

# Accelerated black holes in de Sitter universe

Non-spinning (charged) ***C*-metric** — a solution of Einstein–Maxwell equations with positive cosmological constant  $\Lambda$  — represents a pair of accelerated black holes in asymptotically de Sitter universe.

$$g = \frac{\ell^2}{\omega^2} \left( -\mathcal{F} d\tau^2 + \frac{1}{\mathcal{F}} dv^2 + \frac{1}{\mathcal{G}} d\xi^2 + \mathcal{G} d\varphi^2 \right)$$

$$\mathbf{F} = e \mathbf{d}v \wedge \mathbf{d}\tau$$

## Interpretation of coordinates

- $\tau$  time coordinate of ‘accelerated’ observers
- $v$  radial coordinate
- $\xi$  angular coordinate measured from the axis of symmetry
- $\varphi$  angular coordinate around the axis of symmetry

## Parameters

- |        |   |
|--------|---|
| $m$    | mass parameter                                    |
| $e$    | charge parameter                                  |
| $A$    | acceleration parameter: $\sinh \alpha_o = \ell A$ |
| $C$    | conicity parameter: $\varphi \in (-C\pi, C\pi)$   |
| $\ell$ | cosmological scale: $\ell = \sqrt{3/\Lambda}$     |

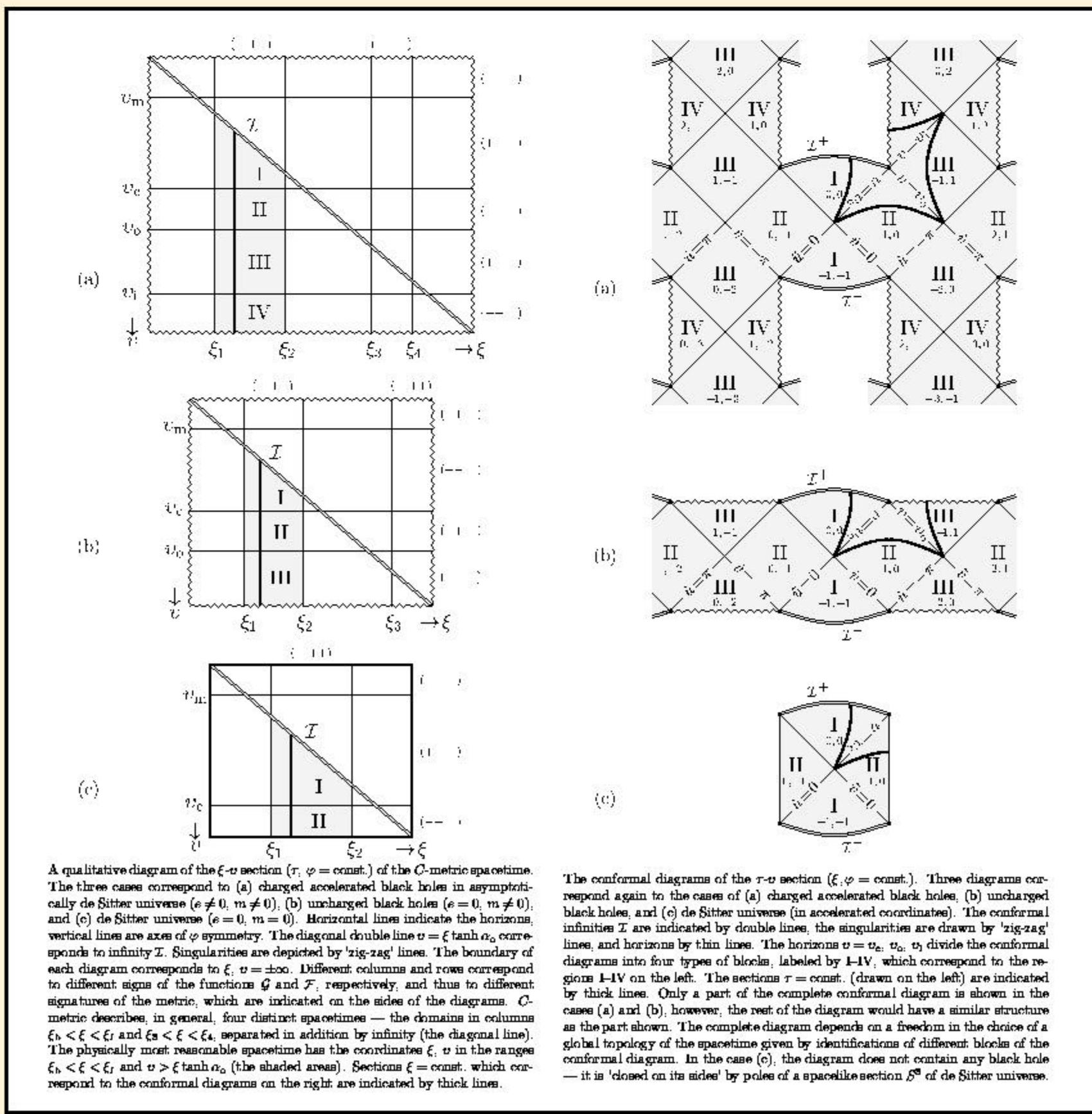
$$\begin{aligned} -\mathcal{F} &= 1 - v^2 + 2 \frac{m}{\ell} \operatorname{ch} \alpha_o v^3 - \frac{e^2}{\ell^2} \operatorname{ch}^2 \alpha_o v^4 \\ \mathcal{G} &= 1 - \xi^2 + 2 \frac{m}{\ell} \operatorname{sh} \alpha_o \xi^3 - \frac{e^2}{\ell^2} \operatorname{sh}^2 \alpha_o \xi^4 \\ \omega &= -v \cosh \alpha_o + \xi \sinh \alpha_o \end{aligned}$$

