

Modelling crystal growth

- growth is relatively complicated process, because several different stochastic processes happen simultaneously (capture of a particle on the surface, diffusion of particles on the surface, evaporation of a particle)
- it can proceed differently on different surfaces, it depends if particles are captured from gas or liquid etc.
- it is often difficult to find out how particular processes contribute to the growth, therefore we use simplified stochastic models based on Monte-Carlo simulations and study their behaviour to determine which processes are relevant
 - it turned out that some models behave similarly and they fall to several classes that can be assigned also to more complicated systems
- there are discrete models on lattices and continuous models
 - here we demonstrate the growth process on two rather simple models in one dimension and introduce also the kinetic Monte-Carlo method using this process, but it has much broader use

Discrete models of growth

- particles can be only in the nodes of a lattice
- basic three processes:
 - deposition of new particles
 - evaporation of particles from the surface
 - migration (diffusion) of particles on the surface

- not all these processes are used in every model
- they happen with probability that depends on the current configuration of the surface
- the simplest models are called solid-on-solid (SOS) models that do not consider "overhangs"  or "cavities" 
- particles lie on top of each other and the surface is characterized by the height of columns of particles above the base layer
- to describe the surface quantitatively during simulations we can determine various mean quantities as functions of time as for example (in 1D case)

mean height

$$\bar{h}(t) = \frac{1}{L} \sum_{i=1}^L h(i, t) \quad \text{where } L \text{ is the lattice size}$$

roughness of the surface (variance of h)

$$w(t) = \left(\frac{1}{L} \sum_{i=1}^L [h(i, t) - \bar{h}(t)]^2 \right)^{1/2}$$

or height correlation function

$$G(r, t) = \frac{1}{L} \sum_{i=1}^L [h(i+r, t) - h(i, t)]^2$$

- in the simple simulations when we only add particles

$$\bar{h}(t) \propto t$$

- for other quantities, dependence on time can be $\propto t^\beta$

and for finite systems, there can be a change

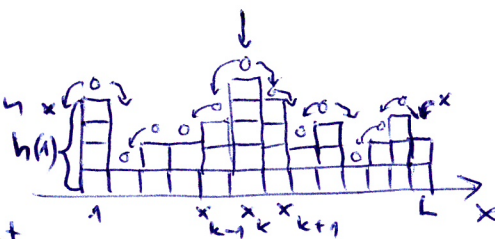
in the exponent after "saturation" of the surface

- quantities can depend also on the lattice size as L^α

Examples: the basic simplest models of growth

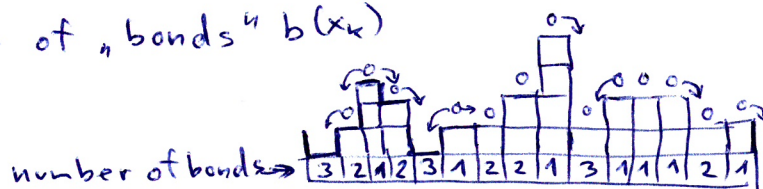
1) Edward-Wilkinson (EW) model

- a simple model of sedimentation
- if a particle (randomly) falls to the site x_k then if the height $h(x_k) \leq h(x_{k-1})$ and $h(x_k) \leq h(x_{k+1})$ then it stays there otherwise it moves (just once) to the neighbouring site with the smallest height or if $h(x_{k-1}) \leq h(x_{k+1}) < h(x_k)$ then randomly to the left or to the right
- periodic boundary conditions



2) Wolf-Villain (WV) model

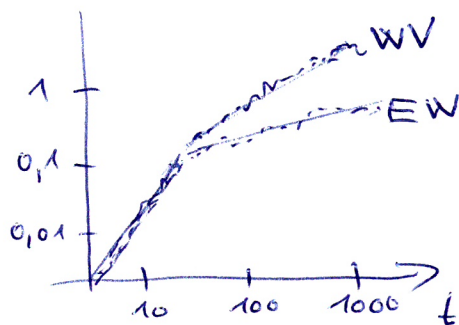
- similar as EW model but the choice of the site where a particle finally deposits is given by the number of "bonds" $b(x_k)$



- the particle goes where $b(x_k)$, $b(x_{k-1})$ or $b(x_{k+1})$ is largest

- they give different exponents for $w(t)$ or $G(t)$

after saturation



Kinetic Monte Carlo - illustration on growth models

- various processes happen with different rates R_α

- for example in growth models we can have

1) different rates for deposition R_1, R_2, R_3

depending on the number of neighbouring particles

(R_i means the number of particles deposited per second to the site with i neighbours)

2) rates of evaporation R_4, R_5, R_6 - also depend on the number of neighbours

3) rates for migration (diffusion) - $R_7, \dots, R_{N_p}$

- depending on neighbourhood / $\nearrow$ total number of possible processes

- all rates R_α determine the model and in principle, these could be calculated from the first principles, but that is usually difficult $\Rightarrow$ estimations or from experiments

- they usually depend on the environment (like temperature etc)

- individual processes are realised by elementary events

(such as deposition to a specific site etc.)

and their number n_α depends on the current configuration A_k

in the k -th iteration of the simulation

(for example if we fill in a hole in A_{k-1} , the neighbour can't migrate there in A_k)

