

Koharentní stavy

máme zvolenou část: rep.

koh. stav

$$|coh: \phi\rangle = \exp\left(\frac{1}{2}\langle\phi, \phi\rangle\right) \hat{W}[\phi] |vac\rangle$$

shift operator

$$\hat{W}[\phi] = \exp(i\phi \cdot \partial \cdot \hat{\Phi}) = \exp(i\hat{L}_\phi)$$

$$\hat{W}[\phi]^\dagger \hat{\Phi} \hat{W}[\phi] = \hat{\Phi} + \phi \hat{1}$$

důkaz:

$$\begin{aligned} [\hat{\Phi}, i\phi \cdot \partial \cdot \hat{\Phi}] &= -i[\hat{\Phi}, \hat{\Phi}] \cdot \partial \cdot \phi = \phi \Rightarrow [\hat{\Phi}, F(i\phi \cdot \partial \cdot \hat{\Phi})] = \phi F'(i\phi \cdot \partial \cdot \hat{\Phi}) \\ \hat{W}[\phi]^\dagger \hat{\Phi} \hat{W}[\phi] &= \hat{\Phi} + \hat{W}[\phi]^\dagger [\hat{\Phi}, \exp(i\phi \cdot \partial \cdot \hat{\Phi})] = \hat{\Phi} + \phi \hat{W}[\phi]^\dagger \hat{W}[\phi] \\ &= \hat{\Phi} + \phi \hat{1} \end{aligned}$$

koharentní stav je tak "posunutá" vakuum
ale! neortonormální, nenormalizované!

pomocné Lemma

$$\begin{aligned} \hat{W}[\phi_1] \hat{W}[\phi_2] &= \exp(i\phi_1 \cdot \partial \cdot \phi_2) \hat{W}[\phi_2] \hat{W}[\phi_1] \\ \exp(\hat{a}[\phi_1]) \exp(\hat{a}[\phi_2]^\dagger) &= \exp(\langle\phi_1, \phi_2\rangle) \exp(\hat{a}[\phi_2]^\dagger) \exp(\hat{a}[\phi_1]) \\ &= \exp\left(\frac{1}{2}\langle\phi_1, \phi_2\rangle\right) \exp(\hat{a}[\phi_1] + \hat{a}[\phi_2]^\dagger) \end{aligned}$$

důkaz:

$$\begin{aligned} \text{důsledky } \exp(\hat{A} + \hat{B}) &= \exp\left(-\frac{1}{2}[\hat{A}, \hat{B}]\right) \exp \hat{A} \exp \hat{B} \\ \text{pro } [\hat{A}, [\hat{A}, \hat{B}]] &= [\hat{B}, [\hat{A}, \hat{B}]] = 0 \end{aligned}$$

normální uspořádaný shift oper.

$$\begin{aligned} \hat{W}[\phi] &= \exp(\hat{a}[\phi]^\dagger - \hat{a}[\phi]) = \\ &= \exp\left(-\frac{1}{2}\langle\phi, \phi\rangle\right) \exp(\hat{a}[\phi]^\dagger) \exp(-\hat{a}[\phi]) \end{aligned}$$

vztah koherentních a částicových stavů

$$|coh \phi\rangle = \exp(\hat{a}(\phi)^\dagger) |vac\rangle$$

$$= \sum_n \frac{1}{n!} |\phi: n\rangle$$

- superpozice částic v módu ϕ
- všechny počty částic (včetně nulové)

ortonormalita a normovanost

$$\langle coh \phi_1 | coh \phi_2 \rangle = \exp(\langle \phi_1 | \phi_2 \rangle)$$

$$\exp(\hat{a}(\phi_1)) \exp(\hat{a}(\phi_2)^\dagger) = \exp(\langle \phi_1, \phi_2 \rangle) \exp(\hat{a}(\phi_2)^\dagger) \exp(\hat{a}(\phi_1))$$