## Relation to Kinnersley–Walker form

$$\begin{aligned} g &= \frac{1}{A^2(x+y)^2} \left( -\mathcal{H} dt^2 + \frac{1}{\mathcal{H}} dy^2 + \frac{1}{G} dx^2 + G d\varphi^2 \right) \\ \mathcal{H} &= -\ell^{-2} A^{-2} - 1 + y^2 - 2mAy^3 + e^2 A^2 y^4 \\ G &= 1 - x^2 - 2mAx^3 - e^2 A^2 x^4 \\ \tau &= \coth \alpha_o t \quad v = \tanh \alpha_o y \quad \xi = -x \end{aligned}$$

## Basic properties of the spacetime

- Two Killing vectors  $\partial_\tau, \partial_\varphi$  and one conformal Killing tensor
- Two double-degenerate principal null directions lying in surfaces  $\xi = \text{const.}$  (Petrov type *D*) pointing ‘radially’ from the black holes
- Spacelike past and future conformal infinities  $\mathcal{I}^-$  and  $\mathcal{I}^+$
- Two disconnected event horizons (a pair of black holes) and one cosmological/acceleration horizon in one asymptotically de Sitter domain
- Conical singularity along the axis of symmetry – cosmic string ‘accelerating’ black holes



## Black hole coordinates

$$\begin{aligned} g &= \frac{\ell^2}{\omega^2 R^2} \left( -\mathcal{H} dT^2 + \frac{1}{\mathcal{H}} dR^2 + R^2(d\Theta^2 + G d\Phi^2) \right) \\ \mathcal{H} &= 1 - \frac{R^2}{\ell^2} - \cosh \alpha_o \frac{2m}{R} + \cosh^2 \alpha_o \frac{e^2}{R^2} \\ T &= \ell \tau \quad R = \frac{\ell}{v} \quad d\Theta = \frac{1}{\sqrt{G}} d\xi \quad \Phi = \varphi \end{aligned}$$

For vanishing acceleration  $\alpha_o = 0$  we get

$$R\omega = \ell \quad \xi = \cos \Theta \quad \mathcal{H} = 1 - \frac{R^2}{\ell^2} - \frac{2m}{R} + \frac{e^2}{R^2}$$

*C*-metric becomes the Reissner–Nordström-de Sitter or Schwarzschild-de Sitter solution, respectively.

## zeros of $\mathcal{G}$ — axes of $\varphi$ symmetry

4 zeros,  $\xi_b < \xi_f$  the smallest ones:

- $\xi_f$  axis in ‘forward’ direction
- $\xi_b$  axis in ‘backward’ direction

## zeros of $\mathcal{F}$ — horizons

4 zeros  $v_i > v_o > v_c > v_m$  (assuming  $m \neq 0, e \neq 0$ ):

- $v_i$  inner black hole horizon
- $v_o$  outer black hole horizon
- $v_c$  cosmological/acceleration horizon
- $v_m$  non-physical value

## zeros of $\omega$ — conformal infinity $\mathcal{I}$

$$v = \tanh \alpha_o \xi$$

## Few references on *C*-metric

$$\Lambda = 0$$

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