

TEORIE POLE V 0+1 DIM (harmonický oscilátor)

Akce a rovnice pohybu

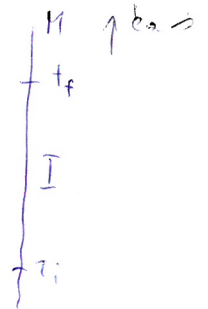
$f: M \rightarrow \mathbb{R}$ vyvoj "pole" v čase

akce

$$S[I](f) = + \frac{1}{2} \int_I (m \dot{f}^2 - k f^2) dt - \int_I f j dt$$

$$= - \frac{1}{2} f \cdot \mathbb{F}[I] \cdot f + f \cdot (dt j)$$

$$\mathbb{F}[I] = - \frac{d}{dt} \cdot (X[I] m dt \delta) \cdot \frac{d}{dt} + (X[I] k dt \delta)$$



$$m = 0$$

pohybové rovnice

$$\vec{\mathbb{F}} = m dt \frac{d^2}{dt^2} + k dt \delta$$

$$\vec{\mathbb{F}} \cdot f = j dt \iff \left[+ \frac{d^2}{dt^2} + \omega^2 \right] f = \frac{j}{m} \quad \omega^2 = \frac{k}{m}$$

0+1 rozstěpení

$$q(t) = q[t] \cdot f = f(t) \quad \text{hodnota v čase } t$$

$$p(t) = p[t] \cdot f = m \dot{f}(t) \quad \text{hyb-ost v čase } t$$

$$L(q, \dot{q}) = \frac{1}{2} m \dot{q}^2 - \frac{1}{2} k q^2 + j q$$

$$H(q, p) = \frac{1}{2m} p^2 + \frac{1}{2} k q^2 - j q$$

obrazové členy

$$\vec{\partial} \mathbb{F}[t] = m dt \delta_t \frac{d}{dt} = q[t] p[t] \quad f_1 \cdot \vec{\partial} \mathbb{F}[t] \cdot f_2 = m f_1(t) \dot{f}_2(t)$$

$$\partial \mathbb{F}[t] = p[t] q[t] - q[t] p[t] \quad f_1 \cdot \partial \mathbb{F}[t] \cdot f_2 = m (\dot{f}_1 f_2 - f_1 \dot{f}_2) |_t$$

Prostor řešení

nehomog. řešení

$$\bar{f} = q_0 \cos(\omega(t-t_0)) - \frac{p_0}{m\omega} \sin(\omega(t-t_0)) + \frac{f}{k}$$

homogem řešení

$$\bar{f} = q_0 \cos(\omega(t-t_0)) - \frac{p_0}{m\omega} \sin(\omega(t-t_0))$$

$$\bar{f} = f_0(t_0)(q_0, p_0)$$

symplekt. form =

$$f_1 \cdot \partial f_2 / \partial t_1 - f_2 \cdot \partial f_1 / \partial t_2 = p_1 q_2 - q_1 p_2$$

$$\{q_1(t_1), p_1(t_1)\} = -1$$

$$\{f(t_2), f(t_1)\} = G_0(t_2, t_1)$$

Pr:

$$f_i = q_i \cos(\omega(t-\tau_i)) + \frac{p_i}{m\omega} \sin(\omega(t-\tau_i))$$

$$f_1 \cdot \mathcal{D}[t_0] \cdot f_2 = m (\dot{f}_1 f_2 - f_1 \dot{f}_2) \Big|_{t_0} =$$

$$= (-mq_1\omega \sin(\omega(t_0-\tau_1)) + p_1 \cos(\omega(t_0-\tau_1))) \\ (q_2 \cos(\omega(t_0-\tau_2)) + \frac{p_2}{m\omega} \sin(\omega(t_0-\tau_2)))$$