- není kolmí: pro kolmé módy $\langle \phi_1, \phi_2 \rangle = 0$
(obsahuje vakuum)

- není normované: pro norm. módy
 $\langle coh \phi | coh \phi \rangle = \exp(\langle \phi | \phi \rangle)$

relace úplnosti

$$\hat{1} = \int_{\phi \in \mathcal{H}} d\mu(\phi) |coh \phi\rangle \langle coh \phi|$$

$$d\mu(\phi) = \exp(-\langle \phi, \phi \rangle) d\Gamma$$

$$d\Gamma = \left(\text{Det} \frac{\partial f}{\partial \bar{u}} \right)^{\frac{1}{2}} = \left(\text{Det} \frac{\partial \bar{f}}{\partial u} \right)^{\frac{1}{2}} = \prod_x \frac{d\varphi^x d\bar{u}^x}{2\pi}$$

$\uparrow$
 $\det J = 1$

viz dále

coherent state is a particle basis:

$u = \{u_k\}$ norm. base in $\mathcal{H}$

$|u: m\rangle$ particle basis

$$|\text{coh } \phi\rangle = \sum_m \frac{1}{\sqrt{m!}} \langle u, \phi \rangle^m |u: m\rangle$$

$\uparrow$
 $\prod_k \langle u_k, \phi \rangle^{m_k}$

directly:

$$\begin{aligned} |\text{coh } \phi\rangle &= \exp(\langle \hat{\Phi}, \phi \rangle) |vac\rangle = \\ &\quad \uparrow \hat{a}[\phi]^\dagger \\ &= \exp\left(\sum_k \langle \hat{\Phi}, u_k \rangle \langle u_k, \phi \rangle\right) |vac\rangle \\ &= \prod_k \exp(\langle u_k, \phi \rangle \hat{a}_k^\dagger) |vac\rangle \\ &= \prod_k \sum_{m_k} \frac{1}{m_k!} \langle u_k, \phi \rangle^{m_k} \hat{a}_k^{+m_k} |vac\rangle \\ &= \sum_m \frac{1}{\sqrt{m!}} \langle u, \phi \rangle^m |u: m\rangle \\ &\quad \uparrow \prod_k \frac{1}{m_k!} \langle u_k, \phi \rangle^{m_k} \quad \uparrow \frac{1}{\sqrt{m!}} \hat{a}[u]^{+m} |vac\rangle = \prod_k \frac{1}{m_k!} \hat{a}_k^{+m_k} |vac\rangle \end{aligned}$$

lastní vlastnosti anihil. operátoru

$$\hat{a}|\phi_0\rangle|\text{coh}:\phi\rangle = \langle\phi_0,\phi\rangle|\text{coh}:\phi\rangle$$

duhaz

$$\begin{aligned}\langle\phi_0,\hat{\Phi}\rangle \exp(\frac{1}{2}\langle\phi,\phi\rangle) \hat{W}|\phi\rangle|vac\rangle &= \\ &= \exp(\frac{1}{2}\langle\phi,\phi\rangle) \hat{W}|\phi\rangle \langle\phi_0,\hat{\Phi}+\phi\hat{1}\rangle|vac\rangle \\ &= \langle\phi_0,\phi\rangle \exp(\frac{1}{2}\langle\phi,\phi\rangle) \hat{W}|\phi\rangle|vac\rangle \\ &= \langle\phi_0,\phi\rangle|\text{coh}:\phi\rangle\end{aligned}$$

neboli $|\text{coh}:\phi\rangle$ je vl. vekt. meg. pro částici $\hat{\Phi}^-$

$$\hat{\Phi}^-|\text{coh}:\phi\rangle = \phi^-|\text{coh}:\phi\rangle$$

ale se breacímí oper.

