=1}^M W_{i \rightarrow j} = 1 \text{ for all } i = 1, \dots, M \quad (*)$$

$$\pi \cdot W = \pi$$

- the last condition is called the detailed balance and it is a necessary condition to get the limit of stationary probability distribution. It will be valid if we suppose the stronger condition of microscopic reversibility

$$\pi_i W_{i \rightarrow j} = \pi_j W_{j \rightarrow i} \quad (**)$$

$$\text{because } \pi \cdot W \Leftrightarrow \sum_{i=1}^M \pi_i W_{i \rightarrow j} = \pi_j \underbrace{\sum_{i=1}^M W_{j \rightarrow i}}_1 = \pi_j \Leftrightarrow \pi$$

- conditions (*) do not determine W uniquely (we have $2M$ equations for $M \times M$ unknowns)

from (**) we have either

$$\frac{W_{j \rightarrow i}}{W_{i \rightarrow j}} = \frac{\pi_i}{\pi_j} = \frac{g(A_i)}{g(A_j)} \quad \left(\text{it is independent of normalization } (Z)! \right)$$

or $W_{j \rightarrow i} = W_{i \rightarrow j} = 0$ (if $g(A_i) > 0$ for all A_i as in the case when $g(A_i) = e^{-\beta H(A_i)} / Z$)

but still there is many possibilities how to choose elements $W_{i \rightarrow j}$

a) Metropolis method, the most common choice of W

$$W_{i \rightarrow j} = \begin{cases} \alpha_{i \rightarrow j} & \text{for } i \neq j, \pi_j \geq \pi_i \\ \alpha_{i \rightarrow j} \frac{\pi_j}{\pi_i} & \text{for } i \neq j, \pi_j < \pi_i \\ 1 - \sum_{k \neq i} W_{i \rightarrow k} & \text{for } i = j \end{cases}$$

where $\alpha_{i \rightarrow j}$ is an arbitrary symmetric stochastic matrix satisfying $\alpha_{i \rightarrow j} \geq 0$

- the first two conditions of (*) are satisfied thanks to $\sum_{j \neq i} \alpha_{i \rightarrow j} = 1$ (necessary for $W_{i \rightarrow i} \geq 0$)

and moreover we have for $\pi_j \geq \pi_i$

$$\pi_i W_{i \rightarrow j} = \pi_i \alpha_{i \rightarrow j} \stackrel{\substack{\uparrow \\ \text{symmetric } \alpha}}{=} \pi_j \alpha_{j \rightarrow i} \frac{\pi_i}{\pi_j} = \pi_j W_{j \rightarrow i} \quad \checkmark$$

and for $\pi_j < \pi_i$

$$\pi_i W_{i \rightarrow j} = \pi_i \alpha_{i \rightarrow j} \frac{\pi_j}{\pi_i} \stackrel{\downarrow}{=} \pi_j \alpha_{j \rightarrow i} = \pi_j W_{j \rightarrow i} \quad \checkmark$$

b) Barker (Glauber) transition matrix:

$$W_{i \rightarrow j} = \begin{cases} \alpha_{i \rightarrow j} \frac{\pi_j}{\pi_i + \pi_j} & \text{for } i \neq j \\ 1 - \sum_{k \neq i} W_{i \rightarrow k} & \text{for } i = j \end{cases}$$

where again $\alpha_{i \rightarrow j}$ is an arbitrary symmetric stochastic matrix

- now

$$\pi_i W_{i \rightarrow j} = \pi_i \alpha_{i \rightarrow j} \frac{\pi_j}{\pi_i + \pi_j} \stackrel{\downarrow}{=} \pi_j \alpha_{j \rightarrow i} \frac{\pi_i}{\pi_i + \pi_j} = \pi_j W_{j \rightarrow i} \quad \checkmark$$

- there are other methods, e.g. heat-bath method
(see Creutz et al., Phys. Rev. Lett. 42 (1979) 1390)

- how to choose $\alpha_{i \rightarrow j}$?

- that depends on a particular system - we have

- these are probabilities of generating A_j from the configuration A_i and there are different

strategies (e.g. changing just one site in the lattice or moving one molecule randomly etc.)