- $1 \leftrightarrow 2$

$$= -m\omega q_1 q_2 (\sin(\omega(t_0-\tau_1)) \cos(\omega(t_0-\tau_2)) - \cos(\omega(t_0-\tau_1)) \sin(\omega(t_0-\tau_2)))$$

$$+ \frac{p_1 p_2}{m\omega} (\cos(\omega(t_0-\tau_1)) \sin(\omega(t_0-\tau_2)) - \sin(\omega(t_0-\tau_1)) \cos(\omega(t_0-\tau_2)))$$

$$+ p_1 q_2 (\cos(\omega(t_0-\tau_1)) \cos(\omega(t_0-\tau_2)) + \sin(\omega(t_0-\tau_1)) \sin(\omega(t_0-\tau_2)))$$

$$- q_1 p_2 (\sin(\omega(t_0-\tau_1)) \sin(\omega(t_0-\tau_2)) + \cos(\omega(t_0-\tau_1)) \cos(\omega(t_0-\tau_2)))$$

$$= m\omega q_1 q_2 \sin(\omega(\tau_1-\tau_2)) + \frac{p_1 p_2}{m\omega} \sin(\omega(\tau_1-\tau_2))$$

$$+ (p_1 q_2 - q_1 p_2) \cos(\omega(\tau_1-\tau_2))$$

Greenovy funkce

$$\omega = \text{konst} \quad (\text{tj. } m, k = \text{konst})$$

$$G_{\text{ret}}^{t_2, t_1} = \frac{\sin(\omega(t_2 - t_1))}{m\omega} \Theta(t_2 - t_1)$$

$$G_{\text{adv}}^{t_2, t_1} = -\frac{\sin(\omega(t_2 - t_1))}{m\omega} \Theta(t_2 - t_1)$$

$$G_{\text{sym}}^{t_2, t_1} = \frac{\sin(\omega(t_2 - t_1))}{m\omega} \text{sign}(t_2 - t_1)$$

$$G_c^{t_2, t_1} = \frac{\sin(\omega(t_2 - t_1))}{m\omega}$$

Candyllo propagátor

$$D_c[t] = -G_c \cdot \partial F[t]$$

$$\begin{aligned} (D_c[t_0] \cdot f)(t) &= -m \frac{\partial G_c(t|t_0)}{\partial t_0} f(t_0) + m G_c(t|t_0) \frac{df}{dt}(t_0) \\ &= \cos(\omega(t-t_0)) f(t_0) + \frac{\sin(\omega(t-t_0))}{m\omega} m \dot{f}(t_0) \end{aligned}$$

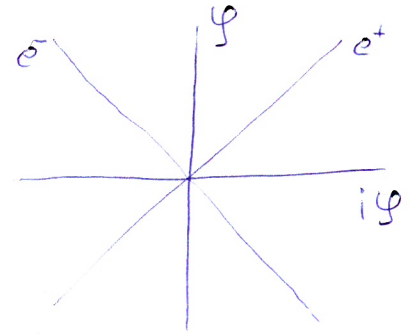
$$\Phi_c[t_0](q_0, p_0) = \cos(\omega(t-t_0)) q_0 + \frac{\sin(\omega(t-t_0))}{m\omega} p_0$$

Poisson. závody

$$\{f(t_2), f(t_1)\} = \frac{\sin(\omega(t_2 - t_1))}{m\omega}$$

Frekvencní rozštěpení

komplexní řešení - prostor \mathcal{G}^c



$$e^+(t) = \frac{1}{\sqrt{2\omega m}} \exp(-i\omega t)$$

$$\omega > 0$$

$$e^-(t) = \frac{1}{\sqrt{2\omega m}} \exp(i\omega t)$$

$$e^{\mp} \cdot \partial \mathcal{F} \cdot e^{\pm} = \frac{1}{2\omega m} m (\dot{e}^- e^+ - e^- \dot{e}^+) \Big|_t = \pm i$$

$$e^{\pm} \cdot \partial \mathcal{F} \cdot e^{\pm} = 0$$

umožňuje defini-ovat "normu" na pos. def. řešeních

$$\langle \phi_1, \phi_2 \rangle_+ = -i \phi_1^* \cdot \partial \mathcal{F} \cdot \phi_2 \quad \phi_1, \phi_2 \sim e^+$$

v 0+1 dim "triviální" řešení

\langle , \rangle_+ je nedeg. i pos. def.