$$\hat{a}[\phi_0]^+|\text{coh}:\phi\rangle = \phi_0 \cdot \frac{\delta}{\delta\phi}|\text{coh}:\phi\rangle$$

duhaz:

$$\begin{aligned}\hat{a}[\phi_0]^+|\text{coh}:\phi\rangle &= \hat{a}[\phi_0]^+ \exp(\hat{a}[\phi]^+) |vac\rangle = \\ &= -i \hat{\Phi}^- \cdot \partial\hat{\Phi} \cdot \phi_0 \exp(-i \hat{\Phi}^- \cdot \partial\hat{\Phi} \cdot \phi) |vac\rangle = \\ &= \phi_0 \cdot \frac{\delta}{\delta\phi} \exp(-i \hat{\Phi}^- \cdot \partial\hat{\Phi} \cdot \phi) |vac\rangle = \phi_0 \cdot \frac{\delta}{\delta\phi} |\text{coh}:\phi\rangle\end{aligned}$$

střední počet částic

$$\frac{\langle\text{coh}:\phi|\hat{N}|\phi_0\rangle|\text{coh}:\phi\rangle}{\langle\text{coh}:\phi|\text{coh}:\phi\rangle} = \frac{\langle\phi,\phi_0\rangle\langle\phi_0,\phi\rangle}{\langle\phi_0,\phi_0\rangle} \quad \text{- ortogonální módy nepřeslávají!}$$

$$\text{tj. } \frac{\langle\text{coh}:\phi|\hat{N}|\phi\rangle|\text{coh}:\phi\rangle}{\langle\text{coh}:\phi|\text{coh}:\phi\rangle} = \langle\phi,\phi\rangle$$

duhaz

$$N|\phi_0\rangle = \frac{\hat{a}[\phi_0]^+ \hat{a}[\phi_0]}{\langle\phi_0,\phi_0\rangle} \hat{a}[\phi_0]|\text{coh}:\phi\rangle = \langle\phi_0,\phi\rangle|\text{coh}:\phi\rangle$$

celkový počet částic

$$\frac{\langle\text{coh}:\phi|\hat{N}|\text{coh}:\phi\rangle}{\langle\text{coh}:\phi|\text{coh}:\phi\rangle} = \langle\phi,\phi\rangle \quad \text{- všechny částice v módu } \phi$$

střední hodnoty pole

$$\frac{\langle \text{coh } \phi | \hat{\Phi} | \text{coh } \phi \rangle}{\langle \text{coh } \phi | \text{coh } \phi \rangle} = \phi$$

důkaz:

$$\frac{\langle \text{coh } \phi | \hat{\Phi} | \text{coh } \phi \rangle}{\langle \text{coh } \phi | \text{coh } \phi \rangle} = \frac{\langle \text{coh } | \hat{\Phi}^- + \hat{\Phi}^+ | \text{coh } \phi \rangle}{\langle \text{coh } \phi | \text{coh } \phi \rangle} = \phi^- + \phi^+ = \phi$$

$$\text{díky } \hat{\Phi}^+ | \text{coh } \phi \rangle = \phi | \text{coh } \phi \rangle \quad \langle \text{coh } \phi | \hat{\Phi}^+ = \phi^+ \langle \text{coh } \phi |$$

komplementarita a čast. jízna

$$\frac{\langle \phi: m | \hat{\Phi} | \phi: n \rangle}{\langle \phi: m | \phi: n \rangle} = 0$$

či

$$\langle u: m | \hat{\Phi} | u: m \rangle = 0$$

asi zgodat

PII-7b

gledam mod ϕ_0

$$\langle \phi_0, \phi_0 \rangle = 1$$

$$\hat{A} = \hat{A}[\phi]$$

$$[\hat{A}, \hat{A}^\dagger] = \hat{1}$$

$$\hat{N} = \hat{A}^\dagger \hat{A}$$

čisti stanj

$$|n\rangle = |\phi_0: n\rangle = \frac{1}{\sqrt{n!}} \hat{A}^{\dagger n} |vac\rangle$$

$$\langle n | n \rangle = 1$$

koher. stanje

$$\phi = r \phi_0$$

$$\langle \phi, \phi \rangle = |r|^2$$

$$|coh \phi\rangle = \exp\left(\frac{1}{2} |r|^2\right) \hat{W}[\phi] |vac\rangle$$

