

AdS/CFT T5 David Kubiznak

Finite temperature

1 Viscous fluids

We have seen in the lecture, that the first derivative correction to the perfect fluid is described by the following tensor:

$$\tau^{\mu\nu} = -2\eta\sigma^{\mu\nu} - 2\zeta\theta P^{\mu\nu}\,,\tag{1}$$

where η is the shear viscosity, ζ is the <u>bulk viscosity</u>, $P^{\mu\nu} = u^{\mu}u^{\nu} + g_0^{\mu\nu}$ is the projection operator, and

$$\sigma^{\alpha\beta} = P^{\alpha\gamma}P^{\beta\delta}\left(\nabla_{(\gamma}u_{\delta)} - \frac{1}{d-1}P_{\gamma\delta}\theta\right), \quad \theta = \nabla_{\gamma}u^{\gamma}.$$
 (2)

- a) Argue that for any CFT, we have to have $\zeta = 0$.
- b) Consider the fluid in the rest frame, $u^{\mu} = (1, 0, 0, 0)$ on the following background:

$$g^{0}_{\mu\nu} = \begin{pmatrix} -1 & 0 & 0 & 0\\ 0 & 1 & h^{0}_{xy}(t) & 0\\ 0 & h^{0}_{xy}(t) & 1 & 0\\ 0 & 0 & 0 & 1 \end{pmatrix}$$
(3)

Show that to the linear order in perturbation, we have

$$\delta\langle \tau^{xy} \rangle \sim -\eta (\nabla_x u_y + \nabla_y u_x) \,. \tag{4}$$

c) Evaluate the Christoffel symbol and go to the Fourier space to show that

$$\delta\langle \tau^{xy}(\omega, q=0)\rangle = -i\omega\eta h_{xy}^0.$$
(5)

2 CFT thermodynamics from gravitational dual

We have learned in the lecture that thermodynamic quantities of the strongly coupled $\mathcal{N} = 4$ SYM can be derived from those of the planar Schwarzschild- AdS_5 black hole (we suppress the S^5 here):

$$ds^{2} = \frac{\ell^{2}}{z^{2}} \left(-fdt^{2} + \delta_{ij}dx^{i}dx^{j} + \frac{dz^{2}}{f} \right), \quad f = 1 - \left(\frac{z}{z_{0}}\right)^{4}.$$
 (6)

This is done via the help of the standard AdS/CFT relations:

$$\frac{\ell^4}{l_s^4} = 2\lambda \,, \quad G_{10} = 8\pi^6 l_s^8 g_s^2 = \frac{\pi^4}{2} \frac{\ell^8}{N_c^2} \,. \tag{7}$$

- a) The temperature of the CFT is determined by the black hole temperature T calculate it!
- b) The entropy of the CFT is given by the black hole entropy. We can calculate it both – from the 5*d*-point of view (suppressing S_5) or the full 10*d* point of view:

$$S = \frac{\text{Area}_5}{4G_5} = \frac{\text{Area}_{10}}{4G_{10}}.$$
 (8)

Use this to derive the expression for the G_5 .

Assuming further that $x^i \in (0, L_i)$ and $V_3 = L_1 L_2 L_3$, calculate the entropy density $s = S/V_3$. Show that it is of the expected CFT form – confirming that the black hole entropy yields the entropy of the dual CFT. Moreover, show that it is proportional to N_c^2 , thence we are in the unconfined (plasma) phase.

c) By calculating the Euclidean action (employing all the counterterms) one obtains the following free energy:

$$F = -\frac{\pi^2}{8} N_c^2 V_3 T^4 \,. \tag{9}$$

Use this to calculate the pressure P, the energy density ϵ , and the entropy density s:

$$P = -\frac{\partial F}{\partial V_3}, \quad \epsilon = \frac{1}{V_3} \frac{\partial(\beta F)}{\partial \beta}, \quad s = -\frac{1}{V_3} \frac{\partial F}{\partial T}, \quad (10)$$

and compare these to your previous tutorial dimensional analysis. Thence you have shown how to calculate the thermodynamic quantities of the strongly coupled CFT at finite temperature.

3 CFT at weak coupling: free thermal gas

Let us finally calculate the entropy density of the SYM at small coupling. Recall from (for example your cosmology/statistical physics course) that we have

$$s_{\text{boson}} = \frac{2\pi^2}{45} T^3, \quad s_{\text{fermion}} = \frac{7}{8} s_{\text{boson}}.$$
 (11)

How many fermionic and bosonic dof do we have for the large N_c , $SU(N_c)$ SYM? Thence show that for the total entropy we find

$$s_{\rm free} = \frac{4}{3} s_{\rm BH} \,. \tag{12}$$

Do we have a match?