báze v \mathcal{G}^c

e^+, e^- tvoří bázi v \mathcal{G}^c

$$f = a e^+ + a^* e^- \quad \text{řešení} \in \mathcal{G}$$

$$a = -i e^- \cdot \partial \mathcal{F} \cdot f = \frac{1}{2} \left(\sqrt{m\omega} q + \frac{i}{m\omega} p \right)$$

mezivlnová +

$$a^* = i e^+ \cdot \partial \mathcal{F} \cdot f = \frac{1}{2} \left(\sqrt{m\omega} q - \frac{i}{m\omega} p \right)$$

$$\{a, a^*\} = -e^- \cdot \partial \mathcal{F} \cdot e^+ \cdot \partial \mathcal{F} \cdot e^- = i$$

Kvantování

$$\begin{aligned}
 q, p &\rightarrow \hat{q}, \hat{p} \\
 \{q, p\} = -1 &\rightarrow \frac{i}{\hbar} [\hat{q}, \hat{p}] = -\hat{1} & [\hat{q}, \hat{p}] = i\hbar \hat{1} \\
 f &\rightarrow \hat{f} \\
 \{f^{i_1}, f^{i_2}\} = G_0^{i_1 i_2} &\rightarrow \frac{i}{\hbar} [\hat{f}^{i_1}, \hat{f}^{i_2}] = G_0^{i_1 i_2} \hat{1} & [\hat{f}, \hat{f}] = i\hbar G_0 \hat{1} \\
 \frac{d^2}{dt^2} f + \omega^2 f = 0 &\rightarrow \frac{d^2}{dt^2} \hat{f} + \omega^2 \hat{f} = 0
 \end{aligned}$$

kreací a anihilací oper.

$$\hat{f}(t) = \hat{a} e^{+it} + \hat{a}^{\dagger} e^{-it}$$

$$\hat{a} = -i e^{-it} \cdot \partial_t \hat{f} \cdot \hat{f} = \frac{1}{\sqrt{2}} \left(\sqrt{m\omega} \hat{q} + \frac{i}{\sqrt{m\omega}} \hat{p} \right)$$

$$\hat{a}^{\dagger} = i e^{+it} \cdot \partial_t \hat{f} \cdot \hat{f} = \frac{1}{\sqrt{2}} \left(\sqrt{m\omega} \hat{q} - \frac{i}{\sqrt{m\omega}} \hat{p} \right)$$

$$[\hat{a}, \hat{a}^{\dagger}] = -e^{-it} \cdot \partial_t \cdot [\hat{f}, \hat{f}] \cdot \partial_t \cdot e^{+it} = -i\hbar e^{-it} \cdot \partial_t \cdot e^{+it} = \hbar \hat{1}$$

Hamiltonian

$$\hat{H} = \frac{\omega}{2} \left(\frac{\hat{p}^2}{m\omega} + m\omega \hat{q}^2 \right) = \frac{\omega}{2} (\hat{a} \hat{a}^{\dagger} + \hat{a}^{\dagger} \hat{a})$$

Fockovské báze v $0+1$

$$[\hat{a}, \hat{a}^\dagger] = \hat{1}$$

$$\hbar = 1$$

$$\hat{H} = \hat{a}^\dagger \hat{a} = \frac{1}{2} \left(\frac{\hat{p}^2}{m\omega} + m\omega \hat{q}^2 - \hat{1} \right)$$

$$\hat{H}^\dagger = \hat{H}$$

hermit. \Rightarrow def.

vlastní vekt \vec{n}

$$\hat{H} |m\rangle = m |m\rangle$$

$$\hat{H} \hat{a} |m\rangle = (m-1) \hat{a} |m\rangle$$

$$\langle m | \hat{a}^\dagger \hat{a} |m\rangle = m \langle m | m \rangle$$

\Rightarrow m kladné

$$\hat{H} \hat{a}^\dagger |m\rangle = (m+1) \hat{a}^\dagger |m\rangle$$

$$|m-1\rangle = \frac{1}{\sqrt{m}} \hat{a} |m\rangle$$

$$|m+1\rangle = \frac{1}{\sqrt{m+1}} \hat{a}^\dagger |m\rangle$$

pozitivita \Rightarrow

$$\exists |vac\rangle \quad \hat{a} |vac\rangle = 0$$

$$|m\rangle = \frac{1}{\sqrt{m!}} \hat{a}^{\dagger m} |vac\rangle$$

$$|vac\rangle = |0\rangle$$

báze stavů

$$|m\rangle \quad m \in \mathbb{N}_0$$

$$|st\rangle = \sum_m \psi_m |m\rangle$$