$r \phi_0$

$$= \exp(\hat{A}[\phi]^\dagger) |vac\rangle =$$

$$= \exp(r \hat{A}^\dagger) |vac\rangle$$

$$= \sum_n \frac{1}{\sqrt{n!}} r^n |n\rangle$$

$$\langle coh \phi | coh \phi \rangle = \exp(|r|^2)$$

$$|\tilde{coh} \phi\rangle = \hat{W}[\phi] |vac\rangle =$$

$$= \exp\left(-\frac{1}{2} |r|^2\right) \sum_n \frac{1}{\sqrt{n!}} r^n |n\rangle$$

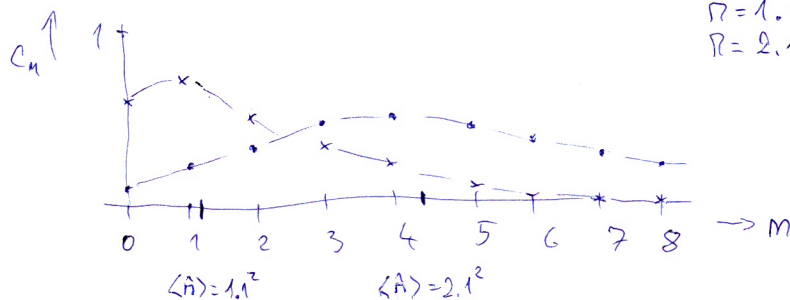
$$\langle \tilde{coh} \phi | \tilde{coh} \phi \rangle = 1$$

$$\langle \tilde{coh} \phi | \hat{N} | \tilde{coh} \phi \rangle = \frac{\langle coh \phi | \hat{N} | coh \phi \rangle}{\langle coh \phi | coh \phi \rangle} = \langle \phi, \phi \rangle = |r|^2$$

obzorica koefic. pro $|\tilde{coh} r \phi_0\rangle = \sum c_n |n\rangle$

$$r = 1.1 \quad \times \times \times$$

$$r = 2.1 \quad \dots$$



Holomorfní reprezentace

prehled pojmů

$f: \mathbb{C} \rightarrow \mathbb{C}$ holomorfní $f(z)$

- existuje komplexní derivace podle z
- f nezávisí na z^*
- f lze napsat jako (meln.) řadu v z

$F: V \rightarrow \mathbb{C}$ holomorfní na komplex. vekt. pr. V

$\equiv F(\phi + z\psi)$ holomorfní v $z \quad \forall \phi, \psi$

polud $V = V_{\mathbb{R}}^{\mathbb{C}} = V_{\mathbb{R}} \oplus iV_{\mathbb{R}}$, pak F holomorfní.

z -mena: $\Rightarrow$ F závisí na ϕ a ne na ϕ^*

$F: V_{\mathbb{R}} \rightarrow \mathbb{C}$ analytická na reál. vekt. pr. $V_{\mathbb{R}}$

$\equiv$ existuje holomorfní rozšíření na $V = V_{\mathbb{R}}^{\mathbb{C}}$
(ne nutně globálně - vždy lok.)

$F: \mathcal{G} \rightarrow \mathbb{C}$ J-holomorfní

$\equiv F(\phi + z\psi)$ holomorfní v $z \quad \forall \phi, \psi$

Plati:

F analytická na $\mathcal{G}$

F je J-holomorfní, pokud $F(\phi) = F(\phi^*)$

F je J-antiholomorfní, pokud $F(\phi) = F(\phi^*)$

zde $F(\phi)$ je holom. rozš. F na $\mathcal{G}^{\mathbb{C}}$

důkaz

$$(z\psi\phi)^+ = z\psi^+\phi^+ \quad (z\psi\phi)^- = z^*\psi^-\phi^- \Rightarrow z\psi\phi = z\psi^+\phi^+ + z^*\psi^-\phi^-$$

$$F(\phi) = F(z\psi^+\phi^+ + z^*\psi^-\phi^-) \Rightarrow \text{závislost na } z/z^* = \text{závislost na } \phi^+/\phi^-$$

holomorfnost Sol. st.