- because R_α are not normalized, to determine probabilities of particular elementary events we have to divide R_α by the total rate

$$R(A_k) = \sum_{\alpha=1}^{N_p} n_\alpha(A_k) R_\alpha$$

- then, a certain elementary event of the process with rate R_α happens with probability

$$\pi_\alpha = \frac{R_\alpha}{R(A_k)}$$

- a possible, but very inefficient algorithm of kinetic MC simulation could be

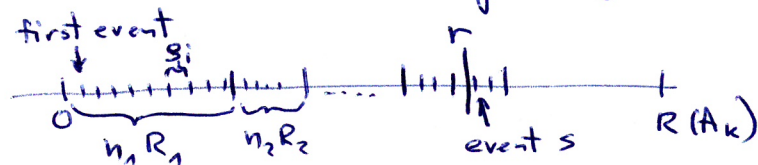
- 1) choose an event that can happen for A_k
- 2) generate $\xi \in (0, 1)$
- 3) if $\xi < \pi_\alpha$ then implement the event and update all necessary variables
otherwise do nothing
- 4) repeat

- the problem is that for many processes and possible events π_α will be too small (or for small R_α also) and we get too many rejections

- n-fold way or BKL algorithm (rejection-free algorithm)

(according to Bortz, Kalos and Lebowitz: J. Comp. Phys. 17 (1975) 10)

- an algorithm that always chooses an event with correct probability
- now we generate a random number $r \in (0, R(A_k))$ and we determine an event corresponding to



- algorithm
 - 1) generate $r \in (0, R(A_k))$
 - 2) find event s using $\sum_{i=1}^s g_i \geq r$, where $g_i = R_i$ for events of process 1, R_2 , etc.
 - 3) implement the event
 - 4) update $R(A_k)$, $n_\alpha(A_k)$ etc.
- in general $O(N)$ → where N is the number of all events
also $O(N)$ → but often faster

- to make it faster, we can first choose a process

using $\sum_{\alpha=1}^{\sigma} n_{\alpha}(A_k) R_{\alpha} \geq r$ where σ is the smallest such index

and then we choose an event for the process

(this can be trivial as often all events for a specific process have the same rate)

- but we still have to update all $n_{\alpha}(A_k)$

- more general algorithm for a lot of processes and events could even group some processes first and choose a group, then a process and finally an event to optimize the number of operations to find an event

Time in kinetic MC simulations

- in standard MC simulations, the number of steps (iterations) does not relate to physical time and generated configurations do not correspond to real evolution of the system

- in kinetic MC, the situation is different because changes of configurations correspond to real processes and events that can happen with corresponding rates

- therefore, it is convenient to introduce time into such simulations using the total rate $R(A_k)$, particularly in each step we increase the time by a random interval given by

$$\Delta t_k = -\frac{1}{R(A_k)} \ln \xi$$

where $\xi \in (0, 1)$

is uniformly distributed random number

and thus, the total time after n steps is

$$t = \sum_{k=1}^n \Delta t_k$$

and mean value of quantity X is

$$\langle X \rangle_{(0,t)} = \frac{1}{t} \sum_{k=1}^n \Delta t_k X(A_k)$$

~ to justify the expression for Δt_k , we can start from an approximative solution of the master equation for Markovian processes for which the change of probability $P(A)$ of finding the system in the configuration A in time t is governed by

$$\frac{\partial P(A,t)}{\partial t} = - \underbrace{\sum_{A' \neq A} W(A \rightarrow A') P(A,t)}_{\text{the change from the state A}} + \underbrace{\sum_{A' \neq A} W(A' \rightarrow A) P(A',t)}_{\text{the change into the state A}}$$

where $W(A \rightarrow A')$ gives probability of transition from A to A' per unit of time, in general dependent on time (it is not the same matrix as in equilibrium simulations)

~ in kinetic MC simulations, we choose $W(A \rightarrow A')$ in such a way to correspond to the correct behaviour of the system in time by considering N_p processes with prescribed rates R_α resulting in $W(A \rightarrow A') = F(R_1, \dots, R_{N_p})$

~ if, in the master equation, the second term is negligible (we consider only changes from the configuration A)

we have

$$\frac{\partial P(A,t)}{\partial t} = -R(A,t) P(A,t)$$

where

$$R(A,t) = \sum_{A' \neq A} W(A \rightarrow A',t) \doteq \sum_{\alpha=1}^{N_p} n_\alpha(A_k) R_\alpha = R(A_k)$$

↑
approximation made in the kinetic MC method

is the total rate of change

- for short times, we can write approximately

$$P(A, t + \Delta t) = \underbrace{e^{-R(A_k)\Delta t}}_{\text{this term expresses the probability that nothing happens during the interval } \Delta t} P(A, t)$$

(the larger is R , the faster decrease towards zero)

- under assumptions that there is one elementary event at a time with the total rate R and that it is a Poisson process (events are independent), the probability that something happens during time t is $1 - e^{-Rt}$. To this cumulative distribution function, the following probability density corresponds

$$g(t) = R e^{-Rt}$$

$$\left(\text{as } \Phi(t) = \int_0^t R e^{-Rt} dt = R \left[-\frac{1}{R} e^{-Rt} \right]_0^t = 1 - e^{-Rt} \right)$$

- therefore, we generate a random time interval after which some event happens in such a way to obtain the corresponding prob. density $g(t)$, using the inverse of the cumulative distribution function $\Phi(t)$ as

$$\Delta t = \Phi^{-1}(v) = -\frac{1}{R} \ln(1-v) \quad \begin{array}{l} \text{where } v \text{ is a random} \\ \text{number from } (0,1) \\ \text{with the uniform distr.} \end{array}$$

- because $1-v$ is also a random variable on $(0,1)$ with a uniform distribution, we can use a simpler formula

$$\Delta t_k = -\frac{1}{R(A_k)} \ln \xi_{(0,1)}$$

for the increment of time during the k -th iteration of kinetic MC simulation