$|\text{coh } \phi\rangle$ je J-holomorfní

$\langle \text{coh } \phi|$ je J-antiholomorfní

$$|\text{coh } \phi\rangle = \exp(\hat{a}^\dagger | \phi \rangle) | \text{vac} \rangle = \exp(-i \hat{\Phi} \cdot \omega \cdot \phi^+) | \text{vac} \rangle$$

holomorf. repr.

$$F[|stau\rangle](\phi) = \langle coh \phi | stau \rangle$$

J-antibal.

$$F[\langle stau|](\phi) = \langle stau | coh \phi \rangle$$

J-bal.

$$F[\hat{A}](\phi_1, \phi_2) = \langle coh \phi_1 | \hat{A} | coh \phi_2 \rangle$$

J-antibal/J-bal.

větná repr. - množina informací
(viz normální uspoř. děle)

vlastnosti:

$$F[\hat{a}[\phi]^\dagger |stau\rangle](\phi_1) = \langle \phi_1, \phi \rangle F[|stau\rangle](\phi_1)$$

$$F[\hat{a}[\phi] |stau\rangle](\phi_1) = \phi \cdot \frac{\delta}{\delta \phi} F[|stau\rangle](\phi_1)$$

indefinite

$$F[|vac\rangle](\phi_1) = 1$$

$$F[\hat{a}[\phi]^\dagger{}^n |vac\rangle](\phi_1) = \langle \phi_1, \phi \rangle^n$$

$$F[|coh: \phi\rangle](\phi_1) = \exp(\langle \phi_1, \phi \rangle)$$

$$F[|vac\rangle \langle vac|](\phi_1, \phi_2) = 1$$

$$F[\hat{1}](\phi_1, \phi_2) = \exp(\langle \phi_1, \phi_2 \rangle)$$

$$F[\hat{A}[\phi]](\phi_1, \phi_2) = \frac{\langle \phi_1, \phi \rangle \langle \phi, \phi_2 \rangle}{\langle \phi, \phi \rangle} \exp(\langle \phi_1, \phi_2 \rangle)$$

$$F[\hat{N}](\phi_1, \phi_2) = \langle \phi_1, \phi_2 \rangle \exp(\langle \phi_1, \phi_2 \rangle)$$

$$F[|coh \phi\rangle \langle coh \phi|](\phi_1, \phi_2) = \exp(\langle \phi_1, \phi \rangle + \langle \phi, \phi_2 \rangle)$$

$$F[\hat{W}[\phi]](\phi_1, \phi_2) = \exp(-\frac{1}{2} \langle \phi, \phi \rangle + \langle \phi_1, \phi \rangle - \langle \phi, \phi_2 \rangle + \langle \phi_1, \phi_2 \rangle)$$

částečné amplitudy v holomorfní repr.

$$\langle u: m | \text{stav} \rangle = \frac{1}{\sqrt{m!}} (u \cdot \delta)^m F[| \text{stav} \rangle](0)$$

$$\uparrow \prod_k (u_k \cdot \delta)^{m_k}$$

dužas:

$$\langle \text{coh } \phi | \hat{a}[\phi_0] | \text{stav} \rangle = \phi_0 \cdot \frac{\delta}{\delta \phi} \langle \text{coh} : \phi | \text{stav} \rangle$$

$$\langle \text{coh } \phi | \hat{a}[u]^m | \text{stav} \rangle = (u \cdot \delta)^m \langle \text{coh} : \phi | \text{stav} \rangle$$

$$\begin{aligned} \langle u: m | \text{stav} \rangle &= \frac{1}{\sqrt{m!}} \langle \text{vac} | \hat{a}[u]^m | \text{stav} \rangle = \frac{1}{\sqrt{m!}} (u \cdot \delta)^m \langle \text{coh} : \phi | \text{stav} \rangle \Big|_{\phi=0} \\ &= \frac{1}{\sqrt{m!}} (u \cdot \delta)^m F[| \text{stav} \rangle](0) \end{aligned}$$

Rekonstrukce stavů a operátorů z holom. repr.

Relace úplnosti

$$\hat{1} = \int_{\phi} d\mu |\text{coh } \phi\rangle \langle \text{coh } \phi|$$

$$d\mu = \exp(-\langle \phi, \phi \rangle) d\Gamma(\phi) = \exp(-\frac{1}{2} \phi \cdot \partial \bar{\partial} \cdot \phi) d\Gamma(\phi)$$

$$d\Gamma = \left(\text{Det} \frac{\partial \bar{\partial}}{2\pi} \right)^{\frac{1}{2}} = \left(\text{Det} \frac{\partial \bar{\partial}}{2\pi} \right)^{\frac{1}{2}}$$

⇓

$$|\text{stav}\rangle = \int d\mu(\phi) |\text{coh } \phi\rangle \langle \text{coh } \phi | \text{stav}\rangle = \int d\mu(\phi) |\text{coh } \phi\rangle F(|\text{stav}\rangle)(\phi)$$

$$\begin{aligned} \hat{A} &= \int d\mu(\phi_1) d\mu(\phi_2) |\text{coh } \phi_1\rangle \langle \text{coh } \phi_1 | \hat{A} | \text{coh } \phi_2\rangle \langle \text{coh } \phi_2 | = \\ &= \int d\mu(\phi_1) d\mu(\phi_2) |\text{coh } \phi_1\rangle F[\hat{A}](\phi_1, \phi_2) \langle \text{coh } \phi_2| \end{aligned}$$

relace dh best:

$$F \left[\int_{\phi} dh \, | \text{coh } \phi \rangle \langle \text{coh } \phi | \right] (\phi_1, \phi_2) =$$

$$= \int_{\phi} d\Gamma \exp(-\langle \phi, \phi \rangle + \langle \phi_1, \phi \rangle + \langle \phi, \phi_2 \rangle)$$

$$= \int_{\phi} \left(\text{Det} \frac{\partial \bar{F}}{\partial u} \right)^{\frac{1}{2}} \exp \left(-\frac{1}{2} \phi \cdot \partial \bar{F} \cdot \phi + \phi \cdot \partial \bar{F} \cdot (\phi_1^- + \phi_2^+) \right) =$$

$$= \exp \left(\frac{1}{2} (\phi_1^- + \phi_2^+) \cdot \partial \bar{F} \cdot (\phi_1^- + \phi_2^+) \right)$$

$$= \exp(\langle \phi_1, \phi_2 \rangle) =$$

$$= F[\mathbb{1}] (\phi_1, \phi_2)$$

$$\langle \phi, \psi \rangle = \phi \cdot \frac{1}{2} (\partial \bar{F} - i \partial \bar{F}) \cdot \psi = \phi^- \cdot \partial \bar{F} \cdot \psi^+$$

$$dh(\phi) = \exp(-\langle \phi, \phi \rangle) d\Gamma(\phi)$$

$$d\Gamma(\phi) = \left(\text{Det} \frac{\partial \bar{F}}{\partial u} \right)^{\frac{1}{2}} = \left(\text{Det} \frac{\partial \bar{F}}{\partial u} \right)^{\frac{1}{2}} = \prod_x \frac{d\psi^x d\pi_x}{2\pi}$$

$\uparrow$
[det] = 1

$$dh(\phi) = \exp \left(-\frac{1}{2} \phi \cdot \partial \bar{F} \cdot \phi \right) \text{Det}^{\frac{1}{2}} \frac{\partial \bar{F}}{\partial u}$$

$\partial \bar{F}$ pos. def. quad. form